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**Faint BVRI Photometric Sequences in Selected Fields**A. Saha<sup>1,2</sup>, A. E. Dolphin<sup>1,3</sup> and F. Thim<sup>1,4</sup>*NOAO, P.O. Box 26732, Tucson, AZ 85726*

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**ABSTRACT**

The results from work done to extend the Johnson-Cousins *BVRI* photometric standard sequence to faint levels of  $V \sim 21$  mag in compact fields is presented. Such calibration and extension of sequences is necessary to fill a calibration gap, if reliable photometry from modest aperture telescopes in space (e.g. *HST*), or terrestrial telescopes with apertures exceeding 4-m is to be done. Sequences like the ones presented here, which cover a large range in brightness as well as color, will allow photometric calibration to be done efficiently, as well as for such work to be less prone to systematic sources of error.

Photometry of stars in approximately  $10 \times 10$  arc-minute fields around 3 globular clusters, NGC 2419, Pal 4 and Pal 14 are presented from data acquired over several photometric nights. In each field, several stars are measured in *B*, *V*, *R* and *I* passbands, with standard errors in the mean less than 0.015 mag from random errors, to levels fainter than  $V = 21$  mag. It is shown that standard errors in the mean from systematic errors when tying to the Landolt standards on the Johnson-Cousins system are typically well below 0.01 mag in all 4 bands (except for *B* in NGC 2419, and *R* in Pal 4), thus justifying the claim that these fields have been correctly calibrated.

The primary context for the work presented here is that parts of these fields were observed repeatedly by the Wide Field Planetary Camera-2 (WFPC2) of the *HST*, and thus these newly calibrated sequences can be used to retro-actively calibrate WFPC2 at over various times of its operating life. In the past, WFPC2 data have had typical photometric zero-point uncertainties of a few hundredths of a magnitude, largely due to a lack of suitable standard stars. The sequences presented here have standard errors at the 0.01 mag level. They agree at the 0.02 mag level with

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other extant calibrations of the targets presented here, except in the  $I$  band, where there are color dependent deviations of up to 0.05 mag versus one other photometric sequence. There is no clear resolution of this difference: we present as much verification of the sequences presented here as possible. We argue that a very likely reason for such discrepancies is differences in the filter bandpass.

*Subject headings:* globular clusters: individual (NGC 2419, Pal 14, Pal 4) — standards

## 1. Introduction

Accurate absolute photometric measurements of stars and galaxies lie at the very heart of modern astronomy. Modern day measurements, both from space with the *Hubble Space Telescope*, as well as with large aperture telescopes from the ground are routinely being used for projects that require accurate photometry. Examples range from the distance scale based on Cepheid or RR Lyrae variables or supernovae, to metallicity estimates from the position of the red giant branch, to the ages of globular clusters and dwarf galaxies throughout the Galactic halo. With each improvement in accuracy it becomes possible to attack whole new classes of problems.

It is fundamental to precision measurement, that calibrators and science targets be observed in exactly the same way to minimize systematic errors. This has not been possible for most photometric observations using the *Hubble Space Telescope (HST)* or large ground-based telescopes, due to the lack of faint standards. For instance, the photometric zeropoints for the Wide Field Planetary Camera 2 (WFPC2) on board *HST* have been monitored using exposures of a few seconds of a  $V = 13$  star, while most science observations utilize exposures of typically 500–2000 s, and are interested in much fainter targets (e.g., Cepheids from  $V = 24$  to 26).

In 1994, Peter Stetson (see Kelson et al. (1996) and Saha et al. (1996)) first noticed that his WFPC2 photometry from short and long exposures of Palomar 4 and NGC 2419 did not agree at the few percent level. This led to a more complete analysis by Casertano & Mutchler (1998)<sup>1</sup>, and by Stetson (1998) which confirmed that the WFPC2 has a non-linearity problem, which is responsible for absolute photometric uncertainty that could be as large as 0.05 mag. Thus far, zero-points reported by various investigators scatter with rms values like 0.02 mag for  $V$  and  $I$ , and 0.03 to 0.04 mag for  $B$  and  $U$  (Heyer et al. 2004)<sup>2</sup>. A brief history of attempts to understand and mitigate the anomalous behavior of WFPC2 can be found in Dolphin (2000b).

The obvious solution is to obtain fainter standards for use with *HST* and large ground-based telescopes which have similar problems. In a recent compilation, Stetson (2000) presented  $> 15,000$  new photometric standards in  $BVRI$ . While Stetson’s current database<sup>3</sup> does contain samples of standard stars in a number of globular clusters and dwarf galaxies that have been observed over the past decade with WFPC2, it is still highly desirable to increase the number density on the sky, the magnitude limit, and the accuracy of faint standard stars in a number of historically popular science fields for two reasons:

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<sup>1</sup>see [http://www.stsci.edu/instruments/wfpc2/Wfpc2\\_isr/wfpc2\\_isr9802](http://www.stsci.edu/instruments/wfpc2/Wfpc2_isr/wfpc2_isr9802)

<sup>2</sup>see [http://www.stsci.edu/instruments/wfpc2/Wfpc2\\_isr/wfpc2\\_isr0401.html](http://www.stsci.edu/instruments/wfpc2/Wfpc2_isr/wfpc2_isr0401.html)

<sup>3</sup><http://cadwww.dao.nrc.ca/cadcbn/wdb/astrocat/stetson/query>

1. Fields suitable for direct observation with *HST* instruments must have sufficient number of stars per square arc-minute to efficiently map the Quantum Efficiency (QE) and Charge Transfer Efficiency (CTE) of all the component CCD chips.
2. Observing established standards *now* with WFPC2 would not tell us anything about the historical calibration of the instrument, since it has been shown (Whitmore et al. 1999) that the calibration of WFPC2 for faint objects has evolved with time. Rather, fields that have been observed with WFPC2 at multiple epochs during its active life should be calibrated as best as possible, so that the time varying behavior of the instrument can be characterized after the fact.

We have therefore used WIYN, with its relatively large aperture and excellent seeing, to extend the Landolt *BVRI* photometric system (Landolt 1992, 1983) to fields that have already been repeatedly observed with *HST* and which are known to have a suitable density and magnitude range of stars to provide an effective retroactive calibration of WFPC2. We are able to obtain consistent photometry to levels of 0.01 mag down to  $V \approx 22$  mag in our targets. In this paper we describe our experiment and present the results of this calibration for three such suitable targets: the globular clusters NGC 2419, Pal 4, and Pal 14. We demonstrate that systematic and random errors have been contained sufficiently well so that stars in these fields have been tied to the Landolt system with accuracy  $\sim 0.01$  mag or better.

The results from this project should improve the magnitude gap between photometric standards and science targets. For instance, typical corrections for WFPC2 linearity will be over a dynamic range of  $\approx 10$  (between photometric standards and extragalactic Cepheids, for instance) rather than the prior range of  $\approx 100,000$ . These calibrated fields will also help calibrate new generations of *HST* imagers. Our photometry of NGC 2419 has already been used to verify the calibration for the *Advanced Camera for Surveys (ACS)* (Sirrianni et al. 2004) and will provide an efficient way to perform nightly photometric calibrations for imagers on 8-m class telescopes.

In the following sections we describe the observations used, followed by a discussion of the peculiarities of the instrument configurations used, and the methods required to properly process the photometry extraction that successfully removes the instrumental signatures. The photometry for four target fields (NGC 2419, Pal 14, Pal 4 and Pal 3) are presented. The exposure to exposure and night to night consistency is examined in a way that gives separate estimates for systematic and random errors in calibration. Lists of stars in each of the four fields are presented, with the measured magnitudes and error estimates. Comments on how they may be best utilized as standard stars is discussed.

## 2. Instrumental Setup

Observations were made primarily with the WIYN 3.5-meter telescope at the Kitt Peak National Observatory. The Mini-mosaic (*MIMO*) camera mounted on one of the Nasmyth focii was used. This instrument has two  $4096 \times 2048$  *SITE* CCD chips, mounted side by side, providing a roughly square field of view (FOV). The 15 micron pixels project to  $\sim 0.14$  arc-seconds on the sky at the  $f/6.4$  Nasmyth focus of the telescope. This is an excellent match to the superb image quality delivered by the telescope (median  $FWHM = 0.7$  arc-sec), and also provides an FOV of about 9.6 arc-minutes on a side. The imaging performance of this telescope and instrument combination is described in Saha et al. (2000). Each chip is read out by two amplifiers as two  $4096 \times 1024$  sub-sections.

The instrument was commissioned with rigorous tests which cover linearity, shutter timing performance,

charge transfer efficiency, dependence of accumulated counts from a star with background level, and geometrical distortion of the focal plane. The complete retinue of tests is described in the commissioning report, which can be found at the website: <http://www.noao.edu/wiyn>. The charge transfer inefficiency is too small to be noticed, even at the lowest background levels. Shutter timing and/or shading corrections are unnecessary at the 0.2% level, even for exposures as short as 0.3s. No measurable dependence of aperture photometry on position in the field of view (due to geometric distortion, for instance) could be detected. With the exception of a few issues that we mention below, the photometric performance of *MIMO* is better than 0.2%. Throughout the course of the observations reported in this paper, one of us (AS) continually monitored the performance of the instrument, to ensure that it remained satisfactory.

The only real issue of note has been the existence of a cross-talk between amplifiers: if a pixel is saturated in the A/D counter of one amplifier, the corresponding pixel in the sub-section read by each of the other amplifiers can be contaminated. This problem is described in the commissioning report, and methods to mitigate it in the data reduction process are provided in an anonymous ftp area that can be reached from the *Documentation/* tag on the website: <http://www.noao.edu/wiyn> under the sub-heading *Saha's notes on reducing MINIMO data*. Pixels that are affected by cross-talk were masked, and not used in subsequent analysis.

The spectral response of the chips from 380 nm to 850 nm varies smoothly from about 80% in the middle of this range to no worse than 50% at the ends. The response drops sharply from 850 to 950 nm, a spectral region that is within the *I*-band. The two chips are very similar in their spectral response, although that does not alleviate the need to evaluate color terms separately for each. The filters used (*BVRI*) are the so called ‘Harris set’, which closely resemble the canonical passbands. The transmission data for the actual filters used are available on the website: <http://www.wiyn.org/filters>, where they are designated *W2*, *W3*, *W4*, and *W5* (*B*, *V*, *R* and *I* respectively). The point to note is that the filters are well bounded, and have no red-leaks even beyond 1000 nm. Photometry of standard stars shows that the overall spectral response (telescope plus filters plus detector) requires only linear color terms with coefficients smaller than 0.05.

There are a few items of note regarding the telescope and instrument combination. Since this is an alt-az telescope, the pupil rotates with respect to the image plane, and any azimuthal asymmetry in pupil illumination will require flat-field correction to be a function of the image rotator position. Dome flats were obtained covering many orientations of the instrument base. The worst peak to valley differences in flat-field images taken in different image rotator positions was 0.5%. The next issue is that the image port was designed to include a field-flattener, which is also an atmospheric dispersion compensator. Unfortunately this optical component produces considerable light loss, especially in the blue, and so is not used. Over the FOV of *MIMO*, the lack of the field flattener can produce small but noticeable changes in the shape of the PSF, especially when the seeing is good. This must be accounted for when doing PSF photometry. It appears that there is also a ‘vertical’ displacement in the relative mounting of the two chips of perhaps a few microns, which is sufficient to produce a subtle discontinuous change in the PSF shape as one steps from one chip to the other. In total, the PSF changes have been confirmed to be a function only of the radial position in the field (as one would expect if the source of the problem is the absence of the field-flattener) plus a step from one chip to the next. The resulting complications for PSF fitting photometry are adequately handled if aperture corrections are determined as a function of radial position on *MIMO* in addition to a step function going from one chip to the other. The third item of note is the imperfect baffling at the Nasmyth port, which allows some stray scattered light to enter the camera. We have ascertained that this produces only an elevation of the background (‘sky’) level, and changes over spatial scales that are much much larger than a stellar PSF. As a result there is no contribution due to this on aperture or PSF photometry of stars.

It is thus clear that there are no issues with this setup that would preclude obtaining photometry with systematic errors much smaller than one percent. The telescope-instrument-site combination has delivered  $R$  band images with  $FWHM$  as small as 0.28 arc-sec (without any adaptive optics corrections) which is as well as one can expect to do from the ground. Median delivered image quality is between 0.6 and 0.7 arc-sec. In order to calibrate faint sequences within small fields of view, which is necessary for the retro-active calibration of the *HST* image archives as well as being best suited for work with current 8-m class telescopes, optimal image quality is essential to minimize confusion noise. The *WIYN/MIMO* combination is thus as good as one might expect to have. The only net disadvantage is the slow read-out time of the chips, which limits the efficiency when observing the bright Landolt standards. It would have been better to observe more standard stars than we were able to do, but on pristine photometric nights the self consistency in the photometry of standard stars was seen to be more than satisfactory, leaving us confident about the results.

In addition, our target fields and standard stars were also observed at the *WIYN* 0.9-m telescope at Kitt Peak, using S2KB, a  $2048 \times 2048$  SITE chip. Photometry of our target fields with this independent setup (telescope, detector, filters, are all different) is restricted to brighter magnitudes, but is adequate for revealing any systematic errors that arise from either telescope, i.e. concordant photometric results is a strong argument for absence of systematic errors in both setups.

## 2.1. Observations

Observations of NGC 2419, Pal 4, and Pal 14, all of which are distant globular clusters in our galaxy and have the additional virtue of having being observed repeatedly with the *WFPC2* camera of the *HST*, were made on nights that appeared to be cloudless. Contemporaneous observations of Landolt equatorial standards were also made. Pointings that include several Landolt stars within the FOV of *MIMO* were chosen, taking care also to choose fields that contain standard stars with a wide range of colors. Exposure times on the Landolt fields were set to optimize the signal-to-noise (S/N) ratios for the standard stars. Varied exposure times were used on the target fields, which allows the calibration of stars over a wider range of magnitudes, while also providing a running check on the linearity of the system (i.e. by comparing the measured magnitudes of stars that appear bright on long exposures, but faint on short ones). On some nights images were taken so that different exposures were taken with two camera base orientations set  $180^\circ$  from one another. This not only puts each star once on each chip, but also reverses the direction of charge transfer with respect to the field, thus providing closure tests not only on chip-to-chip response variations, but also on losses from charge transfer problems. Standard star fields were observed over a wide range of airmasses, so that target field observations were well bracketed. Exposure times for standard stars were always at least one second.

The true test of whether a given night was photometric was decided after reducing the standard star data. On good nights, typical rms residuals in the photometry for standard stars per measurement are  $\sim 0.01$  mag or smaller: errors in the mean are typically  $\sim 0.003$  mag. Only nights which proved to be photometrically pristine were used in the final analysis. A list of these nights, with a summary of the observations made, is given in table 1. The processed images from all the nights are being made available in the STScI archive.

### 3. Data Processing and Photometry

The raw images were pre-processed to identify and mask out known cosmetic flaws on the detectors, and saturated pixels and pixels affected by the cross-talk described in §2. A program written in IDL by one of us (AS) was used. Bias subtraction for the *MIMO* CCDs is best done on a line by line basis, which was also done as part of the pre-processing. Flat-field corrections were done in *IRAF*, using the external *mscred* package.

#### 3.1. Instrumental magnitudes of the Standard Stars

Aperture photometry of the Landolt stars were done interactively, using an IDL program written by AS. A measuring aperture of 35 pixel radius was used, which corresponds to an aperture of  $\sim 5$  arc-sec radius. This is a large aperture in pixels, as a result of the fine spatial sampling of 0.14 arc-sec per pixel, and so the standard stars need to be well exposed in order to avoid read-noise of the CCD (or shot noise from the sky) from contributing uncertainty. An initial background value was estimated from an annular aperture from 40 to 60 pixels in radius, and the value of the sky was interactively adjusted so that an aperture growth curve is flat near 35 pixel radius. This procedure works well for images with seeing  $FWHM \leq 1.8$  arc-sec, which sets the limit for the worst delivered image quality (DIQ) that can be tolerated. There is some low level light in the PSF even beyond the 35 pixel aperture boundary, but as long as the  $FWHM \leq 1.8$  arc-sec, it is entirely due to scattered light in the optics (mostly from the tertiary mirror), which does not change from one exposure to the next in the same passband. In other words, the 35 pixel aperture measures the same fraction of light from a star in different exposures with different  $FWHM$ , as long as the  $FWHM \leq 1.8$  arc-sec.

#### 3.2. PSF Photometry and Instrumental Magnitudes of the Target Stars

PSF photometry was run on the target fields with a modified version of the *DoPHOT* program (Schechter, Mateo & Saha 1993). For a given image, the PSF is not allowed to vary with position in the field of view. The private version of *DoPHOT* that was used, produces two sets of aperture magnitudes (for a range of measuring aperture radii) of the brighter relatively uncrowded stars *in isolation* (i.e. with all other objects subtracted after fitting). One set is for background value chosen so that the growth curve is flat near 15 pixel radius, and the second set is for background value chosen so that it is flat near 35 pixels. Let us designate the instrumental aperture mag at 15 pixel radius from the first set by  $m(15)$ , and that at 35 pixel radius from the second set by  $m(35)$ . The latter is equivalent to the interactively measured instrumental aperture magnitude described above for the standard stars. The advantage here is that we measure the star in isolation (with all neighbors subtracted), and do so automatically for all the stars (from several to several hundreds, depending on the target field), that have adequate S/N. The disadvantage is that an automatic algorithm for flat growth curve background determination can go awry from specious or unsubtracted features in the image. However, this problem is mitigated by the number statistics from a large number of stars.

If  $m(\textit{fit})$  is the PSF fitted magnitude, one needs simply to find the dependence of  $\textit{apcorr}_{35} = m(35) - m(\textit{fit})$  with detector position  $(x, y)$  to account for the subtle spatial variation of the PSF. This is not easy in practice, since there are not always a sufficient number of stars in the FOV for which  $m(35)$  can be measured with sufficient S/N. However, since the PSF variation is primarily due to subtle focus changes and lack of perfect field flattening over the field area, only the core of the PSF varies with position on the FOV, while

the wings of the PSF remain unchanged. Experimentation shows that the PSF variation due to position in the FOV is essentially all contained within a 15 pixel radius of the center of the stellar profile. Thus the spatial variation of the PSF can be compensated by calculating  $apcorr(15) = m(15) - m(fit)$ , and mapping it as a function of position in the FOV. There are more stars for which one can obtain  $m(15)$  with high enough S/N than there are for  $m(35)$ .

Using even a few stars for which both  $m(15)$  and  $m(35)$  are measured with high S/N, we can derive the additive offset required to go from a 15 pixel radius aperture magnitude to a 35 pixel radius aperture mag, since we only need one quantity, which is invariant over the FOV. Thus,

$$offset_{15}^{35} = \langle m(35) - m(15) \rangle \quad (1)$$

Meanwhile,  $apcorr_{15} = m(15) - m(fit)$  is fitted as a function of position  $x$  and  $y$  on the FOV:

$$apcorr_{15} = apcorr_{15}(x, y) \quad (2)$$

With experimentation, we have found that a quadratic polynomial with circular symmetry about the optical axis (center of the FOV) is adequate (a general two dimensional quadratic function has the risk of being unconstrained in the corners if there are not enough stars to define it), plus an additional term to correct the discontinuity going from one chip to the next, as described in §2. The sign and magnitude of the variation depends on the seeing, as well as on the focus setting (i.e. how far, and in which direction, the telescope is from true focus). Once the polynomial description of  $apcorr_{15}$  (in terms of  $x$  and  $y$ ), and the value of  $offset_{15}^{35}$  are known, the PSF fitted magnitude  $m(fit)$  for any star anywhere in the FOV can be put on the system of 35 pixel apertures, which we designate as the instrumental magnitude  $m^{instr}$ .

$$m^{instr} = m(fit) + apcorr_{15} + offset_{15}^{35} \quad (3)$$

This brings the target star instrumental magnitudes to a common footing with the instrumental magnitudes of the standard stars.

Software necessary to determine and apply the aperture corrections as above were custom written in the *IDL* language.

### 3.3. Nightly Photometric Solutions

To account for extinction due to airmass, and to allow for color-terms and zero-point adjustment in order to transform to  $m^{true}$ , the true magnitude on the Landolt system, we can write a system of equations of the kind:

$$m_i^{true} = m_i^{instr} + C_i + \alpha_i X + \beta_{i,j}(m_i^{instr} - m_j^{instr}) \quad (4)$$

where  $i$  and  $j$  are indices for the various passbands ( $i \neq j$ ), and  $X$  is the airmass at which the object is observed. The standard star observations are used to solve for the coefficients  $C, \alpha$ , and  $\beta$ .<sup>4</sup>  $C_i$  and  $\beta_{i,j}$  are in general different for each chip, and were solved for separately for each chip. Higher order color terms might be required, but with our setup, as described in §2, were found to be unnecessary. Also, no cross-terms of airmass with color were used. Such terms are sometimes necessary, since selective absorption by

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<sup>4</sup>Note that in this formulation, one can additionally use the instrumental mags of any stars in the object fields that were observed more than once, and at significantly different air-masses to assist in the determination of  $\alpha$ .

the atmosphere can alter the effective central wavelength for a passband, i.e. a star (particularly if it is of extreme color) observed at high airmass is seen with a different effective wavelength than if observed at low airmass. All of our target observations were made at airmass less than 1.7, and observations of Landolt stars were done up to airmasses of 2.0. The use of such cross-terms did not improve the nightly solutions: typical star by star rms scatter in the standard star solutions was  $\sim 0.015$  mag or smaller, with no discernible trends in the residuals with the product of airmass and color.

In the following set of equations, the solution for the night of 2003 Feb 9 (UT) for one of the CCD chips (using 22 measurements in each passband) is shown, illustrating typical values for the coefficients. The lower case symbols  $b, v, r, i$  denote instrumental magnitudes in units of  $-2.5 \log(ADU s^{-1}) + 30.00$ , whereas  $B, V, R, I$  denote magnitudes on the Landolt system. To derive the coefficients in the equations below, Landolt’s published photometry (Landolt 1992) was used to derive the constants and coefficients in the equations below (and for each night of observations in this program). The errors in the mean are shown in parentheses. Not all color combinations are displayed.<sup>5</sup>

$$B = b - 4.120 - 0.262X + 0.013(b - r) \quad (\pm 0.003) \quad (5)$$

$$V = v - 4.095 - 0.146X - 0.010(v - i) \quad (\pm 0.002) \quad (6)$$

$$R = r - 3.978 - 0.103X - 0.012(b - r) \quad (\pm 0.002) \quad (7)$$

$$I = i - 4.714 - 0.065X - 0.024(v - i) \quad (\pm 0.003) \quad (8)$$

Over the course of the several nights of observations, we noted quite significant changes in the extinction coefficients. The pattern in the night to night differences is that a roughly equal value is added to the extinction coefficients of *all* passbands. For instance, on the night of 2003 June 4 (UT), the extinction coefficients as above for  $B, V, R, I$  were 0.350, 0.221, 0.169, and 0.124 respectively. The extinction in all bands is uniformly increased by  $0.07 \pm .02$  relative to 2003 Feb 9. It is possible that dust, water-vapor and aerosol play a role in this, and it is wise, therefore, to restrict observations to relatively low airmass, since the scale height of these contaminants may be different from other atmospheric sources of scattering.

Once the photometry solution for a night is obtained, as above, the instrumental magnitudes  $m^{instr}$  of each star on a set of target observations (obtained from aperture correction of the DoPHOT PSF fitting magnitudes according to the procedure described in §4.2) in at least 2 passbands can be converted to the Landolt system. DoPHOT reports good estimates of the (random) measurement errors from the photon statistics, read-noise characteristics and fit residuals for each star. These were propagated along with the calibrated magnitudes for the next steps in analysis.

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<sup>5</sup>While it is common practice to express the night constants and color coefficients with the true magnitudes and colors as the independent variables, and instrumental magnitudes as the dependent terms, we have chosen to do the opposite, as in the above, for ease of computing the true magnitudes for the target stars, and because this formulation can be used to estimate atmospheric extinction by comparing the instrumental magnitudes of stars in the target fields when they are observed at different airmasses. This enables the inclusion of the lower end of airmass values not available from Landolt’s equatorial standard stars when observing from mid-latitudes. In the absence of non-linear color terms, the two forms are algebraically identical. Further, the errors in measurement, i.e. of the instrumental magnitudes, are insignificant, and the direction of the regression makes no difference. Nevertheless, we have verified this assertion by computing the regression the other way, and find there is no difference at the 0.001 mag level in how these equations are specified and calculated.



#### 4. Analysis of Errors

The procedures described above yield magnitudes that are nominally calibrated on the Landolt system for each set of exposures (in the various passbands). We see that the nightly photometric solutions themselves carry errors that are smaller than one percent, which gives us the expectation that the systematic errors are also adequately small. However, it is necessary to see how well the systematic errors are indeed contained, by comparing the results obtained across several nights of observations, particularly since extinction coefficients have been seen to vary significantly from night to night. We have data from several nights for each of our targets, including some observations made with an independent setup (0.9m telescope and a different physical set of filters). We also sometimes have several determinations for magnitudes from different exposure sets on the same night. However, there is no external estimate of how well the aperture correction procedures worked, and systematic errors from that are not separately accounted for. In the final analysis, the repeatability of the magnitudes from all of these different measurements is the true test of how well systematic errors have been contained. Also the final reported magnitudes for each star should be a mean that is adjusted and weighted appropriately from all the photometric epochs at which observations were made.

We describe here the adopted procedure, using as an example, the case of  $V$  magnitudes for the NGC 2419 field. There are 14 epochs over 7 nights. After cross-matching the stars from all the epochs, an initial weighted mean value  $\overline{V}_k$  for the  $k^{th}$  star was calculated using the inverse of the variance of individual magnitude errors reported by DoPHOT as weights. The stars that are common to all 14 epochs were then identified: in the case at hand, 74 such stars were found. If  $V_j^i$  denotes the  $i^{th}$  measurement for the  $j^{th}$  such common star, and  $e_j^i$  is the corresponding measurement error estimate reported by DoPHOT, the quantity

$$\Delta V^i = \frac{\sum_j w_j^i (V_j^i - \overline{V}_j)}{\sum_j w_j^i} \quad (9)$$

where

$$w_j^i = \frac{1}{(e_j^i)^2} \quad (10)$$

is the best estimate of the systematic zero-point offset for magnitudes measured in the  $i^{th}$  epoch, relative to the mean of all epochs. More specifically, the error in the mean from the rms scatter of the  $\Delta V^i$ 's is a measure of the systematic error in the values of  $\overline{V}_k$ , and we denote it by  $\epsilon_{sys}$ .<sup>6</sup>

$$(\epsilon_{sys})^2 = \frac{\sum_i (\Delta V^i)^2}{N(N-1)} \quad (11)$$

$\Delta V^i$  was then subtracted from the magnitude of every star as measured in the  $i^{th}$  epoch (for each of the epochs), i.e.:

$$\mathcal{V}_j^i = V_j^i - \Delta V^i \quad (12)$$

from which we recompute the final weighted mean magnitude of the  $k^{th}$  star as:

$$\overline{V}_k = \frac{\sum_i w_k^i \mathcal{V}_k^i}{\sum_i w_k^i} \quad (13)$$

where the  $w$ 's are the same as in equation 10. We also obtain the reduced chi-square for each star as:

$$\chi_{\nu,k}^2 = \frac{\sum_i w_k^i ((\mathcal{V}_k^i - \overline{V}_k)^2)}{\nu - 1} \quad (14)$$

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<sup>6</sup>We adopt the convention of reporting rms scatter with  $\sigma$ , 'errors in the mean' with  $\epsilon$ , and measurement uncertainties with  $e$ .

where  $\nu$  is the number of available observations for star  $k$ .

The weighted random error in the mean for the  $k^{th}$  star,  $\epsilon_k$ , is then given by:

$$(\epsilon_k)^2 = C \cdot \frac{1}{\sum_i w_k^i} \quad (15)$$

where  $C$  is set to unity if  $\chi_{\nu,k}^2 \leq 1$ , and  $C = \chi_{\nu,k}^2$  otherwise. This formalism deals optimally with the fact that the various exposures are of different depth, and thus a given star in general has very different signal-to-noise from one exposure to the next. Further,  $\epsilon_k$ 's are reliable estimates of the uncertainty in the individual  $\bar{V}_k$ 's, irrespective of whether the DoPHOT reported error estimates  $e_k^i$ 's are correct. In the case of a variable object, the value of  $\epsilon$  is the uncertainty in the mean brightness.

In reporting our results, we will keep the systematic errors  $\epsilon_{sys}$  separate from the random/variability uncertainties  $\epsilon$ . The latter are reported star by star, whereas the former are the same for all stars for any target field and filter. Table 2 gives the values of  $\epsilon_{sys}$  for each target and passband. Random errors will be reported in later sections regarding the individual targets.

## 5. Results

For each of 3 targets for which we have sufficient data (namely NGC 2419, Pal 14 and Pal 4) we present here the following data products:

1. A deep reference image of the target field is shown both in print, as well as a *FITS* file in the electronic version of this paper. The *FITS* files are appointed with a ‘world coordinate system’ (WCS) calibrated to J2000 coordinates on the sky. All positions of individual stars for these target fields are referred to positions ( $X$  and  $Y$  in pixels) on the corresponding image and/or their J2000 positions.
2. For each target, a table of the stars for which there are at least 3 measurements in each of the four passbands  $B, V, R$ , and  $I$ , and where  $\epsilon_k$  (as in equation 15) is less than 0.015 mag in each passband. These are the best stars to be used as standards. This table is given both in printed form, and also in the electronic version of this paper. This table is numbered for easy reference to individual stars. In using these values, one should be mindful also of the additional systematic errors tabulated for each passband and target field in Table 2.
3. A table of all available objects with measured magnitudes in at least  $V$  and  $I$  is presented as an electronic table.
4. A table identifying objects flagged as variable stars, where the selection criterion is that the star have  $\chi_{\nu,k}^2$  (as in equation 14) higher than 20 in all of  $B, V, R$ , and  $I$ , with at least 6 measurements in each passband. Also, the rms scatter ( $\epsilon_k$ ) is required to be greater than 0.05 mag in  $V, R$ , and  $I$ , and greater than 0.1 mag in  $B$ . This table is also numbered for easy reference to individual stars.

In the following sub-sections, we identify and comment on these data products target by target.

### 5.1. NGC 2419

The reference field for NGC 2419 is shown in Fig. 1. The *FITS* image supplied in the electronic version reveals much more detail, depth, and dynamic range, and shows individual stars much further into the cluster

core than the printed version.

Table 3 gives the magnitudes and uncertainties (from random errors) in four passbands for all the stars in the field (designated ‘best’ stars) for which the random uncertainty in the mean is less than 0.015 mag in all the passbands. The corresponding data for all stars for which  $V$  and  $I$  magnitudes are available are given in Table 4. Finally, the stars which are most likely variables (at least 6 available measurements in each of the four passbands, and with  $\chi^2_\nu > 20$  in each passband) are listed in Table 5.

A color-magnitude diagram (CMD) in  $V$  and  $I$  for stars in the NGC 2419 field (utilizing all the objects listed in Table 4) is shown in Fig. 2. The ‘best’ stars are marked in bold, and the variables are also identified. In addition to cluster stars, there are clearly several foreground stars, including a few ‘best’ stars with colors redder than  $V - I \sim 1.5$  mag. There are also ‘best’ stars among the blue horizontal branch stars in the cluster to colors as blue as  $V - I \sim 0.0$ . While the ‘best’ stars go no fainter than  $I \sim 20.5$  mag, taking ensemble averages for the large number of fainter stars still permits one to directly calibrate photometry with uncertainties from random errors contained below 0.01 mag.

The systematic error estimates (given in Table 2) in  $V$ ,  $R$  and  $I$  for NGC 2419 are below 0.006 mag rms. This includes 3 nights of observations with the independent setup at the WIYN 0.9-m telescope, and these small values thus strengthen our confidence that the magnitudes given in the tables in this paper are truly on the Landolt system. The systematic error estimates must be added in quadrature to the individual rms estimates for random errors. The larger systematic error of  $\sim 0.02$  mag in  $B$  is disappointing. The seeing was often poorer in  $B$ , and bright uncontaminated stars with sufficient S/N that are common to exposures that range in depth by a factor of 100, and observed in common on two telescopes of very different aperture, were few. This has resulted in poor S/N for anchoring the  $B$  photometry to the Landolt scale.

A total of 92 objects were deemed to be variables. Fig. 2 shows that the majority of these are in the location of the RR Lyrae region in the CMD, with a few that are brighter than the horizontal branch by  $\sim 0.5$  mag. A few lie along a suggestive track that leads up along where the AGB should be, with a few more yet up among the brighter red giants. There is a solitary variable among where one find SX Phe or W UMa stars. No attempt was made to cross identify with known variables.

## 5.2. Pal 4

The reference field for Pal 4 is shown in Fig. 3. As before, the *FITS* image supplied in the electronic version reveals much more detail, depth, and dynamic range.

Table 6 gives the magnitudes and uncertainties from random sources for the ‘best’ stars in Pal 4, i.e. those for which the uncertainties are less than 0.015 mag in each of the four passbands. Table 7 presents the magnitudes for all stars in the field for which at least  $V$  and  $I$  magnitudes are available. The likely variables are identified in Table 8. The systematic uncertainties for each passband listed in Table 2 for Pal 4 must be added (in quadrature) in all cases to get the true uncertainties with respect to the Landolt system.

The CMD for the target field is given in Fig 4. It is much sparser than for NGC 2419: even the numbers of foreground stars is fewer for this field, which is less than 20 degrees from the North Galactic Pole. Also, the data from the 0.9m telescope did not have enough stars with high enough S/N at levels faint enough for a good overlap with the 3.5m data. As a result, there are no calibrated stars in this field that are brighter than  $\sim 16$  mag. There is only one ‘best’ quality star bluer than  $V - I = 0.5$  and only two that are redder than  $V - I = 1.5$ . However, there are several stars in the giant branch and red clump which can be used as

photometric standards to  $I \sim 20$  mag, and fainter calibration can be had from the ensemble average of stars to  $I \sim 22$  mag.

Three variable stars are identified: from their location on the CMD, one is likely an RR Lyrae star, another an SX Phoenicis star, and the brightest object is either an AGB variable (if it is a cluster member—it does lie in the line of sight of the cluster) or a field variable of as yet unknown type.

### 5.3. Pal 14

The reference field for Pal 14 is shown in Fig. 5. Once again, the *FITS* image supplied in the electronic version reveals much more detail, depth, and dynamic range.

Table 9 presents the photometry for the ‘best’ stars, which have uncertainties less than 0.015 mag in each passband. Table 10 gives the available magnitudes for all stars for which  $V$  and  $I$  magnitudes could be measured, and Table 11 lists the likely variables. As before, systematic errors with respect to the Landolt system as listed in Table 2 must be added in quadrature to all listed uncertainties.

The CMD is given in Fig. 6. The field surrounding the cluster is more populous than for Pal 4, and the ‘best’ stars span the range of colors better. Unfortunately again, as in the case of Pal 4, the 0.9m observations do not produce enough stars with sufficient overlap to allow the calibration of standard stars brighter than  $I \sim 15.5$  mag. At the faint end, the sub-giant branch is well defined by numerous stars, and (as in the case for NGC 2419 and Pal 4) even though they do not individually qualify in the ‘best’ category, ensemble averages of several of these stars can carry the calibration well beyond  $I \sim 21$  mag.

Seven variable stars are identified, some of which are likely to be foreground field stars.

## 6. Comparison with other sequences

### 6.1. Comparison with Stetson’s sequence in NGC 2419

Stetson (2000) has presented a photometric sequence in NGC 2419 that spans a similar field size and brightness range. A comparison with this independent calibration is instructive. For NGC 2419, there is also an older unpublished sequence in *UBVRI* by L. Davis in the KPNO consortium fields Christian et al. (1985). Comparison with that is also of interest.

The comparison in all four passbands with the Stetson (2000) sequence for NGC 2419 is shown in Fig. 7: the small dots show the difference in magnitudes star by star versus the mean magnitude of the objects as reported in this paper. The annotations in the figure give the net differences (unweighted means for all objects, and the errors in the mean) in the sense of Stetson minus this work. The mean differences are larger than can be accounted by the sum (in quadrature) of the error in the mean and the systematic error estimates in Table 2 (except in the  $B$  band where the systematic error estimate is atypically high). Thus, they cannot be ignored, and must be discussed. In  $V$  and  $R$  bands, the Stetson sequence is on average brighter by  $\sim 0.015$  mag, but in the  $I$  band it brighter by almost 0.04 mag. In comparison, the 0.9m and 3.5m observations reported in this paper when taken separately, differ (in the sense of 0.9m minus 3.5m) by  $-0.005 \pm .006$ ,  $+0.020 \pm 0.004$ , and  $+0.024 \pm 0.010$  mag in  $V$ ,  $R$ , and  $I$  respectively, which are smaller than the difference between the Stetson sequence and this work. Note also, that there is no obvious brightness dependence in the difference between the Stetson scale and the work in this paper: non-linearity is not the

culprit.

The large dots in Fig. 7 show the difference between the L. Davis unpublished sequence (Christian et al. 1985) and the present work. On average the Davis sequence appears to lie midway between the Stetson and current work values. We should note that formally the Davis sequence is tied to Landolt (1983), not Landolt (1992), though differences due to this are not expected to be significant beyond a few milli-mags.

An important characteristic of the difference between the Stetson (2000) sequence (especially in  $I$ ) and the work in this paper is revealed in Fig. 8, where the differences between the two sequences is shown star by star as a function of color ( $V - I$ ). Note how both in  $V$  and especially in  $I$ , the differences are a function of color. Specifically at  $V - I \sim 0.2$  there is no difference on average in either  $V$  or  $I$ , whereas at  $V - I \sim 1.2$ , The Stetson sequence is brighter in  $I$  by  $\sim 0.05$  mag, and in  $V$  by  $\sim 0.02$  mag. There are not enough data in  $R$  to see if a color dependence exists, and in  $B$  the differences are within the systematic error estimate (which is regrettably large).

## 6.2. Comparison with other photometry in NGC 2419

The equivalent comparison of magnitude difference as a function of color against the Davis standards in NGC 2419 is shown in Fig. 9. The annotations show the net differences and level of significance similar to those given for the Stetson sequence in Fig 7. The color dependent trends seen against the Stetson sequence are not present here.

One of us (AED) has obtained photometry in  $V$  and  $I$  of stars with WFPC2 in a field south of NGC 2419, which overlaps in area with our ground based field here. The photometry were obtained using the *HSTPHOT* program described in Dolphin (2000a), using a procedure that corrects for the CTE anomalies described in Dolphin (2000b). Independent photometry zero-points and color-terms derived from photometry in  $\omega$  Cen by Walker (1994) were used. The comparison of these magnitudes with those in this paper is shown in Fig 10. The data span a smaller range in color than in Fig 8, and the blue stars show significant scatter (because the blue horizontal branch stars are faint on these relatively shallow WFPC2 exposures). Despite these shortcomings, the two sets of photometry show better agreement than the comparison with the Stetson NGC 2419 sequence. Specifically, if the color trend seen against the Stetson sequence were present, it would be revealed by these data, but no such comparably strong trend is seen.

We have private communication from Sirianni that the synthetic photometry calibration of the ACS on *HST* is in close agreement with the results presented in this paper.

The various comparisons detailed above are summarized in Table 12.

We see that there is still a lack of concordance between different independent investigations regarding the calibration of the  $I$  band at the few percent level. Systematic differences of a few hundredths of a magnitude, especially in the  $I$  band, are apparently not uncommon. Many  $I$  filters in use at various observatories allow out of band transmission in the infrared. While detectors like the S20 extended response photomultipliers were blind at near infra-red wavelengths, CCDs can have substantial response at 1.0 to 1.2  $\mu$ . At these wavelengths, the spectra of cool stars have pronounced bands, and simple color transformations (even with higher order terms) may not adequately account for out of band transmission, since color ‘excesses’ become discontinuous. In addition, as we show in the Appendix, the form of the standard color equations used in practice is a poor match to physical reality unless the color response mis-match between the employed and the original measuring systems are small, i.e. correctable by a term that depends only on the first moment

in frequency/wavelength of the spectral energy distribution.

Whether the above is in fact the reason for the difference between the Stetson sequence and the one derived here, is not established beyond doubt. In practice we see much smaller deviations because of bandpass mis-match, which arises out of the fact that the spectral energy distribution of stars vary only in highly constrained ways. Synthetic simulations with a possible extreme variant of the  $I$  band (see §7.3) shows that at most half of the discrepancy seen can be explained this way. The idea of using the 0.9m data as an arbiter, while sound in principle, is only marginally useful because of the lack of S/N at faint magnitudes. While the comparison through the HSTPHOT photometry of WFPC2 observations in the NGC 2419 field against the system of Walker (1994) is consistent with the sequence derived in this paper, it is at least a little bit circular, given that we are trying to set up a sequence to retro-actively calibrate WFPC2. The agreement of the NGC 2419 photometry presented here against the ACS photometry by Sirianni et al. (2004) (priv. comm.) as discussed above, is an endorsement of the sequence derived here. However the agreement may be viewed with some reserve since the ACS Photometry results are from a synthetic calibration.

## 7. Verification of our own photometry

As a sanity check to test that there is not some hidden error in our reduction, one of us (AED) has independently reduced observations on the night of 2003 Feb 09. All measurements were made independently, including PSF fitting, which was done using a modified version of HSTPHOT (Dolphin 2000a) and solving for extinction and color terms. The results from identical data frames were compared. The comparison in  $V$  and  $I$  bands are shown in Fig 11. The ensemble mean differences in photometry, derived from stars that are reported individually (by the respective reduction procedures) to have measurement errors less than 0.05 mag, is only a few milli-mags for both passbands. No color dependent trends are visible.

### 7.1. Testing the Assumptions in the Calibration Procedure

Our discrepancy with Stetson’s sequence begs further introspection. It can be asked whether inclusion of a quadratic term in the color equation for the  $I$  band color-equation would bring the results from our data into better agreement with Stetson sequence. In other words, have we neglected a color term that needed to be included? It can be alleged, that our sparse observations of Landolt standards (due to instrument read-out time limitations) does not permit a quadratic term to be well constrained from a single night’s worth of data. To refute this argument, we demonstrate below that while our observations are much less extensive than Stetson’s, it is adequate for asserting that the discrepancy seen is a real disagreement, and not an artifact of inadequate data.

The first argument is that  $\epsilon_{sys}$  in Table 2 value for  $I$  in NGC 2419 is .0058 mag. This value comes from 14 measurements on 7 different nights on 2 different set-ups. The extinction coefficients and color equations were derived independently for each night using only observations from the same night. Despite noting quite large changes in extinction from one night to another, and observing on two different set-ups, the value of  $\epsilon_{sys}$  is satisfactorily small. Since this is an ‘external’ estimate of the systematic error, it is very unlikely that systematic errors in calibration have been made at the 0.04 mag level.

The second argument involves a combined analysis of standard star observations made on 4 different nights when  $V$  and  $I$  observations of NGC 2419 were made with the WIYN 3.5-m telescope set-up. In this

check, we assume for simplicity that the color equations are the same for both CCD chips (their spectral responses are very close), and that they do not change from one night to another. Allowing for a different zero-point for each night, we wish to derive the color coefficients. We use the extinction coefficients for  $I$  as originally derived: recall that our original procedure corrects the instrumental magnitudes, and the atmospheric coefficients are derived not just from the standards, but from all high S/N stars (including those in our target fields) that were observed at different airmasses. This correction reduces the instrumental magnitudes to zero air-mass. An error as large as 0.04 in the extinction coefficient for  $I$  will produce relative errors between standard stars observed at the extremes of airmass range (1.15 to 1.95) of 0.03 mag. The rms scatter (per star) in determining the extinctions is seen to be at least a factor of 3 smaller, and errors in the mean from all the measured standards are thus a few milli-mags at most. Thus extinction determination cannot itself be a significant source of systematic error. Let  $i$  denote the instrumental magnitudes corrected to zero air-mass. We then write the usual expression:

$$i - I = L_0 + L_1(V - I) \tag{16}$$

which is the linear form of the color equation, and where  $V$  and  $I$  are Landolt’s values for the standard stars. According to the discussion above,  $L_0$  can change from night to night, but  $L_1$  is assumed the same for all observations of standards. Similarly, one can write the quadratic form as:

$$i - I = Q_0 + Q_1(V - I) + Q_2(V - I)^2 \tag{17}$$

We can solve these equations for the standard stars, allowing only one value each for  $L_1$  and for  $Q_1$  and  $Q_2$ , but allowing  $L_0$  and  $Q_0$  to differ from night to night. The residuals for the fit to equation 16 are shown in the top panel of Fig. 12, with data points from the four different nights coded in four different colors. The middle panel shows the residuals for the same observations for the fit to equation 17. In each of the above panels, the dashed lines shows the locus about which the points would lie had we fitted the instrumental mags of stars in NGC 2419 to the Stetson sequence. These figures demonstrate that the difference seen with respect to the Stetson sequence is real, and not the figment of inadequate data. Nor is the difference attributable to using a linear color term when a higher order term is demanded. The continuous black lines in the third panel of Fig. 12 show the fits to equations 16 and 17 (with respective  $L_0$ ’s and  $Q_0$ ’s subtracted). The dashed shows the equivalent quadratic fit, if instead of fitting Landolt standards, the instrumental mags of stars in NGC 2419 are fitted to the Stetson sequence. The red line shows the difference purchased by fitting a quadratic versus a linear color equation to the Landolt stars alone: note that in the range  $-0.1 < V - I < 2.0$ , the difference is always smaller than 0.01 mag. Contrast this with the green curve, which shows the difference between fitting a quadratic color equation to the Stetson sequence in NGC 2419 versus a linear fit to Landolt standards. Since the fit to Stetson’s sequence produces a discrepancy with the Landolt standards at a level of significance that is clear from the 3 panels of Fig. 12, the discrepancy cannot be an artifact of inadequate observations of standards. It is a real discrepancy, clearly evident in the data, even though it lacks a clear physical explanation (except possibly the discussion in the Appendix).

## 7.2. Is the subset of Landolt standards actually used skewed from the Overall Landolt System?

Another possibility is that the subset of Landolt standard stars used in our calibration is somehow skewed from the parent set of all Landolt standards. In particular, Landolt’s photo-electric measurements were made with a 14 arc-sec aperture, whereas here we used 10 arc-sec (diameter). In the present work we

preferred those fields where a number of standard stars are present within the instrument’s field of view, so that we could maximize the number of standards observed while keeping the number of exposures (and accompanying large overhead in read-out time) manageably small. Thus we preferred those fields where the chance of having another star within a 14 arc-sec aperture is increased relative to fields with relatively fewer available stars. We made all our standard star measurements interactively, and any gross cases of such object confusion were immediately apparent, and the offending object was not used further. However, no explicit procedure to guard against such an occurrence was used in a systematic manner, and it is possible, though unlikely, that systematic differences between Landolt’s measurements and ours have been introduced in this manner. It is further unlikely that such an error has occurred only in one of the four bands, and even more so that the difference is correlated with color.

Fig. 13 shows the results of the photometry of Landolt stars. Each point on the ordinate for each pass band is the residual (observed minus calculated, after fitting the photometric solution for the relevant night) for each standard star measurement (one point per star per observation). The filled circles are for observations on nights that have contributed to the calibration of NGC 2419. This is like the first panel of Fig. 12, but for each of the four bands, and using the color terms evaluated independently for each night. No perceptible systematic difference is seen for the observations relevant to NGC 2419 when compared to the rest.

In Fig. 14, the same data are shown, but only for the average residuals of objects that have been observed three or more times. The error bars show the standard errors. Note that some objects have standard errors that are much smaller than their average residual—possibly indicating that there are small but significant differences between the Landolt measurements of these stars and ours. A possible source of such differences might be the difference in aperture sizes used. The mean deviations are less than  $\sim 0.02$  mag.

There is no systematic trend with color (or brightness—not shown), and the differences, if real, are random from star to star. This is a possible indication of the inherent uncertainties in using these standard stars. The worst scatter is in the  $I$  band, but even there, any systematic color dependent trend exceeding  $\sim .01$  mag is ruled out. The same data, and also the mean residuals for stars measured less than thrice, are presented in Table 13. This table also identifies exactly which Landolt stars were used.

### 7.3. Can differences in the $I$ bandpass explain the differences in $I$ band photometry versus Stetson’s magnitudes?

One can ask if red giants in NGC 2419, which are all very metal poor ( $[Fe/H] \sim -2$ ), produce systematically different response through different  $I$  filters, as compared to the Landolt standards at the same  $V - I$  color. This can be studied by synthetic photometry. The standard  $I$  passband is described numerically by Bessell (1990), in his Table 2. The actual filter and CCD response combination used for data in this paper from *MIMO* is a close match to this standard passband: we will refer to it as the WIYN  $I$  passband. There is a wide range of  $I$  filters being vended by commercial suppliers, from filters that are really on the *Johnson* rather than Landolt (based on Cousins) system, to ones touted as “Bessell”  $I$  filters that have quite large red extensions. We have taken one such example as a ‘Bad  $I$ ’ filter. From simulations we find that the WIYN  $I$  filter deviates from the standard filter by  $\sim 0.009$  mag (fainter) for giants with  $V - I \sim 1.5$  mag and  $[Fe/H] = -2.0$  (as compared with giants of the same color with solar metallicity). The ‘Bad’ filter shows a deviation of  $\sim 0.025$  mag in the same sense with respect to the standard passband. The sense of the differences is correct for explaining the discrepancy between our photometry and that of Stetson, but



the magnitude of the predicted difference is only 0.015, whereas the observed difference is nearly 3 times larger. While we have no detailed information on the filter(s) used by Stetson, it is very unlikely that it was as bad a mis-match as the example used here for our case study. Thus, this too is a very unlikely source of the discrepancy.

## 8. Concluding Remarks

Despite the caveats raised in §6, this work is a significant contribution towards reconciling the large archive of *HST* imaging in broadband filters with the *BVRI* photometric system as realized via the Landolt standards (Landolt 1992). The targets presented here have been observed repeatedly with *WFPC2*, and comparing the instrumental photometry from images taken at different times will not only add to what we know of the temporal variation in CTE of that instrument, but will allow the unambiguous calibration of the photometric characteristics over the lifetime of the instrument. In other papers, we will present the confrontation of *WFPC2* photometry with data obtained at various times and with different reduction procedures.

This paper, and our understanding, has benefitted greatly from work done on these data by Peter Stetson, and from our subsequent discussions with him. We are indebted to Marco Sirianni and collaborators for sharing their results of the ACS calibration with us before publication. Support for this work was provided by NASA through grant HST-AR-09216.01-A from Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Inc., under NASA contract NAS 5-26555. The WIYN Observatory is a joint facility of the University of Wisconsin-Madison, Indiana University, Yale University, and the National Optical Astronomy Observatory.

### A. The Consequence of Spectral Response Mis-match between Measuring Setups

Here we examine from first principles, the impact and consequences arising from the usually encountered situation where the measuring system does not have exactly the same wavelength dependence as the system that defined the standards (in this case the latter refers to Landolt’s original setup). We show that the form of the color equation that is usually used does not follow the physical demands, except in the case when the difference in frequency/wavelength response is small.

Consider a source with an energy distribution  $f_\nu$ . Let the overall (telescope, filter and instrument response) response of the original (or standard) system be denoted by  $\alpha_\nu^S$ , and that of the measuring system being used by  $\alpha_\nu$ . The response of the original system is then

$$S^S = \int \alpha_\nu^S f_\nu d\nu \tag{A1}$$

and that of the ‘current’ measuring system is

$$S = \int \alpha_\nu f_\nu d\nu \tag{A2}$$

We can write

$$\alpha_\nu = C_\nu \alpha_\nu^S \tag{A3}$$

where  $C_\nu$  tracks the changes in response of the measuring system with respect to the original standard system. If  $\nu_0$  is the central frequency of the passband under consideration, we can expand  $C_\nu$  as a Taylor

series about  $\nu_0$ :

$$C_\nu = C_0 + C_1(\nu - \nu_0) + C_2(\nu - \nu_0)^2 + \dots \quad (\text{A4})$$

And so we can write:

$$S = C_0 \int \alpha_\nu^S f_\nu d\nu + C_1 \int \alpha_\nu^S (\nu - \nu_0) f_\nu d\nu + C_2 \int \alpha_\nu^S (\nu - \nu_0)^2 f_\nu d\nu + \dots \quad (\text{A5})$$

In magnitudes, the left hand side is the instrumental magnitude on the ‘current’ measuring setup, and the first term on right hand side (RHS) is a constant (zero-point adjustment) plus the true standard magnitude. In the absence of any color dependent terms (i.e. further terms on the RHS) we would just get (by taking the logarithm):

$$m_{observed} = \text{Offset} + m_{standard} \quad (\text{A6})$$

The  $n$  – *th* moment of the spectral energy distribution about the central wavelength of the passband is denoted by  $\mu^n$ , and given by:

$$\mu^n = \frac{\int \alpha_\nu^S (\nu - \nu_0)^n f_\nu d\nu}{\int \alpha_\nu^S f_\nu d\nu} \quad (\text{A7})$$

Equation (A5) can thus be re-written as:

$$S = C_0 \left( \int \alpha_\nu^S f_\nu d\nu \right) (1 + c_1 \mu_1 + c_2 \mu_2 + \dots) \quad (\text{A8})$$

where  $c_r = C_r/C_0$ .

Contrast the above equation, which is derived from physical principles, to the *ad hoc* equation that is used in practice to describe the response variation, which is:

$$m_{observed} = \text{Offset} + m_{standard} + B(\text{color}) + C(\text{color})^2 + \dots \quad (\text{A9})$$

The form of equation (A9) can be strictly derived from (A8) only in the case where when  $c_1 \mu_1 < 1$  and  $c_r \mu_r = 0$  for  $r \geq 2$ .<sup>7</sup> One must also assert that *color* in eqn (A9) is proportional to  $\mu_1$  (which behaves like a change in the effective wavelength for the passband). The failure of either of these conditions will make eqn (A9) inadequate (to varying degrees, depending on the specifics of the situation).

It should be evident that the *ad hoc* form in common use (eqn A9) is thus applicable only if the two measuring systems are a very near match. If the conditions mentioned above are not satisfied, fitting eqn (A9) is tantamount to fitting the wrong function. Depending on the severity of the mis-match, this could result in systematic color-dependent errors in the photometry. If the entire range of colors is evenly sampled, one should notice an increase in the systematic scatter when solving for the night constants using the standard stars, as the severity of the mis-match increases. This scatter is systemic in nature, and will not be mitigated by observing a large number of standards. If the range of colors sampled by the standards is patchy, the fitted solution could ‘run away’ in the unsampled regions of color.

In the presence of out of band leaks, the mis-match could be acute, since one expects the color moments (of progressively higher order) to be large as a consequence of the large moment ‘arms’. In addition, in such situations color will not remain proportional to the first moment of the energy distribution. In practice, however, for stellar work in particular, the progression of the spectral energy distribution (SED) with changing physical parameters in the photosphere is very constrained. This is manifest in the way that

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<sup>7</sup>The McLaurin series expansion of  $\ln(1+x)$  can then be applied to get the log of  $(1 + c_1 \mu_1)$ .

two color diagrams of stars define very constrained loci. For this reason, the standard form of the color equations in fact work “better than they should.” However trouble should be anticipated for non-stellar sources, and in regions of the spectrum where stellar SEDs have large dispersion, e.g. at the Balmer Jump.

The above is the mathematical persuasion that filters being used in combination with CCDs to reproduce a photometric system established with photo-cathodes be very carefully selected, and that they be free of red leaks. We believe that the set used in the study presented in this study is as close an approximation as practical. Further, the *V* and *I* filters used here are a very good match to the *F555W* and *F814W* filters used with the WFPC2 on *HST*.

This analysis strongly urges use of bandpasses that be reliably realized by filters that are well bounded so that they work with pan-chromatic detectors. This is in contrast to *BVRI*, where some of the bands relied on the detector or atmosphere for blocking, since such bandpasses are often hard, if not impossible, to reproduce when using detectors with very different response, or when they are used in space above the atmosphere. Model spectral energy distributions, such as for isochrones and for composite colors of galaxies can be synthesized for any passband. However, the rich empirical legacy of *BVRI* observations—for instance of Period-Luminosity relations for Cepheids—will not be so easily transferred to a new photometric system, and so the problem of designing physical filters and calibrating photometry in this ‘arcane’ system will continue to be of importance.

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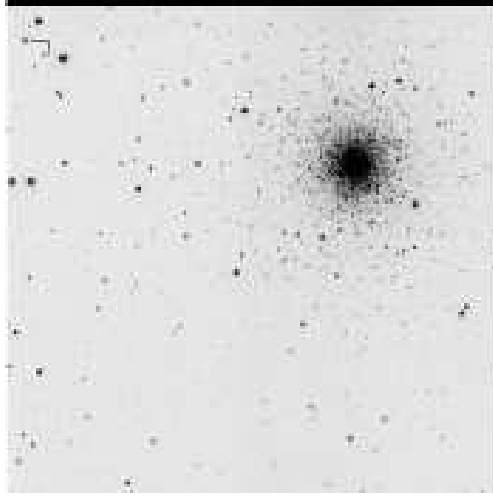


Fig. 1.— Chart of the target field for NGC 2419. North is to the left and East is down. A *FITS* file of this  $4150 \times 4100$  pixel image is given in the electronic edition of the *PASP* and at [ftp://taurus.tuc.noao.edu/pub/saha/Photseq/fg1\\_elec.fits](ftp://taurus.tuc.noao.edu/pub/saha/Photseq/fg1_elec.fits). The positions in pixels in the *FITS* image correspond to the  $X$  and  $Y$  positions of given in the tables for NGC 2419 objects in this paper. The *FITS* image is also appointed with a world coordinate system (WCS) so that sky coordinates can be read directly from the image (using a suitable display program).

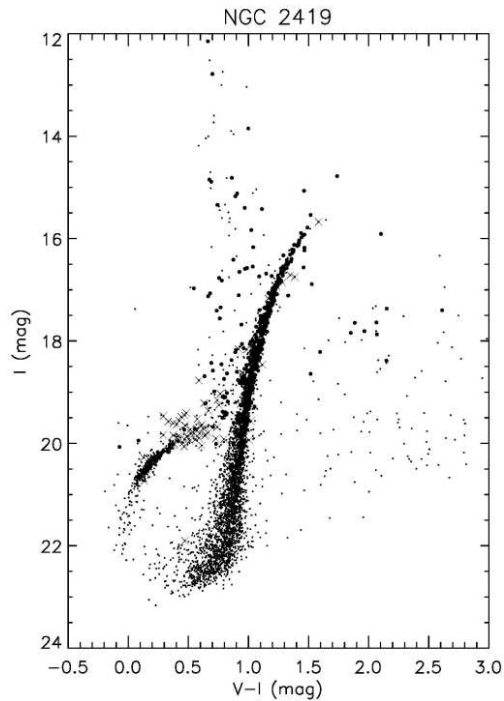


Fig. 2.— The color-magnitude diagram in  $V$  and  $I$  for the NGC 2419 field. Bold points mark objects with photometric uncertainty better than 0.015 mag in all passbands. Crosses mark variable stars. Note that the ‘best’ stars span a range of over 8 magnitudes, and a color range of over 3 mags in  $V - I$ .

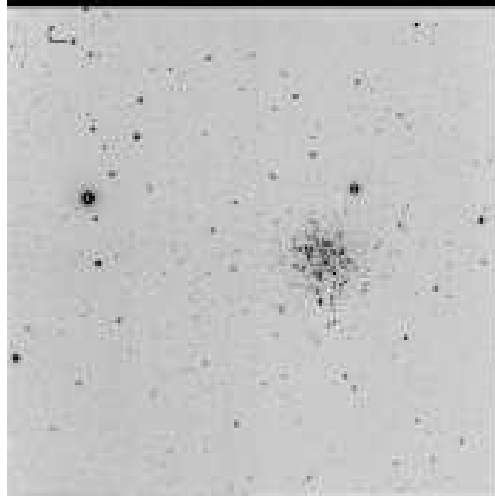


Fig. 3.— Chart of the target field for Pal 4. North is to the right and East is up. A *FITS* file of this  $4150 \times 4100$  pixel image is given in the electronic edition of the *PASP* and at [ftp://taurus.tuc.noao.edu/pub/saha/Photseq/fg3\\_elec.fits](ftp://taurus.tuc.noao.edu/pub/saha/Photseq/fg3_elec.fits). The positions in pixels in the *FITS* image correspond to the  $X$  and  $Y$  positions of given in the tables for Pal 4 objects in this paper. The *FITS* image is also appointed with a world coordinate system (WCS) so that sky coordinates can be read directly from the image (using a suitable display program).

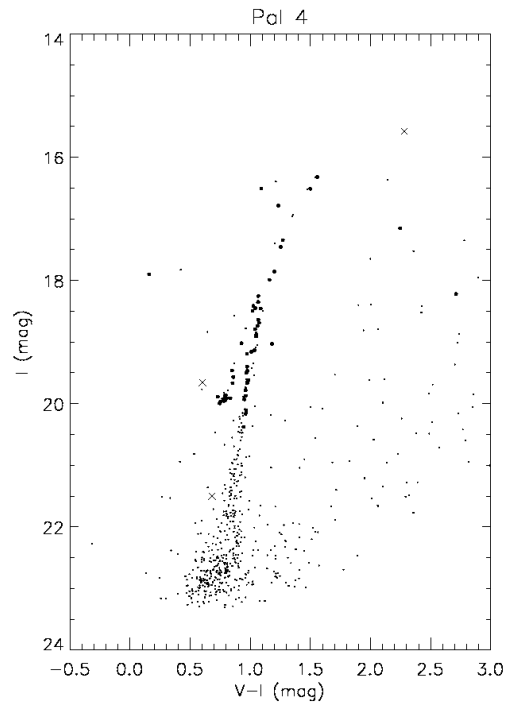


Fig. 4.— The color-magnitude diagram in  $V$  and  $I$  for the Pal 4 field. Bold points mark objects with photometric uncertainty better than 0.015 mag in all passbands. Crosses mark variable stars. Note that the ‘best’ stars span a range of nearly 5 magnitudes.

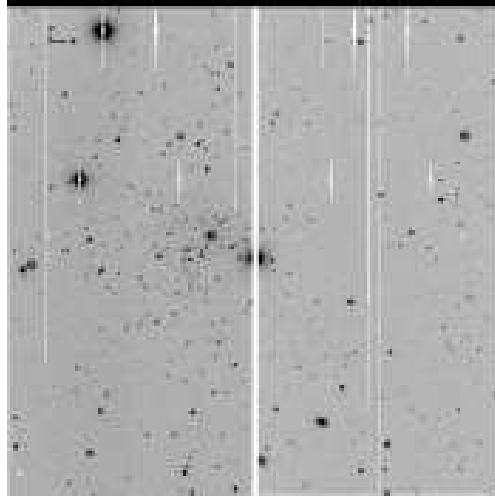


Fig. 5.— Chart of the target field for Pal 14. North is to the right and East is up. Due to the desire to position the bright star near the cluster between the two chips of MIMO, no dithering was done in the ‘X’ direction, and so this stacked deep retains the masked areas. A *FITS* file of this  $4150 \times 4100$  pixel image is given in the electronic edition of the *PASP* and at [ftp://taurus.tuc.noao.edu/pub/saha/Photseq/fg5\\_elec.fits](ftp://taurus.tuc.noao.edu/pub/saha/Photseq/fg5_elec.fits). The positions in pixels in the *FITS* image correspond to the *X* and *Y* positions of given in the tables for Pal 14 objects in this paper. The *FITS* image is also appointed with a world coordinate system (WCS) so that sky coordinates can be read directly from the image (using a suitable display program).

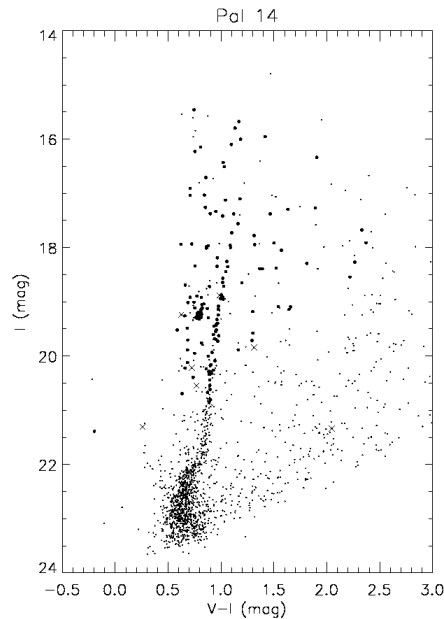


Fig. 6.— The color-magnitude diagram in *V* and *I* for the Pal 14 field. Bold points mark objects with photometric uncertainty better than 0.015 mag in all passbands. Crosses mark variable stars. Note that the ‘best’ stars span a range of over 5 magnitudes.

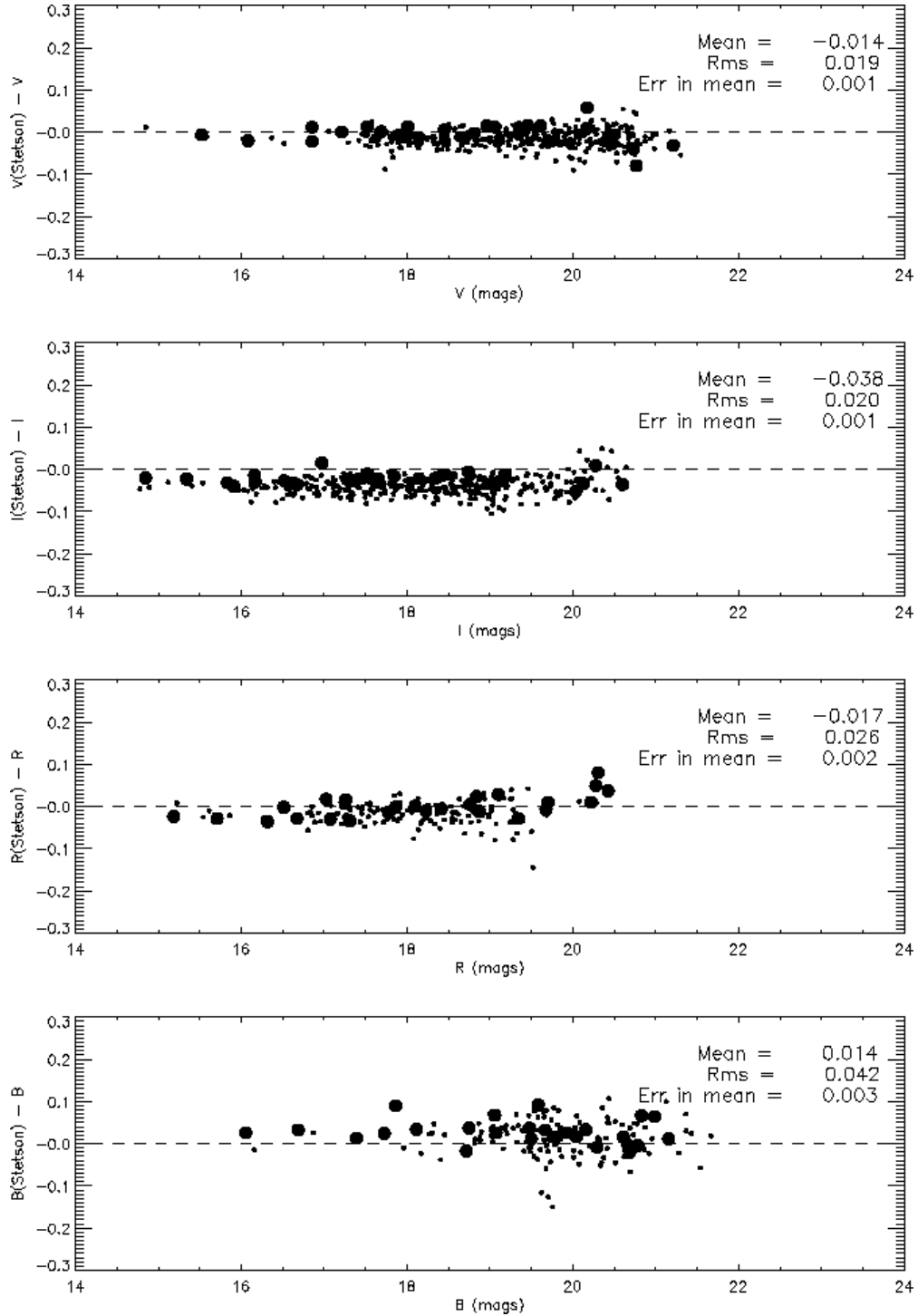


Fig. 7.— The star by star differences (shown as small dots) in all four passbands of the magnitudes of the ‘best’ stars from this study versus the magnitudes reported in Stetson (2000). The differences are shown as a function of object brightness. No obvious trends with brightness are seen. The unweighted mean differences with uncertainties are shown in the annotations. The large dots show the differences versus the Christian et al. (1985) sequence.



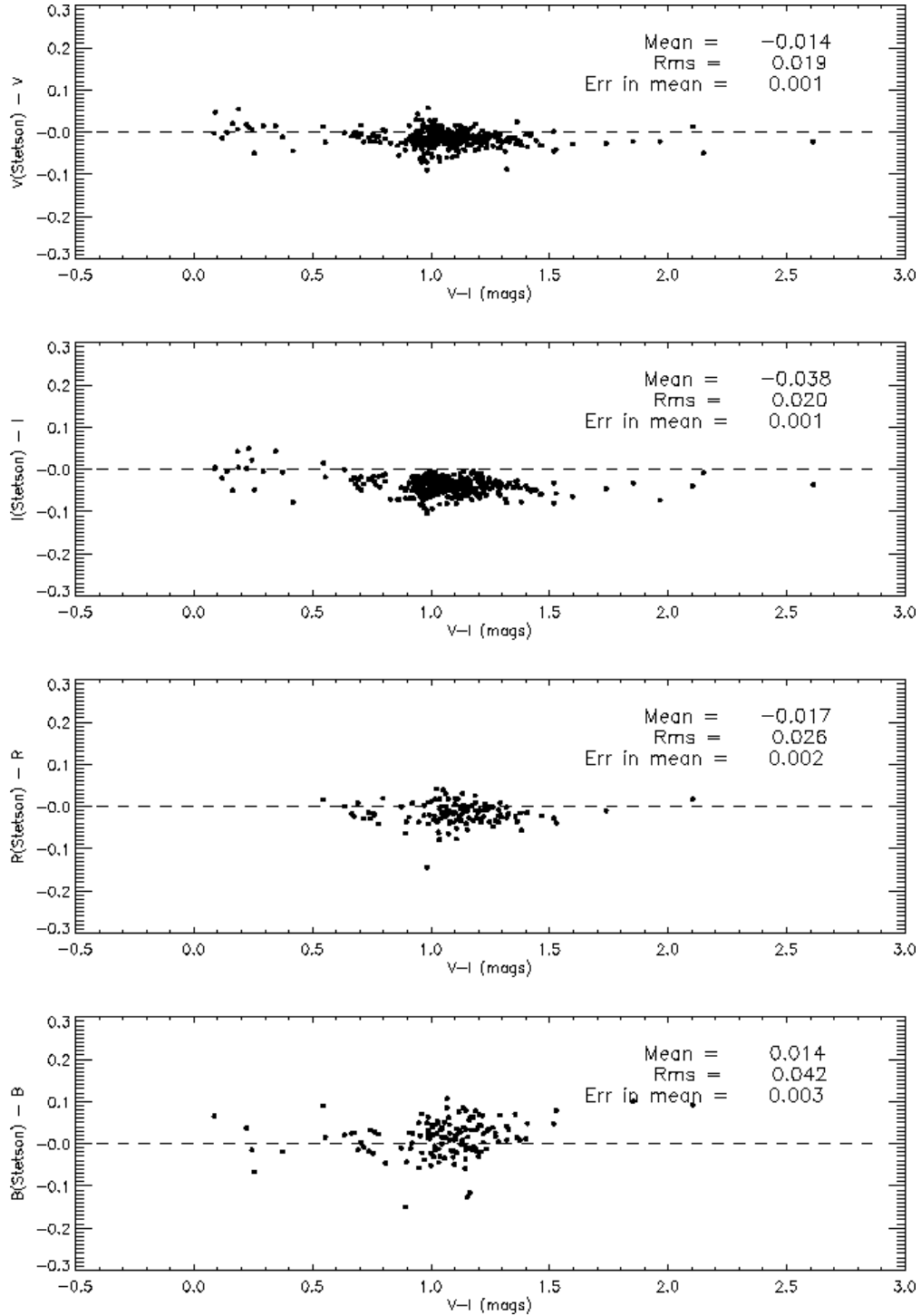


Fig. 8.— The star by star differences in all four passbands of the magnitudes of the ‘best’ stars from this study versus the magnitudes reported in Stetson (2000). The differences are shown as a function of object color. Clear trends are seen in  $V$  and  $I$ , especially in  $I$ , indicating that the source of the differences is likely in how well the two observation sets duplicate the original passband.

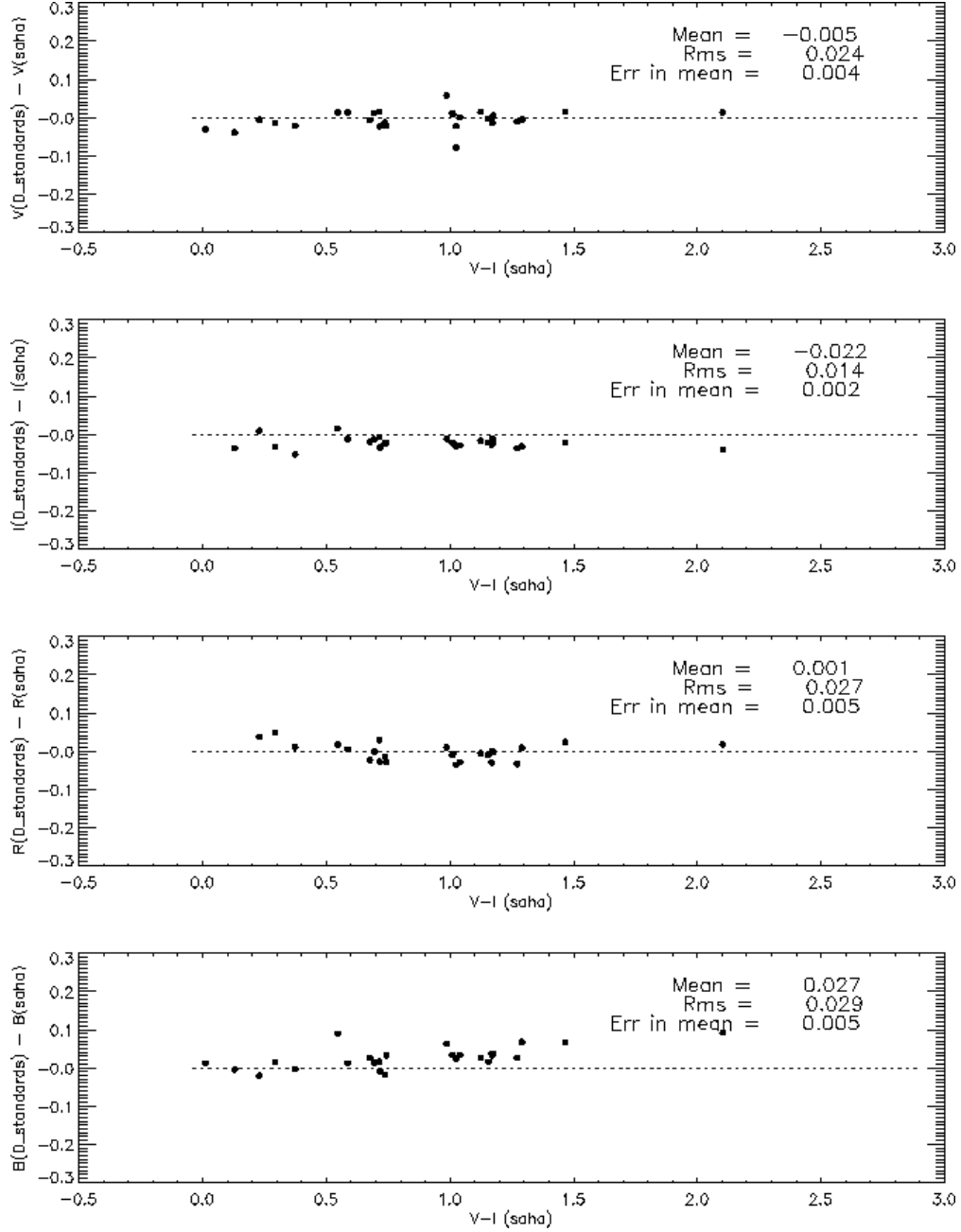


Fig. 9.— Same as Fig 8, but the comparison is against the unpublished sequence by Davis for the Kitt Peak consortium (see text). The trends seen for  $V$  and  $I$  in Fig. 8 are not seen here, but a trend appears for the  $B$  band.

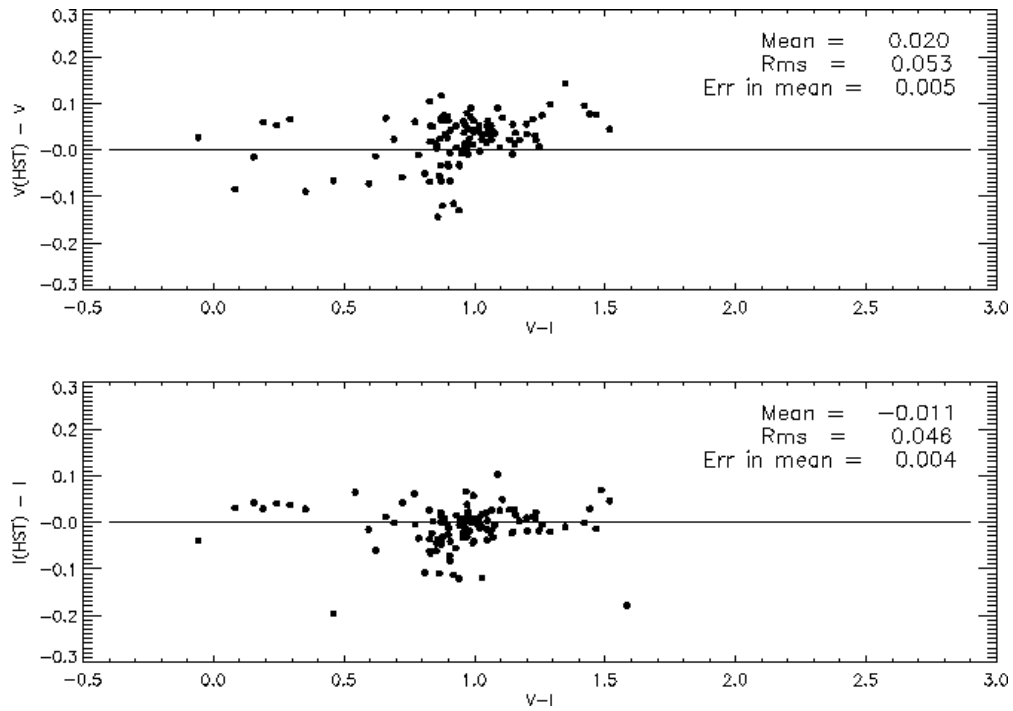


Fig. 10.— Comparison of  $V$  and  $I$  magnitudes against photometry with WFPC2 of stars in an area just south of NGC 2419. The *HST*/WFPC2 data were reduced with HSTPHOT (see text for details) using independent zero-point and color terms. The color trends seen in Fig 8 are not discernible here.

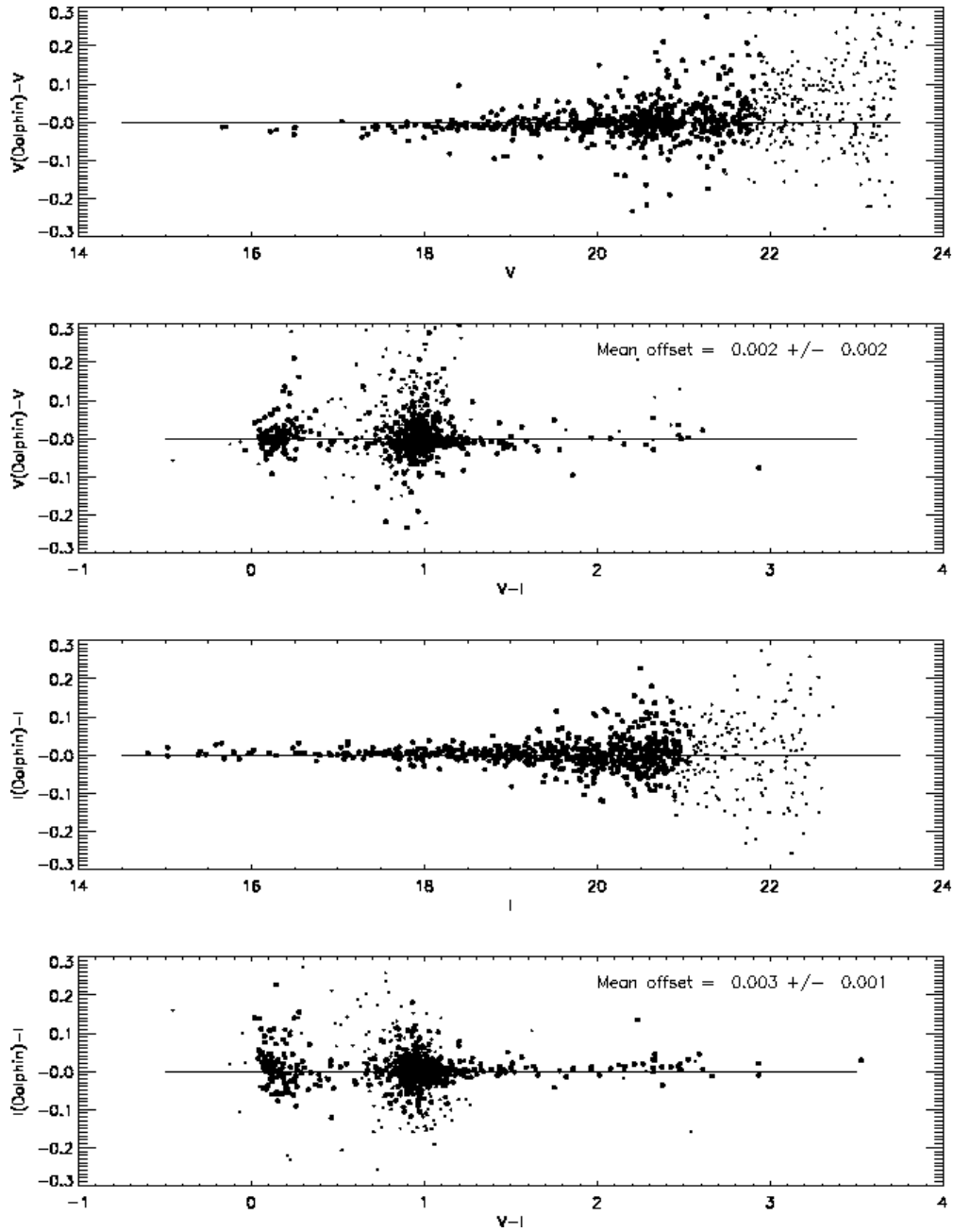


Fig. 11.— Comparison of  $V$  and  $I$  magnitudes from two different independent reduction methods applied to images obtained on 2003 Feb 09.  $V$  and  $I$  refer to magnitudes obtained with the methods described in this paper, whereas  $V(Dolphin)$  and  $I(Dolphin)$  are magnitudes derived using independent methods of PSF fitting, atmospheric extinction estimation, and color equation determination by one of us (AED). The excellent agreement corroborates that systematic errors arising from differences in methodology are insignificant.

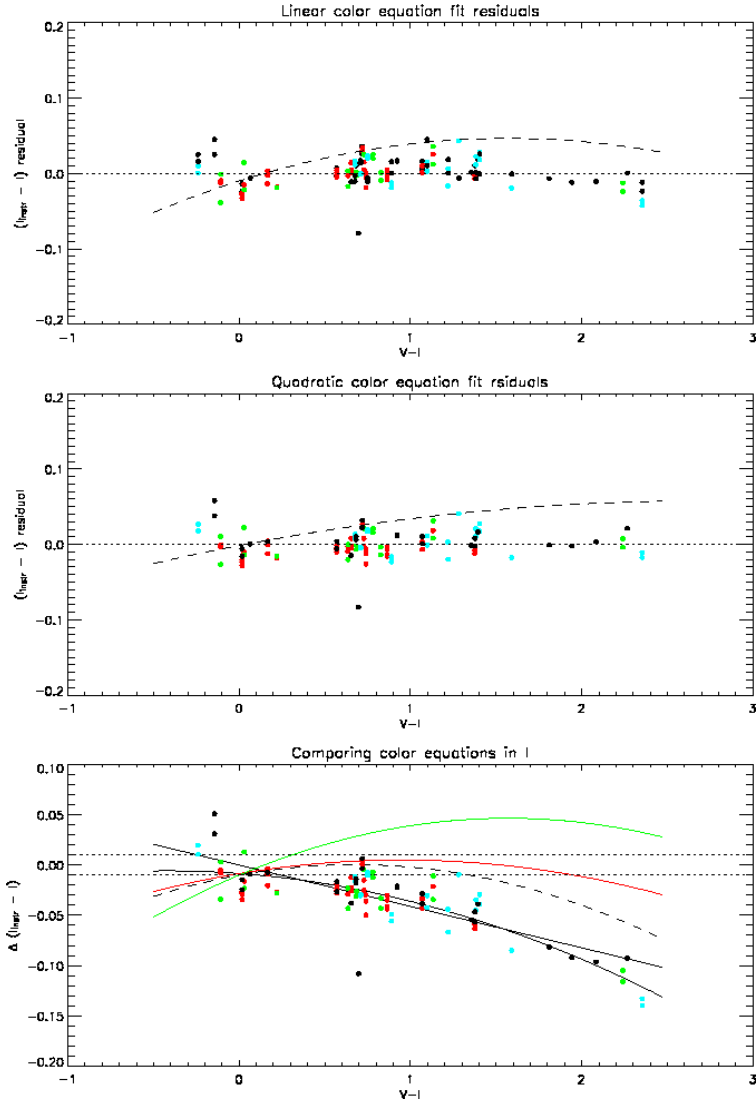


Fig. 12.— Residuals from simultaneously fitting the color equations to standard star data from 4 different nights in the manner described in §7 are shown in the first two panels: the top panel shows the case for a linear fit, and the second panel for a quadratic fit. The various colors mark data from different nights. The dashed line in the first two panels show how much the fit residuals would have to change systematically with color to produce concordance with Stetson’s sequence in NGC 2419. The scatter in the standard star data is sufficiently small to rule out the possibility that the discrepancy is due to uncertainties in our determination of the color correction. In the lowest panel, we show the color equations themselves: the straight and curved continuous lines are our best linear and quadratic solutions respectively; the dashed line is the color equation required to bring our  $I$  band observations in NGC 2419 into agreement with Stetson’s sequence; the red line shows the difference between our quadratic and linear fits and shows that for  $-0.1 < V - I < 2.0$  the difference is less than 0.01 mag, and the green line shows the difference required from our linear fit to match Stetson’s NGC 2419 sequence. More details are in §7.

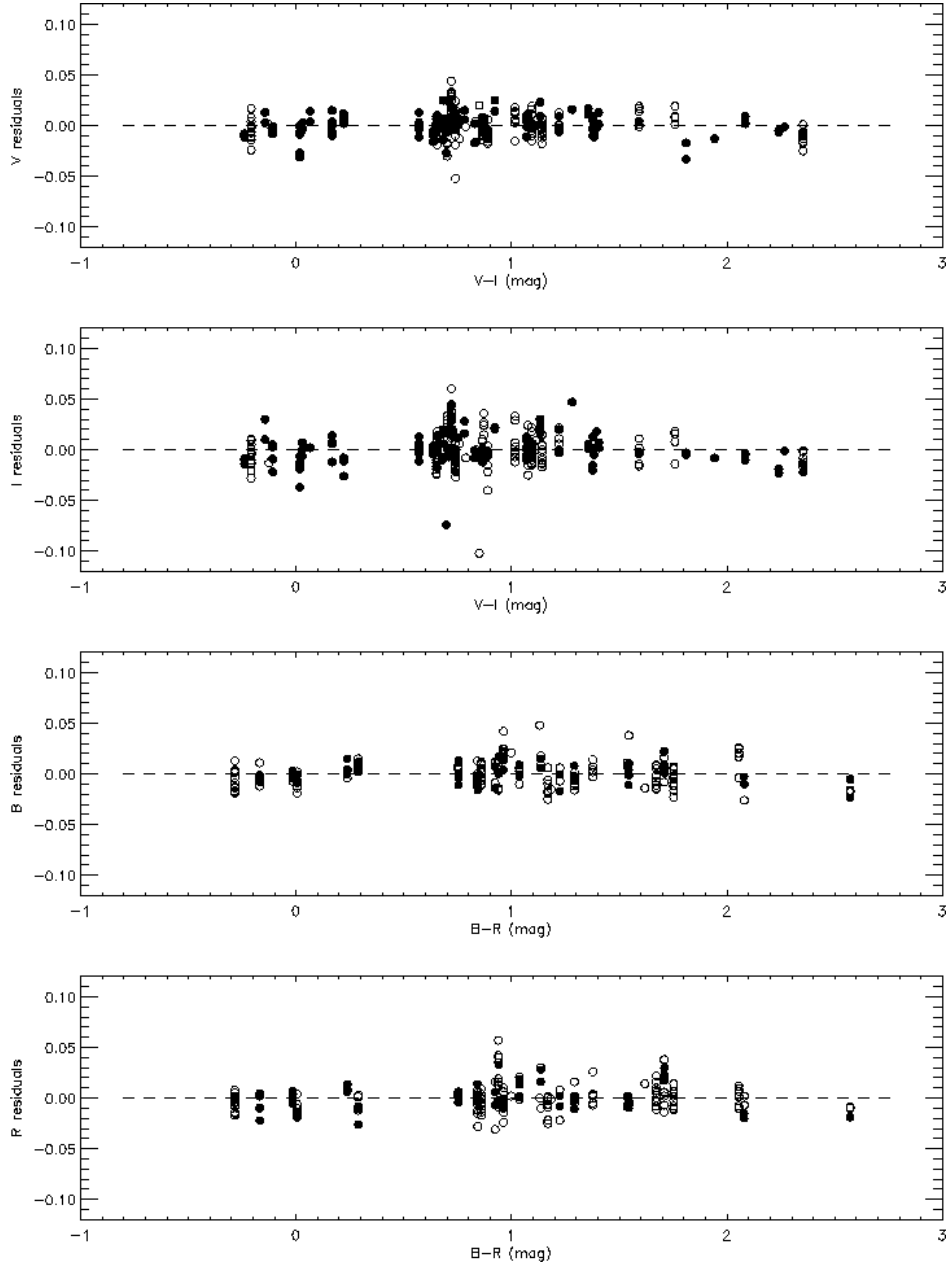


Fig. 13.— Residuals in all 4 passbands for the Landolt stars observed on all photometric nights are shown. The residuals are with respect to the predicted value using the photometric solution for the night in question. Filled circles show those observations that contribute specifically to the calibration of NGC 2419, while open circles contributed only to Pal 4 and/or Pal 14.

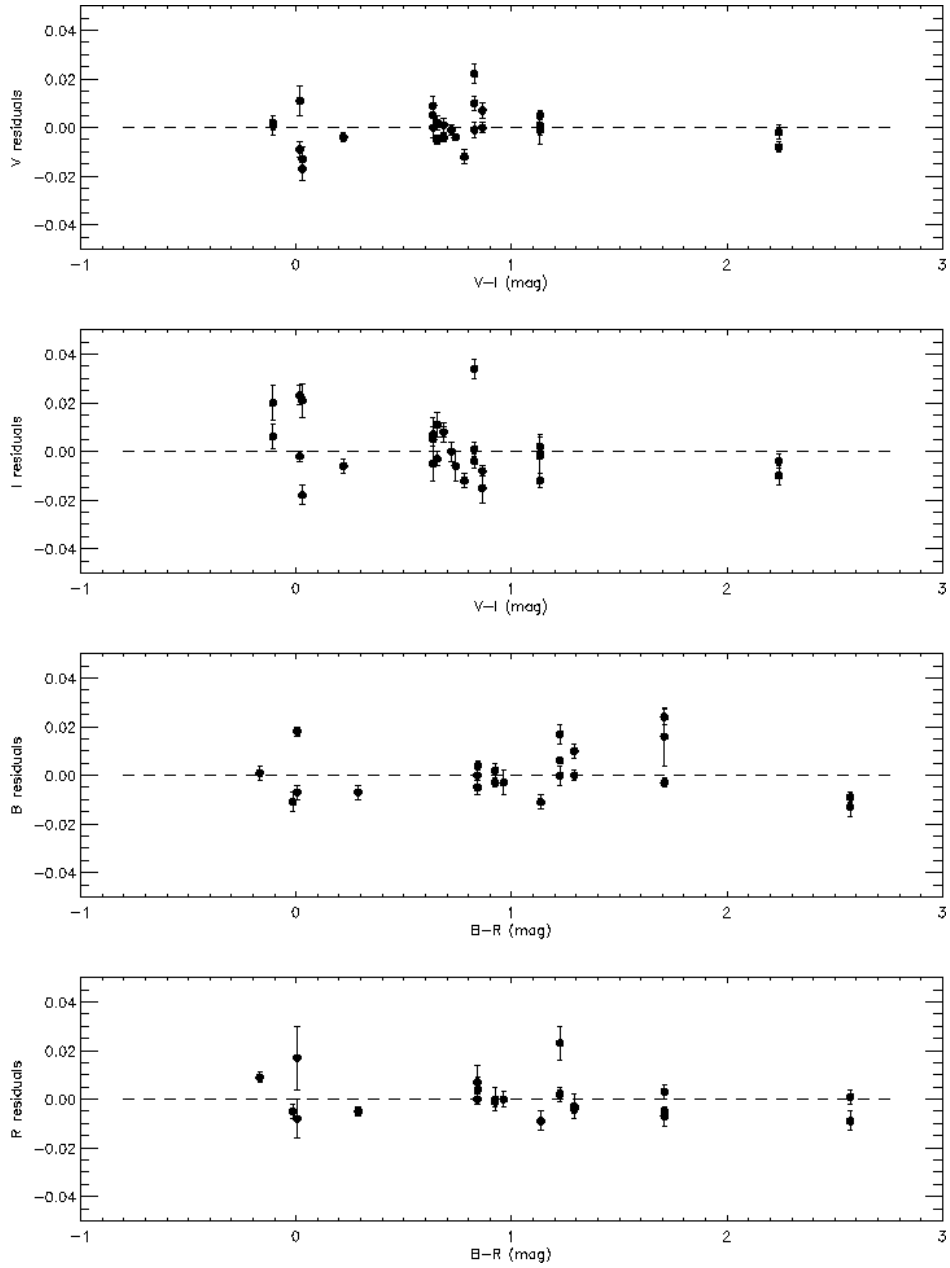


Fig. 14.— Same as Fig. 13, but showing the mean and standard errors in the residuals for only those Landolt standards observed thrice or more. Note the lack of any overall trends with color, although individual stars have mean residuals as large as 0.02 mag.

Table 1. Summary of Photometric Observations

UT Date	Telescope	Targets Observed <sup>a</sup>	Bandpasses	Landolt Fields <sup>b</sup>	Comment
2001 Feb 27	WIYN 3.5m	NGC 2419, Pal 4	<i>BVRI</i>	Rubin 149, SA101-326	–
2001 Apr 13	WIYN 3.5m	Pal 14	<i>VI</i>	PG1633+099	–
2001 Sep 21	WIYN 3.5m	NGC 2419	<i>VI</i>	SA92-430	–
2001 Dec 24	WIYN 3.5m	NGC 2419	<i>VI</i>	SA95-275, SA98-671 Rubin 152	–
2002 Apr 12	WIYN 3.5m	Pal 3, Pal 14	<i>BVRI</i>	Rubin 149, PG0918+029 PG1633+099, SA110	Photometric only in 2nd half of night
2002 Nov 13	WIYN 0.9m	NGC 2419, Pal 3, Pal 4	<i>BVRI</i>	SA92-A, Rubin 149	–
2002 Nov 14	WIYN 0.9m	NGC 2419, Pal 3, Pal 4	<i>BVRI</i>	SA92-A, SA98 Rubin 149, PG1047+003	–
2002 Nov 15	WIYN 0.9m	NGC 2419, Pal 3, Pal 4	<i>BVRI</i>	SA92-A, SA98 Rubin 149, PG1047+003	–
2003 Feb 09	WIYN 3.5m	NGC 2419, Pal 3, Pal 4	<i>BVRI</i>	SA98-671, Rubin 149 SA104-334	–
2003 Jun 02	WIYN 3.5m	Pal 14	<i>BVRI</i>	PG1633+099, SA110-232	–
2003 Jun 03	WIYN 3.5m	Pal 3, Pal 4	<i>VI</i>	PG1633+099, SA110-232	–
2003 Jun 04	WIYN 3.5m	Pal 4	<i>BVRI</i>	PG1633+099, SA110-232 SA104-336	–
2003 Jun 06	WIYN 3.5m	Pal 4, Pal 14	<i>BVRI</i>	SA104-336, PG1633+099	–

<sup>a</sup>Targets were often observed a multiple number of times, with varied positions on the field of view and/or at different air-masses

<sup>b</sup>Several stars around indicated object and within the field of view are Landolt standards. Each field may have multiple observations at different airmasses



Table 2. Summary of Photometric Observations

Target	Passband	$\epsilon_{sys}^a$ (mags)	No. of stars <sup>b</sup>
NGC 2419	V	0.0032	74
	I	0.0058	43
	R	0.0043	87
	B	0.0203	25
Pal 14	V	0.0050	153
	I	0.0057	158
	R	0.0106	262
	B	0.0046	248
Pal 4	V	0.0123	122
	I	0.0072	114
	R	0.0070	122
	B	0.0061	91

<sup>a</sup>See §4 for details

<sup>b</sup>This is the number of objects that have measured magnitudes in this passband for *all* exposures of this target: see §4





















Table 3—Continued

Star ID	X (pix)	Y (pix)	R.A. (J2000)	Decl. (J2000)	$\langle V \rangle$	$\epsilon(V)$	$n_V$	$\langle I \rangle$	$\epsilon(I)$	$n_I$	$\langle R \rangle$	$\epsilon(R)$	$n_R$	$\langle B \rangle$	$\epsilon(B)$	$n_B$
540	3820.864	3265.129	7:38:02.39	38:50:47.6	19.784	0.004	12	18.761	0.005	13	19.276	0.004	12	20.636	0.011	14
541	3834.022	2872.601	7:38:07.11	38:50:45.8	19.832	0.006	12	18.826	0.008	11	19.321	0.005	12	20.652	0.009	13
542	3837.032	3014.499	7:38:05.40	38:50:45.4	18.022	0.003	14	16.564	0.003	14	17.230	0.006	13	19.296	0.005	16
543	3851.007	2900.919	7:38:06.77	38:50:43.4	19.125	0.005	14	18.031	0.003	14	18.578	0.003	13	20.052	0.011	14
544	3872.051	3510.678	7:37:59.44	38:50:40.4	20.048	0.009	13	19.046	0.008	13	19.554	0.005	12	20.861	0.012	12
545	3881.709	3948.466	7:37:54.18	38:50:38.9	19.966	0.008	12	18.953	0.006	11	19.470	0.005	10	20.783	0.014	12
546	3883.491	3185.726	7:38:03.35	38:50:38.8	19.940	0.005	14	17.871	0.004	14	18.939	0.006	12	21.471	0.011	10
547	3889.376	2959.990	7:38:06.06	38:50:38.0	20.747	0.005	10	19.777	0.010	10	20.250	0.006	12	21.514	0.007	12
548	3892.067	2162.206	7:38:15.65	38:50:37.8	20.872	0.006	8	19.925	0.007	11	20.420	0.007	10	21.648	0.010	11
549	3899.065	2573.272	7:38:10.71	38:50:36.7	20.053	0.007	12	19.054	0.006	11	19.556	0.004	12	20.881	0.014	13
550	3920.677	2215.085	7:38:15.02	38:50:33.8	20.524	0.007	10	19.556	0.005	13	20.045	0.005	11	21.308	0.012	10
551	3936.133	2672.994	7:38:09.52	38:50:31.5	18.228	0.003	12	17.004	0.004	14	17.607	0.005	13	19.309	0.008	16
552	3978.136	2954.566	7:38:06.13	38:50:25.6	20.335	0.008	10	19.371	0.006	11	19.849	0.005	11	21.092	0.014	13
553	4002.763	3596.358	7:37:58.42	38:50:22.0	20.416	0.012	12	19.459	0.007	9	19.940	0.007	10	21.185	0.012	11
554	4018.465	2750.180	7:38:08.59	38:50:19.9	18.444	0.003	11	17.114	0.004	12	17.735	0.005	11	19.567	0.005	15
555	4035.005	3960.235	7:37:54.05	38:50:17.3	19.730	0.004	11	18.703	0.007	9	19.235	0.007	9	20.569	0.008	11
556	4058.832	2722.466	7:38:08.93	38:50:14.3	19.511	0.007	11	18.467	0.006	12	18.990	0.004	10	20.380	0.008	14
557	4081.384	3073.213	7:38:04.72	38:50:11.0	20.210	0.005	10	19.209	0.008	10	19.709	0.007	10	21.012	0.009	10
558	4121.021	3384.559	7:38:00.98	38:50:05.4	20.965	0.009	8	20.039	0.008	7	20.504	0.007	8	21.703	0.013	11
559	4124.625	3577.061	7:37:58.66	38:50:04.8	18.621	0.003	10	17.680	0.006	11	18.138	0.007	9	19.487	0.008	13



<sup>a</sup>The full text of this table may be found at: <ftp://taurus.tuc.noao.edu/pub/saha/Photseq/table4.txt>



Table 5—Continued

Star ID	X (pix)	Y (pix)	R.A. (J2000)	Decl. (J2000)	$\langle V \rangle$	$\epsilon(V)$	$n_V$	$\chi^2_V(V)$	$\langle I \rangle$	$\epsilon(I)$	$n_I$	$\chi^2_V(I)$	$\langle R \rangle$	$\epsilon(R)$	$n_R$	$\chi^2_V(R)$	$\langle B \rangle$	$\epsilon(B)$	$n_B$	$\chi^2_V(B)$
V61	3013.056	2092.584	7:38:16.42	38:52:41.3	20.128	0.114	10	1000.50	19.835	0.072	9	129.99	20.048	0.082	12	236.55	20.533	0.106	12	532.85
V62	3027.856	2051.110	7:38:16.92	38:52:39.3	20.332	0.070	11	229.27	19.690	0.052	9	94.66	19.848	0.066	11	319.95	20.590	0.096	13	457.67
V63	3030.575	2586.620	7:38:10.48	38:52:38.8	18.135	0.145	4	1266.44	16.750	0.144	3	282.33	17.379	0.056	6	100.40	19.143	0.029	6	46.90
V64	3042.306	2741.781	7:38:08.62	38:52:37.1	18.333	0.051	12	355.25	17.168	0.008	12	7.37	17.803	0.027	12	62.84	19.364	0.017	12	15.25
V65	3063.239	2200.710	7:38:15.13	38:52:34.3	20.278	0.066	9	197.23	19.863	0.028	9	13.68	19.988	0.031	11	33.88	20.433	0.024	11	22.24
V66	3069.523	2762.122	7:38:08.37	38:52:33.3	19.151	0.043	7	83.52	18.165	0.085	10	142.10	18.671	0.032	8	35.03	19.869	0.019	6	11.73
V67	3070.995	2437.930	7:38:12.27	38:52:33.1	20.357	0.070	10	223.45	19.809	0.042	9	31.06	20.073	0.053	10	101.32	20.521	0.042	10	50.83
V68	3072.023	2795.118	7:38:07.98	38:52:32.9	19.358	0.078	9	330.62	18.769	0.055	9	146.01	19.162	0.039	9	126.43	20.043	0.028	10	31.53
V69	3089.516	2754.098	7:38:08.47	38:52:30.5	19.824	0.030	9	32.00	18.796	0.020	10	18.36	19.449	0.076	6	90.33	20.391	0.180	2	49.07
V70	3093.714	2603.367	7:38:10.28	38:52:29.9	19.578	0.039	9	15.15	18.626	0.053	7	21.27	18.998	0.043	8	26.18	20.403	0.063	6	35.15
V71	3114.330	2578.930	7:38:10.58	38:52:27.0	18.140	0.090	6	230.45	16.855	0.026	5	12.54	17.488	0.051	8	182.79	19.082	0.047	14	188.94
V72	3117.580	2402.180	7:38:12.71	38:52:26.6	19.984	0.092	10	475.19	19.559	0.066	8	156.11	19.969	0.158	8	482.60	20.151	0.220	8	1121.65
V73	3119.460	2669.560	7:38:09.49	38:52:26.3	20.441	0.050	7	25.34	19.509	0.038	6	24.41	19.919	0.048	7	33.01	21.064	0.046	4	4.25
V74	3137.923	2793.357	7:38:08.00	38:52:23.7	19.907	0.081	7	241.95	19.574	0.056	4	48.08	19.800	0.094	4	60.02	19.805	0.042	5	7.68
V75	3138.039	2793.241	7:38:08.00	38:52:23.7	19.907	0.081	7	241.95	19.574	0.056	4	48.08	19.800	0.094	4	60.02	19.805	0.042	5	7.68
V76	3163.560	2694.838	7:38:09.19	38:52:20.1	20.489	0.054	7	29.29	19.561	0.054	8	35.65	20.044	0.046	3	9.91	-10.000	-1.000	0	-100.00
V77	3166.867	2777.204	7:38:08.20	38:52:19.6	20.389	0.051	10	87.63	19.401	0.016	11	10.84	19.853	0.013	7	3.95	21.028	0.026	6	5.74
V78	3179.280	2613.079	7:38:10.18	38:52:17.9	20.609	0.078	6	31.72	19.725	0.050	5	12.71	20.117	0.070	6	32.11	21.467	0.041	4	4.46
V79	3189.541	3144.189	7:38:03.79	38:52:16.4	20.868	0.054	7	80.14	20.733	0.029	7	4.99	20.912	0.044	6	30.86	20.946	0.038	8	50.39
V80	3202.909	3535.976	7:37:59.08	38:52:14.4	20.490	0.055	11	146.65	19.852	0.045	11	76.96	19.988	0.039	12	118.48	20.756	0.069	13	228.23
V81	3234.568	3285.348	7:38:02.09	38:52:10.0	20.839	0.037	4	10.11	19.976	0.050	6	39.39	20.462	0.032	7	20.70	21.496	0.028	6	2.78
V82	3249.691	2517.815	7:38:11.33	38:52:08.0	21.527	0.118	7	28.62	20.680	0.019	7	0.67	21.021	0.027	4	2.02	22.122	0.072	3	6.17
V83	3253.495	2367.815	7:38:13.13	38:52:07.5	20.422	0.062	11	154.01	19.917	0.041	9	30.66	20.167	0.050	11	84.08	20.714	0.066	12	197.09
V84	3282.462	2782.988	7:38:08.14	38:52:03.4	20.321	0.084	7	137.28	19.751	0.044	4	11.76	20.100	0.158	5	142.35	20.788	0.142	4	129.14
V85	3293.646	2360.733	7:38:13.22	38:52:01.9	21.937	0.069	7	21.18	21.178	0.024	6	0.65	21.570	0.024	5	2.00	22.528	0.025	5	1.41
V86	3312.624	2364.114	7:38:13.18	38:51:59.2	22.374	0.129	6	34.15	21.903	0.056	6	2.19	22.225	0.032	4	1.67	22.763	0.104	6	19.69
V87	3318.697	2495.017	7:38:11.61	38:51:58.3	17.257	0.057	12	966.25	15.673	0.036	11	331.44	16.369	0.057	12	892.09	18.807	0.097	14	1080.29
V88	3373.993	3053.476	7:38:04.89	38:51:50.5	18.172	0.137	6	514.47	16.937	0.144	5	570.21	17.559	0.163	5	728.20	19.244	0.095	9	323.26
V89	3374.806	2535.513	7:38:11.12	38:51:50.4	21.734	0.060	6	24.33	20.929	0.042	6	6.62	21.472	0.108	3	18.85	-10.000	-1.000	0	-100.00
V90	3405.096	3525.895	7:37:59.22	38:51:46.0	20.365	0.011	10	5.65	19.988	0.076	9	99.64	20.197	0.014	12	7.24	20.591	0.013	14	8.47
V91	3852.123	2395.051	7:38:12.85	38:50:43.4	20.397	0.068	12	221.54	19.927	0.043	10	70.54	20.138	0.049	12	93.62	20.740	0.051	13	120.08
V92	4002.830	1702.787	7:38:21.18	38:50:22.3	20.218	0.054	8	109.36	19.844	0.050	7	61.43	20.229	0.048	8	67.85	20.837	0.078	12	275.31









<sup>a</sup>The full text of this table may be found at: <ftp://taurus.tuc.noao.edu/pub/saha/Photseq/table7.txt>

Table 8. Likely Variable Stars in the Pal 4 Field

Star ID	X (pix)	Y (pix)	R.A. (J2000)	Decl. (J2000)	$\langle V \rangle$	$\epsilon(V)$	$n_V$	$\chi^2_V(V)$	$\langle I \rangle$	$\epsilon(I)$	$n_I$	$\chi^2_V(I)$	$\langle R \rangle$	$\epsilon(R)$	$n_R$	$\chi^2_V(R)$	$\langle B \rangle$	$\epsilon(B)$	$n_B$	$\chi^2_V(B)$
V1	2014.840	2730.520	11:29:24.27	28:56:50.3	22.185	0.190	5	94.03	21.503	0.142	6	25.96	21.824	0.168	5	135.51	22.697	0.200	5	152.63
V2	2609.960	1846.370	11:29:14.77	28:58:13.7	17.858	0.436	4	30355.61	15.580	0.204	4	6679.30	16.611	0.403	4	26047.48	19.840	0.383	6	17645.34
V3	2699.190	2092.920	11:29:17.41	28:58:26.3	20.258	0.075	5	228.70	19.657	0.061	6	98.17	19.823	0.076	5	374.85	20.641	0.096	6	582.93





Table 9—Continued

Star ID	X (pix)	Y (pix)	R.A. (J2000)	Decl. (J2000)	$\langle V \rangle$	$\epsilon(V)$	$n_V$	$\langle I \rangle$	$\epsilon(I)$	$n_I$	$\langle R \rangle$	$\epsilon(R)$	$n_R$	$\langle B \rangle$	$\epsilon(B)$	$n_B$
120	2515.590	3116.880	16:11:09.52	14:59:43.4	20.048	0.005	9	19.232	0.005	9	19.647	0.014	6	20.663	0.011	5
121	2544.280	2307.310	16:11:01.66	14:59:47.2	19.935	0.003	9	18.938	0.007	9	19.437	0.009	6	20.785	0.009	5
122	2636.280	1366.500	16:10:52.51	14:59:59.8	20.531	0.010	5	18.267	0.007	5	19.480	0.009	6	22.050	0.011	5
123	2674.130	1149.830	16:10:50.40	15:00:05.0	21.324	0.005	6	20.691	0.007	6	21.026	0.008	6	21.784	0.009	5
124	2748.420	1011.600	16:10:49.06	15:00:15.4	20.629	0.006	3	19.088	0.006	3	19.814	0.012	6	21.917	0.014	5
125	2801.070	904.800	16:10:48.02	15:00:22.7	16.955	0.005	3	16.148	0.005	3	16.542	0.007	5	17.676	0.006	4
126	2835.020	1086.690	16:10:49.79	15:00:27.6	19.937	0.009	3	19.115	0.007	3	19.532	0.009	4	20.546	0.006	3
127	2841.220	1163.280	16:10:50.53	15:00:28.5	20.004	0.008	3	17.675	0.003	3	18.924	0.012	4	21.503	0.014	3
128	2910.900	1202.410	16:10:50.91	15:00:38.3	19.847	0.003	3	18.650	0.005	3	19.204	0.012	4	20.843	0.012	3
129	2963.860	3148.640	16:11:09.82	15:00:46.5	20.881	0.009	8	20.002	0.011	7	20.469	0.012	4	21.601	0.013	3
130	3112.300	3124.050	16:11:09.58	15:01:07.3	19.957	0.003	7	19.147	0.008	6	19.567	0.011	4	20.585	0.007	3
131	3123.770	2515.320	16:11:03.66	15:01:08.7	16.844	0.005	3	15.675	0.009	4	16.218	0.006	3	17.850	0.010	3
132	3205.420	1200.940	16:10:50.89	15:01:19.7	17.449	0.006	3	16.428	0.005	3	16.903	0.006	4	18.370	0.010	3
133	3291.180	2101.210	16:10:59.63	15:01:32.1	19.658	0.003	4	18.643	0.008	5	19.167	0.008	4	20.536	0.007	3
134	3370.540	730.300	16:10:46.31	15:01:42.7	20.165	0.006	3	19.485	0.012	3	19.836	0.013	4	20.646	0.011	3
135	3379.470	2202.680	16:11:00.62	15:01:44.5	16.928	0.006	4	15.796	0.008	4	16.315	0.006	4	17.960	0.010	3
136	3625.850	2639.170	16:11:04.85	15:02:19.3	18.431	0.004	4	17.416	0.006	4	17.907	0.006	3	19.330	0.007	3
137	3750.900	3216.830	16:11:10.46	15:02:37.1	19.864	0.003	3	19.111	0.006	4	19.513	0.010	3	20.476	0.007	3
138	3762.800	2516.120	16:11:03.65	15:02:38.5	19.732	0.006	4	18.918	0.006	4	19.341	0.010	3	20.436	0.010	3
139	3769.830	2784.180	16:11:06.26	15:02:39.6	19.158	0.003	4	17.267	0.008	4	18.218	0.009	3	20.666	0.007	3
140	3789.520	2271.990	16:11:01.28	15:02:42.2	18.925	0.004	4	17.294	0.006	4	18.088	0.005	3	20.250	0.010	3
141	3790.740	3825.410	16:11:16.38	15:02:42.9	17.195	0.009	3	16.098	0.010	3	16.614	0.011	3	18.193	0.005	3
142	4026.450	2440.300	16:11:02.91	15:03:15.6	19.617	0.003	4	18.047	0.007	4	18.818	0.009	3	20.902	0.007	3
143	4074.200	3466.410	16:11:12.88	15:03:22.6	18.870	0.008	3	18.008	0.004	3	18.460	0.011	3	19.637	0.003	3





<sup>a</sup>The full text of this table may be found at: <ftp://taurus.tuc.noao.edu/pub/saha/Photseq/table10.txt>

Table 11. Likely Variable Stars in the Pal 14 Field

Star ID	X (pix)	Y (pix)	R.A. (J2000)	Decl. (J2000)	$\langle V \rangle$	$\epsilon(V)$	$n_V$	$\chi^2_V(V)$	$\langle I \rangle$	$\epsilon(I)$	$n_I$	$\chi^2_V(I)$	$\langle R \rangle$	$\epsilon(R)$	$n_R$	$\chi^2_V(R)$	$\langle B \rangle$	$\epsilon(B)$	$n_B$	$\chi^2_V(B)$
V1	750.640	3667.320	16:11:14.92	14:55:35.6	23.382	0.051	6	1.71	21.336	0.078	6	23.62	22.389	0.043	4	7.03	-10.000	-1.000	1	-100.00
V2	776.070	1393.620	16:10:52.83	14:55:38.3	19.869	0.233	6	4969.24	18.881	0.034	6	95.32	19.357	0.005	4	2.88	20.700	0.013	3	4.93
V3	841.310	1739.040	16:10:56.18	14:55:47.6	19.866	0.068	6	436.13	19.241	0.010	6	5.52	19.554	0.005	6	3.03	20.341	0.015	5	10.00
V4	888.810	883.730	16:10:47.87	14:55:54.0	20.943	0.100	6	240.46	20.220	0.034	6	32.05	20.560	0.005	6	0.45	21.587	0.007	5	1.17
V5	1334.060	1962.690	16:10:58.34	14:56:57.0	21.155	0.051	9	94.90	19.843	0.005	9	1.69	20.441	0.014	6	6.72	22.200	0.031	5	6.69
V6	1487.820	1992.960	16:10:58.63	14:57:18.6	21.568	0.062	6	34.19	21.307	0.024	4	1.97	21.501	0.073	6	89.40	21.849	0.047	4	29.80
V7	1867.460	2161.920	16:11:00.26	14:58:12.0	21.315	0.022	8	13.13	20.546	0.062	8	44.30	20.881	0.020	4	8.61	22.033	0.045	4	8.70

Table 12. Mean Offsets of Various Photometry of NGC 2419 vs. this paper

Study	Mean Offsets (Study - this paper) <sup>a</sup>			
	<i>B</i>	<i>V</i>	<i>R</i>	<i>I</i>
Stetson sequence	+0.014 (.021)	−0.014 (.004)	−0.017 (.005)	−0.038 (.006)
Davis sequence	+0.027 (.021)	−0.005 (.005)	+0.001 (.007)	−0.022 (.007)
Dolphin WFPC2 <sup>b</sup>	—	+0.020 (.006)	—	−0.011 (.007)
Dolphin’s independent calib of WIYN data	—	+0.002 (.004)	—	+0.003 (.007)
WIYN 0.9m data only	+0.055 (.031)	−0.002 (.012)	+0.021 (.009)	+0.032 (.018)
Sirianni et al. <i>ACS</i> <sup>c</sup>	Offsets within $\pm 0.02$ in <i>V</i> , <i>R</i> , and <i>I</i>			

<sup>a</sup>Values in parentheses indicate errors in the mean differences of the stars compared, to which the systematic error estimates from Table 2 have been added in quadrature. These may be slight underestimates, since systematic error estimates for the sequences being compared to, which should also be added in quadrature, are not generally available.

<sup>b</sup>The calibration is based on WFPC2 observations of  $\omega$  Cen, and the photometric sequence by Walker (1994).

<sup>c</sup>Estimates are from privately communicated figures. Numbers not available



<sup>a</sup>The sense of  $\Delta mag$  is "Landolt value" minus "Mean from WIYN 3.5m observations"