

# Radio, Millimeter and Optical Monitoring of GRB030329 Afterglow: Constraining the Double Jet Model

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**Abstract.** We present radio, millimeter and optical observations of the afterglow of GRB030329. *UBVR<sub>C</sub>I<sub>C</sub>* photometry is presented for a period of 3 hours to 34 days after the burst. Radio monitoring at 1280 MHz has been carried out

using the GMRT for more than a year. Simultaneous millimeter observations at 90 GHz and 230 GHz have been obtained from the Swedish-ESO Submillimeter Telescope (SEST) and the IRAM-PdB interferometer over more than a month following the burst. We use these data to constrain the double jet model proposed by Berger et al. (2003) for this afterglow. We also examine whether instead of the two jets being simultaneously present, the wider jet could result from the initially narrow jet, due to a fresh supply of energy from the central engine after the “jet break”.

**Key words.** Gamma rays: bursts – GRB030329 – afterglow – photometry – radio/mm observations

## 1. Introduction

The Gamma Ray Burst of 29th March 2003 has been an unique event. At a distance of  $\sim 870$  Mpc (assuming a cosmology of  $\Omega_\Lambda = 0.7$  and  $\Omega_m = 0.3$ , and a redshift of 0.1685 (Greiner et al. 2003a)) it is the second nearest GRB for which an afterglow has been observed. The optical and radio afterglow of this burst has been one of the brightest detected till date (Peterson & Price 2003). The spectral signature of a supernova (SN2003dh) emerged in the optical transient a few days after the burst (Stanek et al. 2003) and thus provided the first unambiguous evidence of the long suspected association between Gamma Ray Bursts and Supernovae. Multifrequency observations indicated that the GRB consisted of at least two jet-like components of ejection with different opening angles and Lorentz factors, in addition to the supernova component (Berger et al. 2003).

GRB030329 was detected and localized by the HETE-II satellite (Vanderspek et al. 2003a). The trigger H2652 occurred on 29th March 2003, at UT 11:37:14.7 and lasted more than 100 s. This was one of the brightest bursts detected by the instrument, with a 30 – 400 KeV fluence of  $1.1 \times 10^{-4}$  erg cm $^{-2}$ . The Soft X-ray Camera on board HETE-II localized the burst to be at RA (J2000) =  $10^{\text{h}}44^{\text{m}}49^{\text{s}}$  and Dec (J2000) =  $+21^\circ28'44''$  within an error circle of radius 2 arcmin. The temporal profile showed two distinct peaks in the burst, separated by  $\sim 11$  seconds. Fluence in the lower energy band,  $S_{(7-30\text{KeV})}$ , was  $5.5 \times 10^{-5}$  erg cm $^{-2}$ , which implies a hardness ratio of  $S_{(7-30\text{KeV})}/S_{(30-400\text{KeV})} > 0.33$ , classifying this GRB into the ‘X-ray rich’ category (Vanderspek et al. 2004). In a worldwide observational campaign, the afterglow was detected in all possible wavebands. The first X-ray detection by *RXTE*  $\sim 5$  hours after the burst found the source to be extremely bright with a 2 – 10 KeV flux of  $1.4 \times 10^{-10}$  erg cm $^{-2}$  sec $^{-1}$  (Marshall & Swank 2003). The optical transient had an R-band magnitude  $\sim 12$  when it was reported by Peterson & Price (2003) and Torii (2003). The VLA detected a bright 3.5 mJy radio

afterglow at 8.46 GHz (Berger et al. 2003) on 2003 March 30.06 UT. Later follow up observations in other radio frequencies were reported by Pooley et al. (2003), Rao et al. (2003a, b), Hoge et al. (2003) and Kuno et al. (2003). Around 7 days after the burst the optical spectrum showed the signature of an underlying supernova emission (Stanek et al. 2003; Hjorth et al. 2003; Matheson et al. 2003), and the presence of the associated supernova SN2003dh was confirmed later by spectroscopic measurements. Continued monitoring has provided an unprecedentedly rich temporal coverage of the transient in all wavebands (Lipkin et al. 2004; Tiengo et al. 2003; Sheth et al. 2003; Berger et al. 2003; Guziy et al. 2005; Gorosabel et al. 2005a).

We present in this paper the observations done by an Indo-European GRB collaboration at radio, millimeter and optical wavelengths over a total observing span of nearly one year. The millimeter observations were conducted for more than a month at ESO and IRAM. Radio observations at a frequency of 1280 MHz were carried out with the Giant Meter Wave Radio Telescope (GMRT) (<http://www.ncra.tifr.res.in>) operated by the National Center for Radio Astrophysics, Pune. Optical follow-up was conducted till 34 days after the burst using the 2.01m Himalayan Chandra Telescope (HCT) of the Indian Astronomical Observatory (IAO), Hanle and the 1.04m Sampurnanand Telescope (ST) at the State Observatory, Naini Tal, now renamed as Aryabhata Research Institute of Observational Sciences (ARIES). The GMRT observations represent the lowest frequency detection and the longest follow-up of the afterglow reported so far. This is also the first detection of a GRB afterglow by the GMRT. Our optical observations in  $UBVR_CI_C$  pass-bands started  $\sim 3$  hours after the burst. Except in the  $R$  band, these represent the earliest photometry of the optical transient. Our optical data also fill many temporal gaps existing in the literature (e.g. Lipkin et al. 2004). At millimeter waves, we used the SEST to make an early detection ( $\sim 0.6$  days) at 86 GHz, albeit at low statistical significance, and subsequently monitored the burst until June 19, 2003, using the IRAM Plateau de Bure interferometer, making several simultaneous detections at frequencies ranging from 86 to 240 GHz. The early evolution of the optical afterglow followed the familiar behavior of a power-law decay in the light curve, with a steepening (“jet break”) around 0.5 days (Garnavich et al. 2003; Smith 2003; Price et al. 2003). The nature of the optical light curve, however, deviated from this simple model after  $\sim 1.5$  days, displaying substantial variability and some change in average slope (see Lipkin et al. 2004 for a full compilation). At radio frequencies, the evolution of the afterglow flux was much slower than the early optical decay, and a steepening of the light curve was observed at  $\sim 10$  days after the burst (Berger et al. 2003). In order to explain this behavior, Berger et al. (2003) introduced a double-jet model, with a narrow jet responsible for the early optical emission and a wider jet contributing to the radio and late-time optical and X-ray emission. In addition, the optical light curve also has a significant contribution from the underlying supernova SN2003dh after about a week following the burst. The sharp bump

in the optical light curve seen at  $\sim 1.5$  days has been attributed by Berger et al. (2003) to the deceleration epoch of the wider jet.

Granot, Nakar & Piran (2003) have instead proposed that the bump in the R-band lightcurve at 1.5 days and at three successive epochs could be attributed to refreshed shocks (see also Guziy et al. 2005). In their original model, however, the second jet break at  $\sim 10$  days was not expected (Piran, Nakar & Granot 2003).

In this paper, we examine the ability of the double jet model of Berger et al. to fit our observations along with other multi-band observations reported in the literature. We present a refined set of parameters for the two jets resulting from our fits. We also examine whether the two jets are in fact distinct entities or whether a refreshed shock might have converted the decelerating narrow jet to a wider, more energetic jet. In sect. 2, 3 and 4 we describe respectively our radio, millimeter and optical observations, in sect. 5 we present theoretical model fits and in sect. 6 we discuss their implications. In sect. 7 we calculate the reverse shock emission expected from the model. In sect. 8 we estimate the possible contribution from SN2003dh. Sec. 9 summarizes our results.

## 2. Radio Observations

We obtained radio observations at the center frequency of 1280 MHz using the Giant Meter Wave Radio Telescope (GMRT) located at Khodad, near Pune in Western India. GMRT is operated by the National Center for Radio Astrophysics. The telescope is an interferometric array of 30 fully steerable parabolic dishes of 45m diameter each, spread over a distance of up to 25km. The telescope operates at several spot frequencies between 150 to 1500 MHz, with a maximum bandwidth of 32 MHz at any given band. All feeds provide dual polarization outputs. The highest angular resolution achievable ranges from about 20 arc-sec at the lowest frequencies to about 2 arc-sec at 1.5 GHz. More details about the GMRT can be found in Swarup et al. (1991).

We made the first detection of the GRB030329 radio transient in 1280 MHz band on 31st March 2003, 2.3 days after the GRB trigger (Rao et al. 2003a). The source was bright, with a flux of 0.33 mJy,  $8\sigma$  above the background rms noise. Since then we have carried out a series of monitoring observations at the center frequency of 1280 MHz, at nine epochs until April 18, 2004. We used a bandwidth of 32 MHz during these observations. The flux scale is set by observing the primary calibrator 3C286, 3C147 or 3C48. A phase calibrator was observed before and after a 30 to 45 min scan on GRB030329. The integration time was 16 s. The data recorded from GMRT have been converted to FITS and analyzed using Astronomical Image Processing System (AIPS). At each epoch, the stability of the flux scale was checked by measuring the fluxes of a few background sources. The flux of background sources were consistently stable at all epochs, and the rms fluctuation in their flux (estimated to be 8%) has been taken as the error in the

flux calibration. The final error quoted is the quadrature sum of the measurement error and the flux calibration uncertainty. The rms noise in the image is estimated from a sufficiently large area where no source was visible. The rms noise ranged from 35 to 80  $\mu\text{Jy}$  on different epochs. The radio transient was detected in all our observations. The flux showed an initial rise, reaching a peak of 2.5 mJy at  $\sim 133$  days after the burst followed by a secular decay (Table 1). As mentioned in sect. 5, the observed nature of the 1280-MHz light curve tightly constrains the non-relativistic transition of the expanding jet.

### 3. The millimeter face of the afterglow

The first, single dish, observations were carried out on March 29–30, 2003 at La Silla, Chile, with the 15m Swedish-ESO Submillimeter Telescope (SEST, Booth et al. 1989). We used the dual channel IRAM SIS-receiver, which allows simultaneous observations at 1.3 and 3 mm wavelengths. The 1.3 mm receiver was tuned to 215.0 GHz and the 3 mm receiver was tuned to 86.243 GHz. The backend used was a 3-level correlator with 2000 channels and the bandwidth was 1028 MHz for the 215 GHz receiver, and 512 MHz for the 86 GHz receiver. We used a dual beam switch mode (12 arcmin throw in azimuth) and integrated for 60 seconds (on source) per measurement. Observations were performed over three nights (see Table 3). During the observations the sky was clear and stable. The pointing was checked before and after the observations each night and was found to be stable within  $\pm 3''$ .

The intensity scale was calibrated using the standard chopper wheel method, with an internal calibration error of  $\sim 10\text{-}20\%$ . The intensity scale is converted from K to Jy using 40 Jy/K and 25 Jy/K for the 86 GHz and 215 GHz measurements respectively. Two blank sky integrations were done in order to test the performance of the receiver system (see Table 3). These observations were done during day time, resulting in higher noise levels ( $16 \pm 164$  and  $4 \pm 492$  mJy respectively) than the observations of GRB030329 which were done at night time. The blank sky observations did not produce reliable results at 215 GHz.

In order to derive the continuum level, each individual scan was inspected for spikes (which were removed) and abnormal baseline curvature. The average continuum and an estimate of the noise rms were derived for the central 400 MHz of the 86 GHz spectra, and 700 MHz of the 215 GHz spectra. The subscans were subsequently added together, weighted with their individual noise rms. Scans deviating by more than 3 sigma from this initial analysis were removed and the procedure repeated with the remaining scans. This generally does not change the average continuum flux, but decreases the dispersion. The final average continuum fluxes are given in Table 4. The dispersion in the continuum fluxes among individual scans are given as the  $1\sigma$  error in Table 4.

The dispersion in the measured continuum level, using the present technique, reflects the variable sky background and the inherent calibration uncertainty. It does not follow a normal distribution and increasing the number of subscans does not necessarily decrease the dispersion. The dispersion of continuum fluxes follows more a top-hat distribution. As such, it is not a good measure of the precision of the final continuum flux value, but nevertheless the only measure of the uncertainty at our disposal.

Under clear and stable weather conditions, as was the case for the SEST observations of GRB030329, the average flux value is usually a robust estimate if the number of subscans is sufficiently large (e.g.  $>20$ ). The flux value derived for the night of March 29 is higher than any of the following measurements, while the noise rms remains more or less constant. This is a strong indication that continuum flux was detected during the first night but not during the following nights. The weather conditions at the higher frequency band (215 GHz) were worse than at 86 GHz. We therefore conclude that at least at 3 mm (86 GHz), we did tentatively detect continuum flux during March 29, at a level around 80 mJy. The uncertainty associated with this, however, remains poorly defined and the quoted flux value should be regarded as approximate. At a  $2\sigma$  level, the quoted uncertainty gives an upper bound of  $\sim 104$  mJy on the flux of the source.

At Plateau de Bure interferometer (Guilloteau et al. 1992), we carried out simultaneous observations in two bands, around 90 GHz and 230 GHz, the actual band center frequencies being slightly different on different days. The instrument was operating on compact D configuration with six antennas (6Dp) until May 3rd and with fewer antennas later (see Table 4). Data reduction and analysis was performed with the GLIDAS (Grenoble Image and Line Data Analysis) software. The target was observed at coordinates RA (2000) =  $10^h44^m50.030^s$ , Dec (2000) =  $+21^\circ31'18.15''$  and the best position offsets were found on 31 Mar 2003 to be:  $-0.87'' \pm 0.03''$  and  $-0.78'' \pm 0.02''$  respectively. Therefore, the absolute coordinates are RA (2000) =  $10^h44^m49.968^s(\pm 0.002s)$ , Dec (2000) =  $+21^\circ31'17''.37(\pm 0.02'')$ . Thus, within the error limits, the measured position definitely stays constant throughout the observing period. The source has been detected up to May 3rd, after which there are only upper limits available.

#### 4. Optical Observations and Data Reduction

Starting about 3 hours after the burst, we obtained a total of 13, 70, 87, 167, and 39 photometric observations in Johnson  $U$ ,  $B$ ,  $V$  and Cousins  $R$  and  $I$  bands respectively, from both ST and HCT. The CCD used at HCT was  $1024 \times 1024$  pixel<sup>2</sup> with the entire chip covering a field of  $\sim 4'.7 \times 4'.7$  on the sky. It has a read out noise of  $11 e^-$  and gain of  $4.8 e^-/ADU$ . A CCD chip of size  $2048 \times 2048$  pixel<sup>2</sup> was used at ST, which covers a field of  $\sim 13' \times 13'$ . The gain and read out noise are  $10 e^-/ADU$  and  $5.3 e^-$  respectively. The frames were binned in  $2 \times 2$  pixel<sup>2</sup> to improve the signal-to-noise ratio of the source.

Several twilight flat field and bias frames were obtained for the CCD images at both the telescopes. While imaging the optical transient (OT), several short exposures upto a maximum of 15 min were taken in various filters. We used MIDAS, IRAF and DAOPHOT softwares to process the CCD frames in the standard fashion. The bias subtracted and flat fielded CCD frames were co-added, whenever found necessary.

The *BVRI* magnitudes of the OT obtained from the Sampurnanand Telescope at Naini Tal were calibrated differentially using secondary standard stars no. 1, 11, 14, 19, 37 and 57 in the list of Henden (2003), while the *U* magnitudes were determined using reference stars 1, 11 and 37 from Henden (2003). *BVR* magnitudes obtained at the IAO were also calibrated differentially, using reference stars no. 14, 18 and 19 of Henden (2003). OT magnitudes at similar epochs, determined using photometry from the two sites, are consistent with each other within the limits of uncertainty. A full compilation of our data is presented in Table 2.

Figure 1(a) displays our *UBVRI* photometry (open circles) along with some of the data published by other authors (filled circles). As mentioned above, in *UBVI* bands we report the earliest observations and in all the bands our observations fill in gaps in photometric data reported in the literature so far, at epochs beyond one day after the burst.

Reported X-ray observations of this afterglow are shown in an accompanying figure 1(b). The role of the X-ray observations in constraining the model, along with observations in other bands, is discussed in sect. 5.

## 5. Modeling the Multiband Observations

It was proposed by Rhoads (1999) that the explosion which makes the GRB may not be isotropic. Ejection of matter in the explosion could in fact be collimated within narrow cones. If so, then a signature of this collimation is expected in the form of an achromatic steepening (“jet break”) in the afterglow light curve (Rhoads 1999, Sari, Piran & Halpern 1999). The jet break appears when the Lorentz factor of the expanding outflow drops below the inverse of the initial angle of collimation, causing the lateral expansion of the jet to dominate over its radial motion. This behavior has now been observed in a number of afterglows (see reviews by Piran 2004 and Sagar 2002).

### 5.1. The double jet model for GRB030329

The X-ray and optical lightcurves of GRB030329 afterglow had an initial temporal slope of  $\sim -0.9$ . Around half a day, the optical decay steepened to an index of  $\sim -1.9$ . Sampling of the X-ray evolution was poor, but an interpolation of the *RXTE* and *XMM* data obtained during  $\sim 0.1$  to  $\sim 100$  days indicated a break almost at the same time as the optical steepening mentioned above (Tiengo et al. 2003). A nearly simultaneous

break observed in frequencies separated by four orders of magnitude suggested that the break observed at  $\sim 0.5$  days was a jet break.

At radio frequencies, the first observations obtained were around 1 days, somewhat later than the epoch of the jet break obtained from optical and X-ray observations. The radio and millimeter light curves, however, did not display the behavior expected after a jet break at 0.5 days. Instead, they were rather well described by a second jet, with a jet break around 10 days (Berger et al. 2003; Sheth et al. 2003). The optical light curve showed a re-brightening around 1.5 days, followed by a slower decay consistent with that expected from the second jet to which the radio emission was attributed (Berger et al. 2003; Lipkin et al. 2004). This led Berger et al. (2003) to propose a double-jet model for GRB030329. The optical re-brightening at 1.5 days was attributed to the epoch of deceleration of the second jet. Berger et al. (2003) estimated the initial opening angle of the first (narrow) jet to be  $\sim 5^\circ$  and that of the second (wide) jet to be  $\sim 18.4^\circ$ . The energy contents of the two jets were estimated to be  $\sim 6.7 \times 10^{48}$  erg for the narrow jet and  $\sim 10^{49}$  erg for the wide jet. Compared to the narrow jet, a later deceleration epoch implied a smaller initial Lorentz factor for the wide jet, and correspondingly a much higher initial baryon load. The double jet model has been found to be consistent with most observations reported till date. The optical emission at late times ( $t > 10$  days) is, however, dominated by the associated supernova SN2003dh.

We attempted modeling our observations along with multiwavelength data available in the literature within the ambit of this double-jet model (model 1). The basic quantity in our model is the synchrotron source function, which we consider to have appropriate power-law forms between the usual break frequencies. Transition from one power-law phase to another is made gradual through a Band type smoothening (Band et. al. 1993) at the peak frequency  $\nu_m$  and the cooling frequency  $\nu_c$ . The self-absorption frequency  $\nu_a$  is not treated as a break; instead the absorption is incorporated into the expression for synchrotron optical depth, which, along with the source function, yields the flux at any given frequency. In the co-moving frame of the shock, the optical depth is set to unity at  $\nu = \nu_a$  (comoving). We incorporate transition to a non-relativistic phase of expansion, as in Frail, Waxman & Kulkarni (2000). The non-relativistic transition is treated as a sharp break at a time  $t = t_{\text{nr}}$ . We obtain our fits through the usual  $\chi^2$  minimization procedure, using  $\nu_a, \nu_m, \nu_c, t_{\text{nr}}$ , the electron distribution index  $p$  and the jet break time  $t_j$  as the fit parameters. Here we are trying to model the underlying smooth power law behavior rather than the short time scale variabilities in the lightcurve. Since our model does not include the short-term variabilities, the nominal  $\chi^2$  obtained is relatively high. The best fit from this model is shown in figures 1 and 2. We derived the physical parameters  $E_{\text{iso}}$  (the isotropic equivalent energy of the burst),  $n$  (number density of the ambient medium),  $\epsilon_e$  and  $\epsilon_B$  (fractional energy content in the electrons and in the magnetic field respectively) using the expressions in Wijers & Galama (1999) from these fitted

parameters, with appropriate modifications to place  $\nu_a$  at  $\tau_\nu = 1$  instead of 0.35 used by Wijers & Galama (1999).

### 5.1.1. The wide jet

The parameters of the wide jet are well constrained by data in the 4–250 GHz range, as discussed by Berger et al. (2003) and Sheth et al. (2003). At the jet break time of 9.8 days, we find the self absorption frequency  $\nu_a$  to be  $1.3_{-0.06}^{+0.25} \times 10^{10}$  Hz, the synchrotron peak frequency  $\nu_m$  to be  $3.98_{-0.1}^{+0.5} \times 10^{10}$  Hz with a peak flux  $F_m$  of  $44.7_{-2.0}^{+1.0}$  mJy. The post jet-break decay in radio, the peak optical flux at 1.5 days and the late X-ray observations at 37 and 61 days together constrain the electron energy distribution index  $p$  to  $2.3_{-0.02}^{+0.05}$  and the cooling frequency  $\nu_c$  to  $3.98_{-2.0}^{+1.3} \times 10^{14}$  Hz after the jet break. These parameter values are very similar to those derived by Berger et al. (2003). Our 1280 MHz observations show a gradual rise of flux to a peak around 133 days. This behavior can be reproduced by a non-relativistic transition of the jet at  $t_{nr} = 42_{-7}^{+17}$  days.

### 5.1.2. The narrow jet

The jet break time for the narrow jet is derived to be  $0.69_{-0.06}^{+0.08}$  days from the optical light curve. Using a galactic extinction  $E(B - V) = 0.025$  mag in the direction of the GRB (Schlegel et al. 1998), the early optical and X-ray observations, both before and after this break, are well described by an electron energy distribution index  $p$  of  $2.12 \pm 0.05$ . Some authors (e.g. Sato et al. 2003, Lipkin et al. 2004) have conjectured that there occurs a passage of the cooling break ( $\nu_c$ ) through optical bands within the first few hours after the GRB, based on a derived change in slope of R-band light curve around  $\sim 0.25$  days, and a small color evolution. Our own observations have a continuous coverage from  $\sim 0.15$  days to  $\sim 0.3$  days after the burst, taken from the same instrument and calibrated uniformly on the same scale. We find that our data over this interval are fit very well by a single power law. Even the data reported by Lipkin et al. (2004) and Sato et al. (2003) do not conclusively demonstrate a secular steepening at  $\sim 0.25$  days; the effect could be easily mimicked by short-term variability riding on a single, underlying power law. Colors derived from our own multiband observations have somewhat large errors ( $\sim 0.08$  mag  $1-\sigma$ ), and we are unable to discern the  $\sim 0.1$  mag systematic change in  $B - R$  color reported by Lipkin et al. (2004). As mentioned by Lipkin et al. (2004), there could be various reasons for this early color evolution. We do not feel that the passage of cooling break through optical bands at early times can be conclusively established from the existing observations. From multiband fits, however, we estimate  $\nu_c$  to be at  $1.0_{-0.5}^{+1.0} \times 10^{16}$  Hz at 0.5 days. At times  $> 1.5$  days, the contribution from the wide jet is sufficient to reproduce the radio light curves (see fig. 2), so the radio emission from the narrow jet is constrained to be almost negligible. In order to achieve this, we need

to have the peak frequency  $\nu_m$  and the self absorption frequency  $\nu_a$  of the narrow jet to be as high as possible. The passage of  $\nu_m$  is not observed through the optical band, so we chose it to be just below the  $R$ -band at the earliest epoch ( $\sim 0.05$  days) at which data are available. This results in  $\nu_m$  regressing to  $\sim 10^{13}$  Hz at the jet break epoch of 0.5 days. The fitted peak flux  $F_{(\nu=\nu_m)}$  is  $19.8_{-2.4}^{+9.3}$  mJy at this time. The density of the ambient medium can be derived from the rather well-constrained parameters of the wide jet, and works out to be  $n \sim 8$  atom/cc. This, along with the narrow jet parameters mentioned above, predicts the self absorption frequency of the narrow jet to be  $\nu_a$  equal to  $3.1_{-0.63}^{+0.14} \times 10^9$  Hz at 0.5 days. We find that this value of  $\nu_a$  yields adequate suppression of the narrow jet flux at low frequencies for the model to be consistent with our 1280 MHz data.

As mentioned earlier, the  $\chi_{\text{DOF}}^2$  is somewhat high due to short-term variabilities in the observed light curve. From the fit we excluded the first seven days of data at 4.86 and 8.46 GHz, which appear to have been affected by scintillations. The first data point at 250 GHz ( $\sim 1.5$  days) and at 100 GHz (0.8 days) were also removed from the fit. We did not consider five more data points in radio bands (of 2.6 days in 43 GHz, days 1 and 12 in 22 GHz and days 3.5 and 4.7 in 15 GHz) which produced high  $\chi^2$  values due to scatter. We exclude the optical data from discussion for epochs larger than  $\sim 5$  days because of the dominant contribution from SN2003dh. A  $\chi_{\text{DOF}}^2$  of 23.3 is obtained for the best fit with this model. The optical (mostly  $V$  and  $B$ ) bands dominate the contribution to  $\chi^2$  along with the lower radio frequencies (4 GHz, 8 GHz and 15 GHz). The number density of the ambient medium is inferred to be  $8.6_{-5}^{+12}$ . We infer the fractional energy content in relativistic electrons and magnetic field to be  $0.56_{-0.5}^{+0.4}$  and  $4_{-1.8}^{+1.9} \times 10^{-4}$  respectively for the narrow jet, and  $9.0_{-1}^{+3} \times 10^{-2}$  and  $11.9_{-7}^{+10} \times 10^{-4}$  for the wide jet. We derive  $1.4_{-0.8}^{+1.3} \times 10^{51}$  erg for the isotropic equivalent energy and  $6.2_{-0.03}^{+0.02}$  degrees for the opening angle of narrow jet. This corresponds to a total energy content of  $3.3_{-2.4}^{+4.8} \times 10^{48}$  erg in the jet. For the wide jet, we derive an isotropic equivalent energy of  $1.2_{-0.2}^{+0.4} \times 10^{50}$  erg, opening angle of  $23.3_{-0.04}^{+0.07}$  degrees, and a total energy of  $5.0_{-2.1}^{+3.3} \times 10^{48}$  erg.

We calculated the rising flux from the wide jet at  $t < 1.5$  days, assuming time evolution of the spectral parameters to be of the form  $\nu_m \propto t^0$ ,  $\nu_c \propto t^{-2}$ ,  $f_{\nu_m} \propto t^3$  (Peng, Königl & Granot 2004) and  $\nu_a \propto t^1$  and normalizing them at 1.5 days.

We also explored the possibility of the ambient medium of the burst being generated by a stellar wind, with a density profile of  $n(r) \propto r^{-2}$ , but were unable to obtain consistent fits with the double-jet model. If the model is tuned to reproduce radio data in the 8–43 GHz frequency range, it leads to an overprediction of fluxes in millimeter bands and an underprediction at 1280 MHz (cf. Fig. 2).

## 5.2. Refreshed Jet ?

Most of the observations are well reproduced by a model which sums the contributions from the wide and narrow jet. We note that the contribution of the narrow jet is almost negligible at radio bands; in fact the wide jet alone is quite sufficient to account for the observed flux after  $\sim 1.5$  days. It therefore appears to us that the data could be well described if, instead of both jets contributing simultaneously to the emission, the narrow jet alone contributes at epochs earlier than  $\sim 1.5$  days, and only the wide jet after that time. This suggests that a possible re-energization event around or before  $\sim 1.5$  days could have refreshed the initially narrow jet, which had entered the lateral expansion phase, and given it additional forward momentum, converting it into the second, ‘wide’ jet. The opening angle of the laterally expanding, initially narrow jet around  $\sim 1.5$  days is estimated to be  $\sim 20^\circ$ , not far from the initial opening angle inferred for the wide jet ( $\sim 23^\circ$ ). We therefore consider it possible that a re-energization event occurred after the jet break of the narrow jet as suggested by Granot, Nakar & Piran (2003) and the double-jet model for GRB030329 could represent the conversion of an initially narrow jet to a wide one by the re-energization.

In a simple representation of such a re-energization, we assume that the physical parameters of the fireball, namely  $E$ ,  $\epsilon_e$  and  $\epsilon_B$  undergo a change after re-energization, while the ambient density  $n$  remains the same. We allow these physical parameters to be determined by the model fits.

Fits to multiwavelength observations with this model (model 2) are shown in figs. 3 and 4. We excluded the same set of data in these fits as done in model 1. The minimum  $\chi^2_{\text{DOF}}$  we obtained with this model is 24.5, slightly higher than the value for the previous model. Here again, the  $\chi^2$  is dominated by the optical band as well as the low radio frequencies. Both model 1 and model 2 underpredict the flux at 1280 MHz peak by a factor of 2. But it is difficult to make a distinction between model 1 and model 2 from the small difference in the  $\chi^2$ . Parameters for the initial, narrow jet are nearly the same as those listed in sect. 5.1.2. except the self absorption frequency  $\nu_a$  ( $2.9^{+0.8}_{-0.06} \times 10^9$  Hz), which is implied by the value of  $n$  inferred from the parameters of the jet after re-energization. Parameters of the refreshed jet are only marginally different from that of the wide jet discussed in sect. 5.1.1. We obtain a  $p$  of  $2.24 \pm 0.02$ ,  $t_j$  of  $10^{+2.3}_{-1.0}$  days, and at 9.8 days, a cooling frequency of  $5.0^{+2.1}_{-1.5} \times 10^{14}$  Hz, and a marginally reduced self absorption frequency of  $1.1^{+0.3}_{-0.05} \times 10^{10}$  Hz. The non-relativistic transition time  $t_{\text{nr}}$  is  $63^{+13.5}_{-30}$  days. The number density of the ambient medium is inferred to be  $6.7^{+13.5}_{-3.0}$ . We infer the fractional energy content in relativistic electrons ( $\epsilon_e$ ) and magnetic field ( $\epsilon_B$ ) to be  $0.53^{+0.5}_{-0.4}$  and  $4^{+2.5}_{-1.3} \times 10^{-4}$  respectively for the initial (narrow) jet, and  $0.103^{+0.05}_{-0.01}$  and  $1.0^{+1}_{-0.5} \times 10^{-3}$  for the refreshed (wide) jet. We derive an isotropic equivalent energy of  $1.2^{+1.6}_{-0.6} \times 10^{51}$  erg for the original jet, which has an initial opening angle of  $5.9^{+0.3}_{-0.1}$

degrees. This corresponds to a total energy content of  $3.2_{-2.8}^{+6} \times 10^{48}$  ergs in the jet. After re-energization, which is assumed to end around  $\sim 1.5$  days, the total energy content of the jet increases to  $5.8_{-1.7}^{+5.9} \times 10^{48}$  erg, and the jet widens to  $20.5_{-0.03}^{+0.1}$  degrees.

The transition to the refreshed physical parameters of the jet is expected to be gradual, over the time required to establish a new equilibrium. The timescale for achieving a new Blandford-McKee structure will be roughly equal to the time the second wave requires to cross the existing shocked shell. This time, in the co-moving frame, can be estimated as the thickness of the matter in the co-moving frame of the shock, divided by  $c$ , which in the observer's frame will be  $\Delta t \approx R/(a\Gamma_{\text{old}}\Gamma_{\text{new}}c)$ , where  $a \sim 5 - 10$ . At 1.5 days, when the new power law phase begins, we calculate the bulk Lorentz factor ( $\Gamma_{\text{new}}$ ) to be  $\sim 2.3$ , by extrapolating the value of  $\Gamma$  from the jet break time. Extrapolation for the initial jet produces  $\Gamma_{\text{old}}$  to be close to this value. Hence  $\Delta t$  will be of the order of  $0.2 - 0.3$  days. From a close examination of the optical lightcurve, we find that the refreshment episode could begin at  $\sim 1.2$  days and at  $\sim 1.5$  days, the new power law phase begins.

## 6. Discussion

We have seen above a comparison between the parameters of the two models. We now discuss their compatibility with the constraints imposed by other available observations of the afterglow. The angular size of the fireball estimated by Taylor et al. (2004) can be reproduced by models with isotropic equivalent energy to external density ratio in the range  $10^{50} - 10^{52}$  erg  $\text{cm}^{-3}$  (Oren, Nakar & Piran 2004). The parameters extracted from both the models fall close to this range.

Polarization measurements of the afterglow are available in optical (Greiner et al. 2003b) and in radio (Taylor et al. 2004) bands. In the optical, the degree of polarization decreases shortly after the jet break at  $\sim 0.5$  days and rapid variations in polarization start occurring around 1.5 days, which according to Greiner et al. could be the beginning of a new power law phase. The linear polarization at 8 GHz ( $< 0.1\%$ ) around 8 days is significantly lower than the optical polarization ( $\sim 2\%$ ), which could be due to the fireball being optically thick at this frequency. This polarization behavior has been thought to support the double jet model (Greiner et al. 2003b; Taylor et al. 2004); however this is equally applicable to the refreshed jet.

We point out that the first millimeter-wave observation at 86 GHz by the SEST at 0.6 days and the 100 GHz observation reported by Sheth et al. (2003) at 0.8 days are not well fit by this simple model. Nor does the double-jet model of Berger et al. (2003) succeed in reproducing this well.

## 7. Reverse Shock

It has been pointed out (Piran, Nakar & Granot 2003) that the deceleration of the wide jet at 1.5 days is expected to be accompanied by a strong radio flash from the reverse shock, which is not observed. This concern remains for the refreshed jet model, too. The two models differ in the nature of the medium in which the second shock front decelerates (in model 1, a normal ISM while in model 2, it is the material already shocked by the first shell). The Sedov length in these two cases are also to be evaluated differently, since in model 1, the shell responsible for the wide jet encounters matter all the way from the progenitor star while in model 2, the second wave of energy passes through the region evacuated by the first jet.

We estimated the reverse shock emission expected in either model, assuming that the shock is ultrarelativistic (thick shell case) (Sari & Piran 1995; Kobayashi 2000). For model 1 there are four relevant regions, namely, normal ISM (0), ISM shocked by the wide jet (1), reverse shocked ejecta (2) and the cold second shell (3). Region 1 and region 2 are separated by a contact discontinuity (CD). We followed the formulation of Kobayashi (2000) to obtain the flux expected from reverse shock. In model 2, since the reverse shock originates when the second wave of energy decelerates into the already shocked material, the space ahead of the CD will be divided into three regions instead of two (Kumar & Piran 2000). The five relevant regions in this case are: the ISM (0), the ISM shocked by the first wave (1), ISM additionally shocked by the second wave (2), reverse shocked ejecta (3) and the cold second shell (4). A contact discontinuity separates regions 2 and 3. Assuming pressure balance at the CD, and for the ejecta using the assumption that  $n_4 R^2$  is constant, (where  $n_4$  is the number density of the cold shell and  $R$  is the distance to the CD), one obtains the bulk Lorentz factor ( $\gamma_3$ ) and thermodynamic quantities (density and pressure) of region 3. These quantities allow one to estimate the synchrotron emission from that region. The thickness ( $\Delta$ ) of the ejecta is an unknown parameter in both models. For a given  $\Delta$  the peak flux produced by model 1 is two orders of magnitude lower than model 2 at the deceleration time of 1.5 days. The computed flux is inversely related to  $\Delta$ , and the minimum value  $\Delta$  can reach without overpredicting the flux observed at various bands is  $\sim 10^{10}$  cm for model 1, and  $\sim 10^{13}$  cm for model 2. More detailed investigations taking into account detailed hydrodynamics in non-spherical geometry as well as the density structure of a post-jet break fireball may be necessary to get a better estimate.

## 8. SN 2003dh

The optical emission observed at times later than  $\sim 7$  days cannot be fully accounted for by the afterglow models discussed so far. We attribute the excess emission at these late epochs primarily to the associated supernova, SN2003dh. We subtract the afterglow flux

predicted by the model and the flux due to the host galaxy (Gorosabel et al. 2005b) from the observed data to estimate the contribution from the supernova. Fig. 5 displays the flux attributed to the supernova in the two models discussed above. The K-corrected light curve of SN1998bw (Galama et al. 1998) appropriate for the redshift of GRB030329 is shown along with the residuals for comparison. While being similar in temporal behavior, the residuals from model 1 are fainter by  $\sim 0.3$  mag in comparison to an equivalent SN1998bw lightcurve. In case of model 2, this difference is  $\sim 0.4$  mag. These results compare well with the estimate of Lipkin et al. who find that an SN1998bw lightcurve, diminished by 0.3 magnitude, is required to fit the observed data.

## 9. Summary

We have presented low frequency radio, millimeter wave and optical observations of the afterglow of GRB030329, and interpreted them in terms of the double-jet model discussed earlier in the literature. Our main conclusions are summarized below.

- The 1280 MHz GMRT observations, starting  $\sim 2$  days after the burst and continuing for over a year, provide constraints on the self absorption frequency of the emission region as well as the epoch of non-relativistic transition of the fireball.
- We report several simultaneous two-band ( $\sim 90$  and  $\sim 250$  GHz) detections of the afterglow for over a month after the burst.
- Our optical observations span a temporal range of 3 hours to 34 days and fill in many gaps in the coverage reported so far in the literature. In *UBVI* bands our data represent the earliest photometry of the optical transient.
- We find that the data can be well fit by a model where the two jets are present either simultaneously or in exclusion of each other. We derive a non-relativistic transition time of  $\sim 42$  days for model 1 and  $\sim 63$  days for model 2. Although concerns with reverse shock emission remain in both models, we deem it possible that the optical re-brightening seen at the epoch of  $\sim 1.3$  days could be a re-energization of the jet, which resulted in the initially narrow jet being converted into a more energetic wider jet.

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**Table 1.** 1280 MHz radio Observations at GMRT

	Mid UT	Flux (mJy)
2003		
	Mar 31 18:30	0.33 +/- 0.09
	Apr 1 20:30	0.33 +/- 0.09
	May 30 16:30	1.10 +/- 0.15
	Jun 28 11:00	1.40 +/- 0.16
	Aug 9 07:30	2.50 +/- 0.25
	Oct 20 05:00	1.50 +/- 0.17
	Dec 3 00:20	1.30 +/- 0.14
2004		
	Jan 29 22:12	1.20 +/- 0.18
	Apr 18 14:32	1.10 +/- 0.12

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**Table 2.** Optical Transient

Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope	
<b><i>U</i>– passband</b>					
2003	March				
	29.6459	0.1619	13.606 ± 0.051	200	ST
	29.6918	0.2078	13.824 ± 0.052	200	ST
	29.7042	0.2202	13.831 ± 0.060	200	ST
	29.7186	0.2346	13.910 ± 0.061	200	ST
	29.7345	0.2505	14.016 ± 0.053	200	ST
	29.7477	0.2637	14.069 ± 0.051	200	ST
	30.7558	1.2718	16.454 ± 0.053	300	ST
	30.7738	1.2898	16.486 ± 0.057	300	ST
	30.7909	1.3069	16.483 ± 0.056	300	ST
	30.8084	1.3244	16.507 ± 0.063	300	ST
2003	April				
	1.6524	3.1684	17.156 ± 0.061	400	ST
	1.7046	3.2206	17.153 ± 0.058	400	ST
	1.7626	3.2786	17.201 ± 0.055	400	ST
<b><i>B</i>– passband</b>					
2003	March				
	29.6429	0.1589	14.189 ± 0.026	150	ST
	29.6483	0.1643	14.1463 ± 0.070	900	HCT
	29.6552	0.1712	14.248 ± 0.032	100	ST
	29.6603	0.1763	14.1933 ± 0.046	900	HCT
	29.6891	0.2051	14.409 ± 0.031	100	ST
	29.7013	0.2173	14.440 ± 0.022	100	ST
	29.7153	0.2313	14.534 ± 0.012	200	ST
	29.7312	0.2472	14.599 ± 0.016	200	ST
	29.7445	0.2605	14.653 ± 0.014	200	ST
	29.7574	0.2734	14.704 ± 0.028	200	ST
	30.7378	1.2538	17.0433 ± 0.033	900	HCT
	30.7598	1.2758	17.137 ± 0.015	200	ST
	30.7730	1.289	17.0573 ± 0.060	900	HCT
	30.7776	1.2936	17.170 ± 0.016	200	ST
	30.7948	1.3108	17.167 ± 0.014	200	ST
	30.8123	1.3283	17.192 ± 0.016	200	ST
	30.8359	1.3519	17.0923 ± 0.032	900	HCT
	31.5978	2.1138	17.464 ± 0.032	300	ST
	31.6707	2.1867	17.3683 ± 0.053	900	HCT
	31.6763	2.1923	17.421 ± 0.032	300	ST
	31.6984	2.2144	17.472 ± 0.030	500	ST
	31.7192	2.2352	17.496 ± 0.031	300	ST

**Table 2.** Optical Transient (contd.)

31.7410	2.257	$17.4503 \pm 0.028$	900	HCT
31.7414	2.2574	$17.509 \pm 0.030$	300	ST
31.7598	2.2758	$17.508 \pm 0.029$	300	ST
31.7740	2.29	$17.4403 \pm 0.029$	900	HCT
31.7769	2.2929	$17.485 \pm 0.031$	300	ST
31.7970	2.313	$17.520 \pm 0.029$	300	ST
31.8046	2.3206	$17.4553 \pm 0.020$	900	HCT
31.8149	2.3309	$17.550 \pm 0.028$	400	ST
31.8343	2.3503	$17.4853 \pm 0.028$	900	HCT
31.8387	2.3547	$17.586 \pm 0.029$	400	ST
31.8600	2.376	$17.595 \pm 0.031$	400	ST
31.8806	2.3966	$17.617 \pm 0.022$	400	ST
2003	April			
1.5967	3.1127	$17.866 \pm 0.033$	300	ST
1.6277	3.1437	$17.749 \pm 0.038$	200	ST
1.7153	3.2313	$17.826 \pm 0.029$	400	ST
1.7207	3.2367	$17.831 \pm 0.027$	400	ST
1.7356	3.2516	$17.7193 \pm 0.031$	900	HCT
1.8582	3.3742	$17.791 \pm 0.037$	300	ST
1.9000	3.416	$17.7843 \pm 0.034$	900	HCT
2.6386	4.1546	$18.3483 \pm 0.036$	900	HCT
2.6436	4.1596	$18.419 \pm 0.041$	300	ST
2.6771	4.1931	$18.4003 \pm 0.029$	900	HCT
2.6862	4.2022	$18.534 \pm 0.039$	300	ST
2.7439	4.2599	$18.528 \pm 0.038$	300	ST
2.7707	4.2867	$18.506 \pm 0.039$	300	ST
4.6673	6.1833	$18.787 \pm 0.039$	400	ST
4.6966	6.2126	$18.748 \pm 0.034$	400	ST
4.7246	6.2406	$18.852 \pm 0.033$	400	ST
5.6125	7.1285	$19.2573 \pm 0.067$	900	HCT
5.6511	7.1671	$19.2373 \pm 0.053$	900	HCT
5.7416	7.2576	$19.2963 \pm 0.046$	900	HCT
5.7702	7.2862	$19.3183 \pm 0.044$	900	HCT
5.8037	7.3197	$19.3153 \pm 0.046$	900	HCT
5.8293	7.3453	$19.3683 \pm 0.030$	900	HCT
5.8618	7.3778	$19.3873 \pm 0.045$	900	HCT
6.6438	8.1598	$19.6213 \pm 0.027$	900	HCT
6.6559	8.1719	$19.6443 \pm 0.028$	900	HCT
6.7522	8.2682	$19.6063 \pm 0.030$	900	HCT
6.8384	8.3544	$19.6183 \pm 0.039$	900	HCT
6.8505	8.3665	$19.6553 \pm 0.034$	900	HCT

**Table 2.** Optical Transient (contd.)

	7.6122	9.1282	$19.761 \pm 0.144$	400	ST
	8.6260	10.142	$19.983 \pm 0.069$	600	ST
	8.6641	10.1801	$19.943 \pm 0.076$	600	ST
	8.6933	10.2093	$19.919 \pm 0.064$	600	ST
	10.6368	12.1528	$20.287 \pm 0.087$	900	ST
	10.7059	12.2219	$20.1933 \pm 0.030$	600X4	HCT
	11.6782	13.1942	$20.432 \pm 0.177$	600	ST
2003	May				
	1.6271	33.1431	$22.214 \pm 0.156$	3*500+900	ST
<b>V – passband</b>					
2003	March				
	29.6310	0.147	$13.8983 \pm 0.029$	600	HCT
	29.6363	0.1523	$13.848 \pm 0.020$	200	ST
	29.6515	0.1675	$13.930 \pm 0.023$	75	ST
	29.6872	0.2032	$14.087 \pm 0.021$	50	ST
	29.6991	0.2151	$14.142 \pm 0.017$	100	ST
	29.7127	0.2287	$14.213 \pm 0.018$	100	ST
	29.7284	0.2444	$14.266 \pm 0.033$	100	ST
	29.7416	0.2576	$14.331 \pm 0.018$	100	ST
	29.7552	0.2712	$14.371 \pm 0.032$	100	ST
	29.7985	0.3145	$14.569 \pm 0.032$	200	ST
	29.8773	0.3933	$14.9323 \pm 0.050$	600	HCT
	30.6674	1.1834	$16.7423 \pm 0.047$	600	HCT
	30.7486	1.2646	$16.7943 \pm 0.050$	600	HCT
	30.7632	1.2792	$16.790 \pm 0.013$	200	ST
	30.7809	1.2969	$16.789 \pm 0.013$	200	ST
	30.7836	1.2996	$16.8153 \pm 0.032$	600	HCT
	30.7980	1.314	$16.773 \pm 0.013$	200	ST
	30.8202	1.3362	$16.825 \pm 0.013$	200	ST
	30.8256	1.3416	$16.8363 \pm 0.040$	600	HCT
	30.8680	1.384	$16.8183 \pm 0.046$	600	HCT
	30.8998	1.4158	$16.8353 \pm 0.050$	600	HCT
	31.6026	2.1186	$17.111 \pm 0.013$	300	ST
	31.6531	2.1691	$17.1033 \pm 0.064$	600	HCT
	31.6811	2.1971	$17.094 \pm 0.024$	300	ST
	31.7041	2.2201	$17.126 \pm 0.018$	300	ST
	31.7263	2.2423	$17.148 \pm 0.017$	300	ST
	31.7308	2.2468	$17.1783 \pm 0.055$	600	HCT
	31.7459	2.2619	$17.150 \pm 0.016$	300	ST
	31.7589	2.2749	$17.1653 \pm 0.052$	600	HCT

**Table 2.** Optical Transient (contd.)

Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope	
	31.7643	2.2803	17.148 ± 0.018	300	ST
	31.7814	2.2974	17.157 ± 0.020	300	ST
	31.7894	2.3054	17.1883 ± 0.040	600	HCT
	31.8015	2.3175	17.179 ± 0.014	300	ST
	31.8211	2.3371	17.212 ± 0.014	400	ST
	31.8216	2.3376	17.2143 ± 0.062	600	HCT
	31.8444	2.3604	17.227 ± 0.016	400	ST
	31.8498	2.3658	17.2443 ± 0.069	600	HCT
	31.8656	2.3816	17.253 ± 0.014	400	ST
	31.8781	2.3941	17.2633 ± 0.055	600	HCT
	31.8862	2.4022	17.253 ± 0.020	400	ST
	31.9177	2.4337	17.2843 ± 0.059	600	HCT
2003	April				
	1.6139	3.1299	17.438 ± 0.017	300	ST
	1.6727	3.1887	17.453 ± 0.017	300	ST
	1.6758	3.1918	17.4463 ± 0.051	450	HCT
	1.7215	3.2375	17.4203 ± 0.051	450	HCT
	1.7326	3.2486	17.453 ± 0.016	400	ST
	1.7381	3.2541	17.464 ± 0.016	400	ST
	1.8841	3.4001	17.427 ± 0.018	300	ST
	1.8847	3.4007	17.4233 ± 0.040	450	HCT
	2.6234	4.1394	18.0493 ± 0.061	450	HCT
	2.6595	4.1755	18.048 ± 0.021	300	ST
	2.6612	4.1772	18.0873 ± 0.065	450	HCT
	2.6903	4.2063	18.0913 ± 0.064	450	HCT
	2.6909	4.2069	18.068 ± 0.020	300	ST
	2.7221	4.2381	18.112 ± 0.024	300	ST
	4.6767	6.1927	18.370 ± 0.024	400	ST
	4.7069	6.2229	18.393 ± 0.025	400	ST
	4.7302	6.2462	18.481 ± 0.023	400	ST
	5.5977	7.1137	18.8733 ± 0.036	450	HCT
	5.6670	7.183	18.8953 ± 0.039	450	HCT
	5.6997	7.2157	18.9383 ± 0.052	450	HCT
	5.7303	7.2463	18.9283 ± 0.055	450	HCT
	5.7573	7.2733	18.9353 ± 0.044	450	HCT
	6.6149	8.1309	19.1813 ± 0.027	600X2	HCT
	6.7218	8.2378	19.2393 ± 0.027	600X3	HCT
	6.8111	8.3271	19.2183 ± 0.046	600X3	HCT
	6.8767	8.3927	19.2253 ± 0.064	600X2	HCT
	7.6073	9.1233	19.294 ± 0.073	300	ST

**Table 2.** Optical Transient (contd.)

Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope	
	8.6170	10.133	19.364 ± 0.033	600	ST
	8.6567	10.1727	19.406 ± 0.037	600	ST
	8.6856	10.2016	19.235 ± 0.035	600	ST
	9.5952	11.1112	19.414 ± 0.070	400	ST
	9.6050	11.121	19.5083 ± 0.032	900	HCT
	9.6439	11.1599	19.5233 ± 0.040	900	HCT
	9.6730	11.189	19.5233 ± 0.035	300X2	HCT
	9.7002	11.2162	19.5283 ± 0.044	700	HCT
	9.8625	11.3785	19.5463 ± 0.046	900	HCT
	10.6424	12.1584	19.585 ± 0.067	900	ST
	10.6461	12.1621	19.6303 ± 0.036	900X2	HCT
	10.6656	12.1816	19.6513 ± 0.058	450X3	HCT
	10.7440	12.26	19.6313 ± 0.060	900	HCT
	11.6590	13.175	19.7303 ± 0.043	450X2	HCT
	11.6925	13.2085	19.7763 ± 0.039	450X3	HCT
	11.6975	13.2135	19.951 ± 0.174	600	ST
	22.6520	24.168	20.5103 ± 0.048	600X3	HCT
	23.6646	25.1806	20.6343 ± 0.050	300X2+600	HCT
2003	May				
	1.6383	33.1543	20.836 ± 0.123	900	ST
<b><i>R</i>- passband</b>					
2003	March				
	29.6133	0.1293	13.3803 ± 0.031	120	HCT
	29.6171	0.1331	13.3743 ± 0.041	300	HCT
	29.6227	0.1387	13.431 ± 0.032	100	ST
	29.6232	0.1392	13.4123 ± 0.060	450	HCT
	29.6264	0.1424	13.422 ± 0.025	200	ST
	29.6397	0.1557	13.506 ± 0.024	200	ST
	29.6483	0.1643	13.577 ± 0.024	50	ST
	29.6530	0.169	13.619 ± 0.034	50	ST
	29.6572	0.1732	13.636 ± 0.040	100	ST
	29.6701	0.1861	13.7243 ± 0.034	600	HCT
	29.6784	0.1944	13.7233 ± 0.044	600	HCT
	29.6839	0.1999	13.724 ± 0.025	50	ST
	29.6863	0.2023	13.7893 ± 0.031	600	HCT
	29.6967	0.2127	13.829 ± 0.026	100	ST
	29.7103	0.2263	13.854 ± 0.026	100	ST

**Table 2.** Optical Transient (contd.)

Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope
29.7262	0.2422	13.916 $\pm$ 0.026	100	ST
29.7396	0.2556	13.982 $\pm$ 0.027	100	ST
29.7526	0.2686	14.030 $\pm$ 0.027	100	ST
29.7650	0.281	14.098 $\pm$ 0.052	200	ST
29.7937	0.3097	14.210 $\pm$ 0.032	100	ST
29.8675	0.3835	14.4893 $\pm$ 0.050	600	HCT
29.8857	0.4017	14.5483 $\pm$ 0.046	600	HCT
29.9079	0.4239	14.6333 $\pm$ 0.032	450	HCT
29.9140	0.43	14.6533 $\pm$ 0.033	450	HCT
29.9201	0.4361	14.6613 $\pm$ 0.040	450	HCT
29.9275	0.4435	14.6753 $\pm$ 0.070	450	HCT
30.5833	1.0993	16.249 $\pm$ 0.070	100	ST
30.5875	1.1035	16.295 $\pm$ 0.056	200	ST
30.6518	1.1678	16.2963 $\pm$ 0.070	180	HCT
30.6586	1.1746	16.3343 $\pm$ 0.042	600	HCT
30.6747	1.1907	16.3533 $\pm$ 0.039	450	HCT
30.6879	1.2039	16.3733 $\pm$ 0.028	450	HCT
30.7268	1.2428	16.3723 $\pm$ 0.049	450	HCT
30.7637	1.2797	16.4133 $\pm$ 0.040	450	HCT
30.7665	1.2825	16.420 $\pm$ 0.021	200	ST
30.7841	1.3001	16.438 $\pm$ 0.018	200	ST
30.7971	1.3131	16.4223 $\pm$ 0.050	450	HCT
30.8012	1.3172	16.439 $\pm$ 0.018	200	ST
30.8176	1.3336	16.4453 $\pm$ 0.041	450	HCT
30.8238	1.3398	16.460 $\pm$ 0.017	200	ST
30.8457	1.3617	16.4593 $\pm$ 0.039	450	HCT
30.8603	1.3763	16.4683 $\pm$ 0.034	450	HCT
30.8759	1.3919	16.4623 $\pm$ 0.032	450	HCT
30.8914	1.4074	16.4613 $\pm$ 0.039	450	HCT
30.9076	1.4236	16.4353 $\pm$ 0.057	450	HCT
30.9142	1.4302	16.4703 $\pm$ 0.034	450	HCT
30.9212	1.4372	16.4303 $\pm$ 0.031	450	HCT
30.9276	1.4436	16.4013 $\pm$ 0.063	450	HCT
31.6449	2.1609	16.7163 $\pm$ 0.040	450	HCT
31.6610	2.177	16.7163 $\pm$ 0.031	450	HCT
31.6653	2.1813	16.711 $\pm$ 0.032	100	ST
31.6804	2.1964	16.7103 $\pm$ 0.041	450	HCT
31.6898	2.2058	16.744 $\pm$ 0.033	200	ST

**Table 2.** Optical Transient (contd.)

Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope
31.7083	2.2243	16.771 ± 0.031	200	ST
31.7148	2.2308	16.783 ± 0.027	300	ST
31.7231	2.2391	16.7823 ± 0.026	450	HCT
31.7303	2.2463	16.772 ± 0.027	200	ST
31.7500	2.266	16.779 ± 0.026	200	ST
31.7507	2.2667	16.7953 ± 0.040	450	HCT
31.7655	2.2815	16.7933 ± 0.035	300	HCT
31.7682	2.2842	16.788 ± 0.029	200	ST
31.7826	2.2986	16.7953 ± 0.043	300	HCT
31.7853	2.3013	16.798 ± 0.029	200	ST
31.7943	2.3103	16.8203 ± 0.025	300	HCT
31.8054	2.3214	16.820 ± 0.025	200	ST
31.8139	2.3299	16.8083 ± 0.034	300	HCT
31.8264	2.3424	16.8443 ± 0.043	300	HCT
31.8280	2.344	16.844 ± 0.024	300	ST
31.8431	2.3591	16.8503 ± 0.040	300	HCT
31.8495	2.3655	16.867 ± 0.024	300	ST
31.8545	2.3705	16.8583 ± 0.048	300	HCT
31.8705	2.3865	16.874 ± 0.024	300	ST
31.8713	2.3873	16.8573 ± 0.044	300	HCT
31.8832	2.3992	16.8683 ± 0.045	300	HCT
31.8912	2.4072	16.900 ± 0.023	300	ST
31.8999	2.4159	16.8723 ± 0.059	300	HCT
31.9111	2.4271	16.9013 ± 0.037	300	HCT
31.9223	2.4383	16.9163 ± 0.028	300	HCT
2003	April			
1.6008	3.1168	17.071 ± 0.027	200	ST
1.6237	3.1397	17.045 ± 0.025	300	ST
1.6569	3.1729	17.059 ± 0.025	200	ST
1.6694	3.1854	17.0553 ± 0.046	300	HCT
1.6766	3.1926	17.075 ± 0.026	200	ST
1.6821	3.1981	17.0453 ± 0.045	300	HCT
1.6950	3.211	17.079 ± 0.026	200	ST
1.7274	3.2434	17.0473 ± 0.037	300	HCT
1.7445	3.2605	17.0583 ± 0.032	300	HCT
1.7487	3.2647	17.075 ± 0.029	300	ST
1.7531	3.2691	17.071 ± 0.026	300	ST
1.8628	3.3788	17.058 ± 0.023	300	ST
1.8891	3.4051	17.039 ± 0.024	300	ST
1.8906	3.4066	17.0273 ± 0.035	300	HCT

**Table 2.** Optical Transient (contd.)

Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope
1.9103	3.4263	$17.101 \pm 0.028$	300	ST
2.6158	4.1318	$17.6723 \pm 0.046$	300	HCT
2.6391	4.1551	$17.694 \pm 0.028$	300	ST
2.6425	4.1585	$17.6933 \pm 0.029$	300	HCT
2.6531	4.1691	$17.7113 \pm 0.029$	300	HCT
2.6553	4.1713	$17.707 \pm 0.025$	300	ST
2.6697	4.1857	$17.680 \pm 0.027$	300	ST
2.6843	4.2003	$17.7193 \pm 0.033$	300	HCT
2.6953	4.2113	$17.689 \pm 0.028$	300	ST
2.6962	4.2122	$17.7023 \pm 0.033$	300	HCT
2.7124	4.2284	$17.706 \pm 0.034$	300	ST
2.7265	4.2425	$17.709 \pm 0.030$	300	ST
2.7484	4.2644	$17.721 \pm 0.030$	300	ST
2.7615	4.2775	$17.737 \pm 0.030$	300	ST
2.7752	4.2912	$17.718 \pm 0.032$	300	ST
4.6831	6.1991	$18.030 \pm 0.031$	400	ST
4.7125	6.2285	$18.085 \pm 0.035$	400	ST
4.7357	6.2517	$18.130 \pm 0.031$	400	ST
5.5919	7.1079	$18.4583 \pm 0.049$	300	HCT
5.6044	7.1204	$18.5093 \pm 0.028$	300	HCT
5.6209	7.1369	$18.5033 \pm 0.025$	300	HCT
5.6340	7.15	$18.4963 \pm 0.032$	300	HCT
5.6603	7.1763	$18.4803 \pm 0.036$	300	HCT
5.6858	7.2018	$18.5763 \pm 0.029$	300	HCT
5.7054	7.2214	$18.5873 \pm 0.024$	300	HCT
5.7231	7.2391	$18.5243 \pm 0.029$	300	HCT
5.7507	7.2667	$18.5683 \pm 0.027$	300	HCT
5.7790	7.295	$18.5873 \pm 0.027$	300	HCT
5.7949	7.3109	$18.5803 \pm 0.036$	300	HCT
5.8130	7.329	$18.6003 \pm 0.027$	300	HCT
5.8454	7.3614	$18.5893 \pm 0.035$	300	HCT
5.8704	7.3864	$18.6283 \pm 0.037$	300	HCT
5.9002	7.4162	$18.6643 \pm 0.049$	300	HCT
6.5919	8.1079	$18.8233 \pm 0.039$	300	HCT
6.5995	8.1155	$18.8373 \pm 0.029$	300	HCT
6.6234	8.1394	$18.8053 \pm 0.036$	300	HCT
6.6849	8.2009	$18.8453 \pm 0.025$	600	HCT
6.6932	8.2092	$18.8503 \pm 0.030$	600	HCT

**Table 2.** Optical Transient (contd.)

Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope
6.7327	8.2487	$18.8503 \pm 0.032$	600	HCT
6.7413	8.2573	$18.8443 \pm 0.022$	600	HCT
6.7792	8.2952	$18.8363 \pm 0.029$	600	HCT
6.7878	8.3038	$18.8433 \pm 0.034$	600	HCT
6.8205	8.3365	$18.8523 \pm 0.029$	600	HCT
6.8286	8.3446	$18.8473 \pm 0.047$	600	HCT
6.8851	8.4011	$18.8823 \pm 0.043$	600	HCT
6.8935	8.4095	$18.8773 \pm 0.038$	600	HCT
6.9016	8.4176	$18.8183 \pm 0.031$	600	HCT
7.5936	9.1096	$18.897 \pm 0.052$	300	ST
7.5980	9.114	$18.945 \pm 0.067$	300	ST
8.6070	10.123	$19.017 \pm 0.046$	400	ST
8.6483	10.1643	$19.049 \pm 0.041$	600	ST
8.6778	10.1938	$19.064 \pm 0.032$	600	ST
9.5890	11.105	$19.078 \pm 0.057$	400	ST
9.5905	11.1065	$19.1263 \pm 0.033$	900	HCT
9.6321	11.1481	$19.0963 \pm 0.037$	900	HCT
9.6875	11.2035	$19.1403 \pm 0.037$	300X2	HCT
9.7242	11.2402	$19.1743 \pm 0.039$	300X3	HCT
9.8502	11.3662	$19.1663 \pm 0.050$	900	HCT
10.6068	12.1228	$19.3113 \pm 0.030$	900	HCT
10.6157	12.1317	$19.440 \pm 0.095$	400	ST
10.6437	12.1597	$19.2923 \pm 0.052$	900	HCT
10.6483	12.1643	$19.294 \pm 0.057$	900	ST
10.6825	12.1985	$19.2553 \pm 0.050$	300X3	HCT
10.7272	12.2432	$19.3483 \pm 0.039$	300+450	HCT
10.7390	12.255	$19.3503 \pm 0.041$	300	HCT
10.7574	12.2734	$19.3283 \pm 0.049$	900	HCT
10.8689	12.3849	$19.3463 \pm 0.060$	900	HCT
11.6743	13.1903	$19.4273 \pm 0.038$	300X3	HCT
11.7064	13.2224	$19.527 \pm 0.154$	600	ST
11.7082	13.2242	$19.2813 \pm 0.053$	300X2	HCT
22.6028	24.1188	$20.1233 \pm 0.043$	600	HCT
22.6279	24.1439	$20.0573 \pm 0.040$	600X2	HCT
22.8033	24.3193	$20.0633 \pm 0.025$	600X3	HCT
23.6311	25.1471	$20.0983 \pm 0.032$	300X5	HCT
2003	May			
	1.6069	$\pm 20.692 \pm 0.090$	2*300	ST

**Table 2.** Optical Transient (contd.)

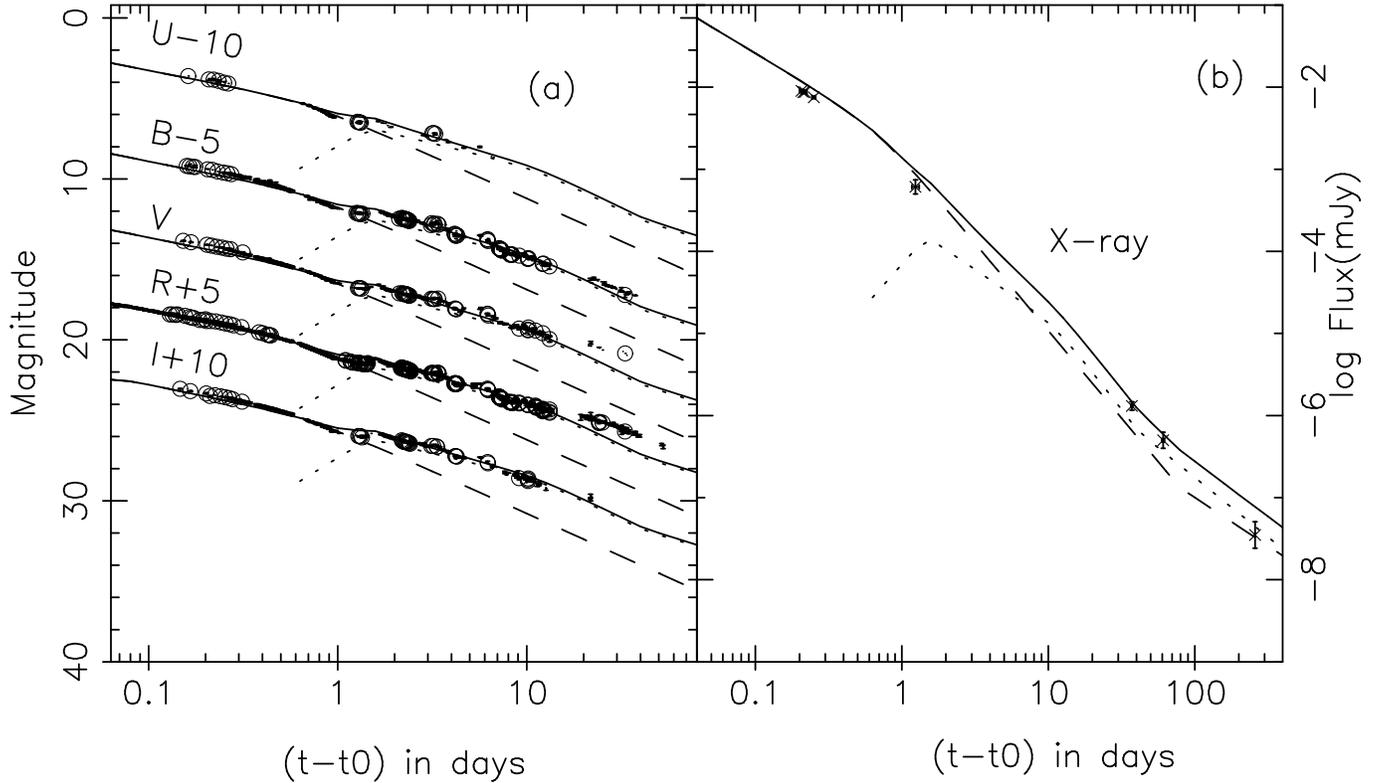
Date (UT)	Time since Burst in days	Magnitude (mag)	Exposure time (Seconds)	Telescope
<i>I</i> – passband				
2003	March			
	29.6306	13.037 ± 0.026	200	ST
	29.6499	13.185 ± 0.029	50	ST
	29.6855	13.328 ± 0.029	50	ST
	29.6948	13.487 ± 0.031	50	ST
	29.7083	13.437 ± 0.030	100	ST
	29.7241	13.501 ± 0.032	100	ST
	29.7374	13.559 ± 0.032	100	ST
	29.7504	13.618 ± 0.032	100	ST
	29.7616	13.685 ± 0.032	200	ST
	29.7959	13.837 ± 0.049	100	ST
	30.7698	15.972 ± 0.023	200	ST
	30.7872	15.985 ± 0.021	200	ST
	30.8046	15.978 ± 0.021	200	ST
	30.8272	16.058 ± 0.021	200	ST
	31.6683	16.240 ± 0.034	100	ST
	31.6933	16.290 ± 0.033	200	ST
	31.7110	16.305 ± 0.034	200	ST
	31.7338	16.342 ± 0.031	200	ST
	31.7535	16.329 ± 0.031	200	ST
	31.7715	16.326 ± 0.033	200	ST
	31.7887	16.341 ± 0.030	200	ST
	31.8088	16.351 ± 0.030	200	ST
	31.8326	16.387 ± 0.029	300	ST
	31.8539	16.411 ± 0.028	300	ST
	31.8750	16.404 ± 0.028	300	ST
	31.8960	16.460 ± 0.028	300	ST
2003	April			
	1.6086	16.591 ± 0.032	200	ST
	1.6890	16.586 ± 0.033	200	ST
	1.8294	16.639 ± 0.029	300	ST
	1.8338	16.631 ± 0.029	300	ST
	2.6738	17.211 ± 0.031	300	ST
	2.7077	17.224 ± 0.032	300	ST
	2.7572	17.254 ± 0.035	300	ST
	4.6886	17.586 ± 0.046	400	ST
	4.7180	17.663 ± 0.042	400	ST
	7.6029	18.595 ± 0.098	300	ST
	8.5965	18.674 ± 0.077	400	ST

**Table 3.** Results of the ESO/SEST scans on GRB 030329

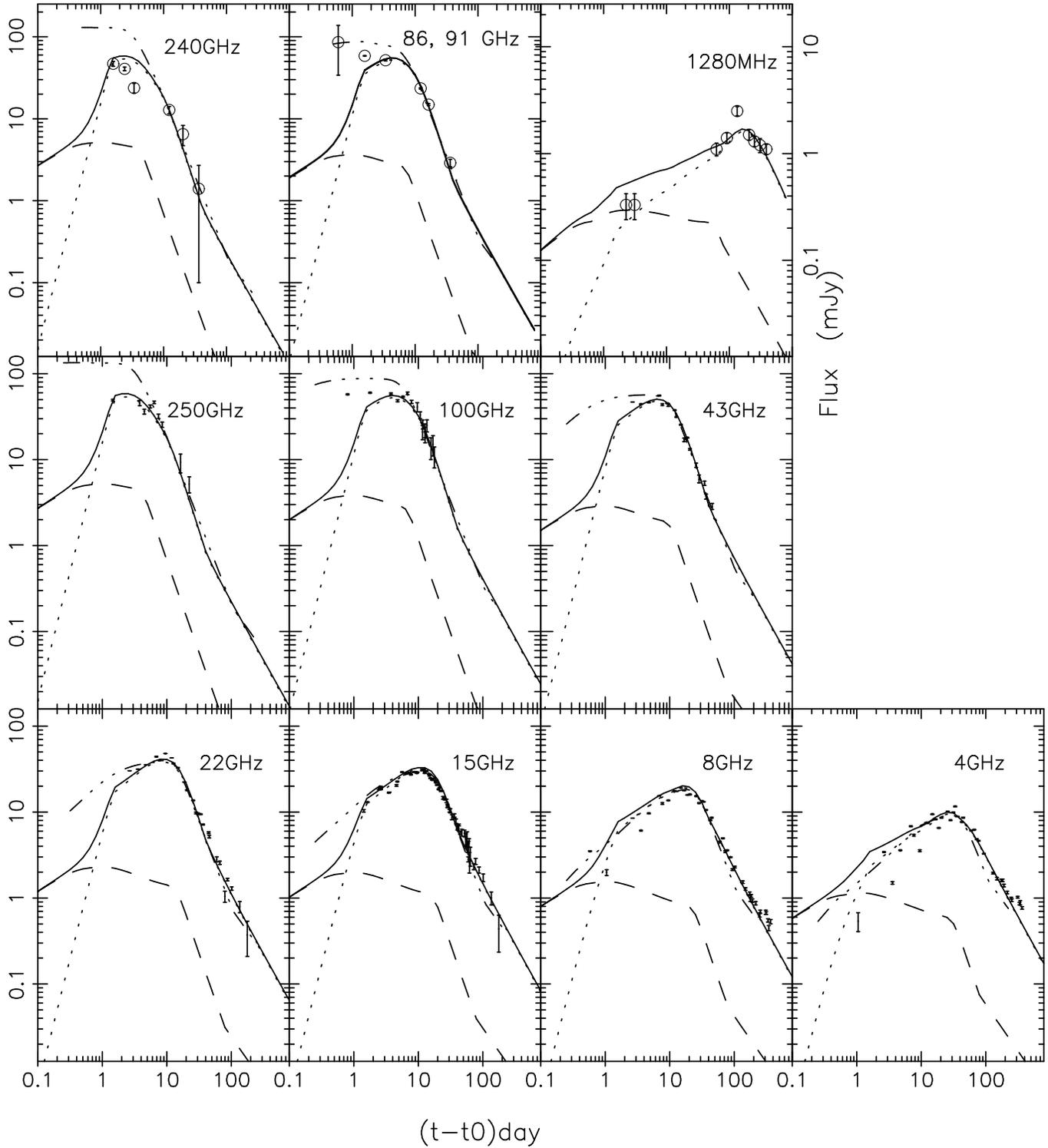
UT date of 2003	freq (GHz)	flux (mJy)	Beam
Mar 29/30	86	$82 \pm 52$	$57''$
	215	$201 \pm 308$	$23''$
Mar 30/31	86	$-2 \pm 78$	$57''$
	215	$-49 \pm 262$	$23''$
Mar 31/Apr 1	86	$38 \pm 65$	$57''$
	215	$-144 \pm 250$	$23''$

**Table 4.** Results of the IRAM/PdB scans on GRB 030329

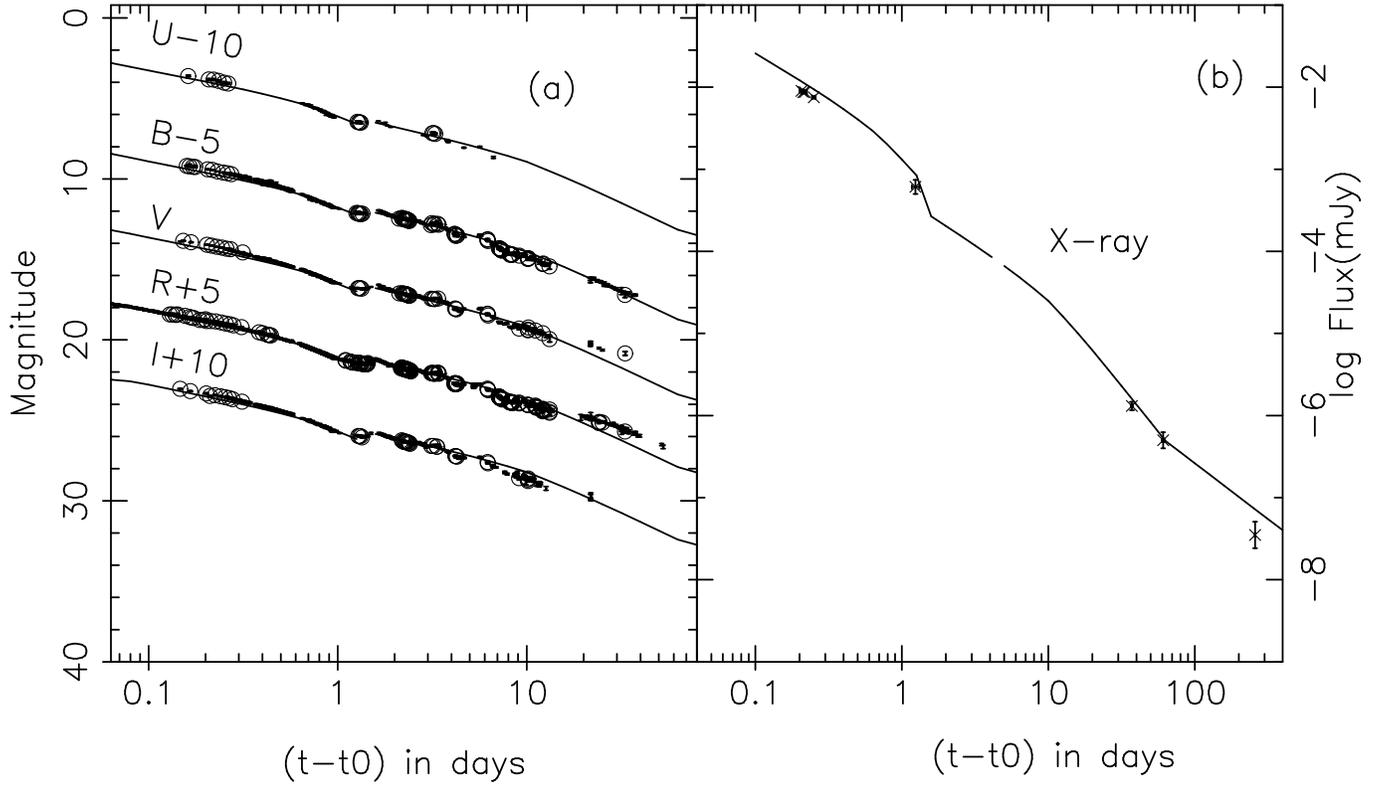
UT date of 2003	Config.	freq (GHz)	flux (mJy)	Beam and P.A.
31 Mar 00:04 to 03:07	6Dp	86.253	$58.6 \pm 0.5$	$14.1'' \times 4.0''$ at $55^\circ$
		232.032	$46.8 \pm 3.1$	$5.2'' \times 1.5''$ at $55^\circ$
31 Mar 18:33 to 23:57	6Dp	98.473	$58.2 \pm 0.6$	$6.4'' \times 4.2''$ at $98^\circ$
		238.500	$40.6 \pm 2.1$	$2.8'' \times 1.6''$ at $84^\circ$
01 Apr 20:20 to 21:55	6Dp	86.673	$51.7 \pm 0.9$	$9.1'' \times 4.2''$ at $98^\circ$
		240.528	$23.8 \pm 3.2$	$3.9'' \times 1.6''$ at $100^\circ$
05 Apr 16:54 to 18:04	6Dp	115.447	$40.4 \pm 3.7$	$10.3'' \times 3.1''$ at $-53^\circ$
10 Apr 17:42 to 19:12	6Dp	86.244	$23.5 \pm 0.4$	$11.0'' \times 4.2''$ at $-62^\circ$
		232.171	$12.8 \pm 1.3$	$4.5'' \times 1.5''$ at $-65^\circ$
14 Apr 20:10 to 22:16	6Dp	91.333	$14.9 \pm 0.4$	$8.1'' \times 4.2''$ at $84^\circ$
		217.029	$9.2 \pm 2.2$	$3.6'' \times 1.7''$ at $80^\circ$
18 Apr 19:39 to 21:10	6Dp	115.271	$7.7 \pm 1.0$	$6.3'' \times 3.3''$ at $90^\circ$
		232.032	$6.5 \pm 1.8$	$3.3'' \times 1.5''$ at $87^\circ$
24 Apr 17:36 to 19:32	6Dp	96.250	$4.7 \pm 0.7$	$8.2'' \times 4.0''$ at $-67^\circ$
		241.480	$0.0 \pm 3.8$	$3.5'' \times 1.4''$ at $105^\circ$
03 May 15:59 to 19:49	6Dp	86.243	$2.9 \pm 0.3$	$9.5'' \times 4.4''$ at $108^\circ$
		233.467	$1.4 \pm 1.3$	$3.6'' \times 1.6''$ at $-76^\circ$
16 May 13:20 to 15:57	5Dp	86.243	$1.1 \pm 0.8$	$12.4'' \times 5.3''$ at $-54^\circ$
		231.490	$0.0 \pm 10.1$	$4.6'' \times 1.9''$ at $-59^\circ$
28 May 21:29 to 23:22	5Dp	84.443	$0.4 \pm 0.7$	$16.6'' \times 5.5''$ at $52^\circ$
		238.500	$0.0 \pm 8.0$	$5.5'' \times 1.9''$ at $52^\circ$
20 June 12:41 to 17:49	4Dp	95.434	$0.6 \pm 1.3$	$5.92'' \times 5.63''$ at $64^\circ$



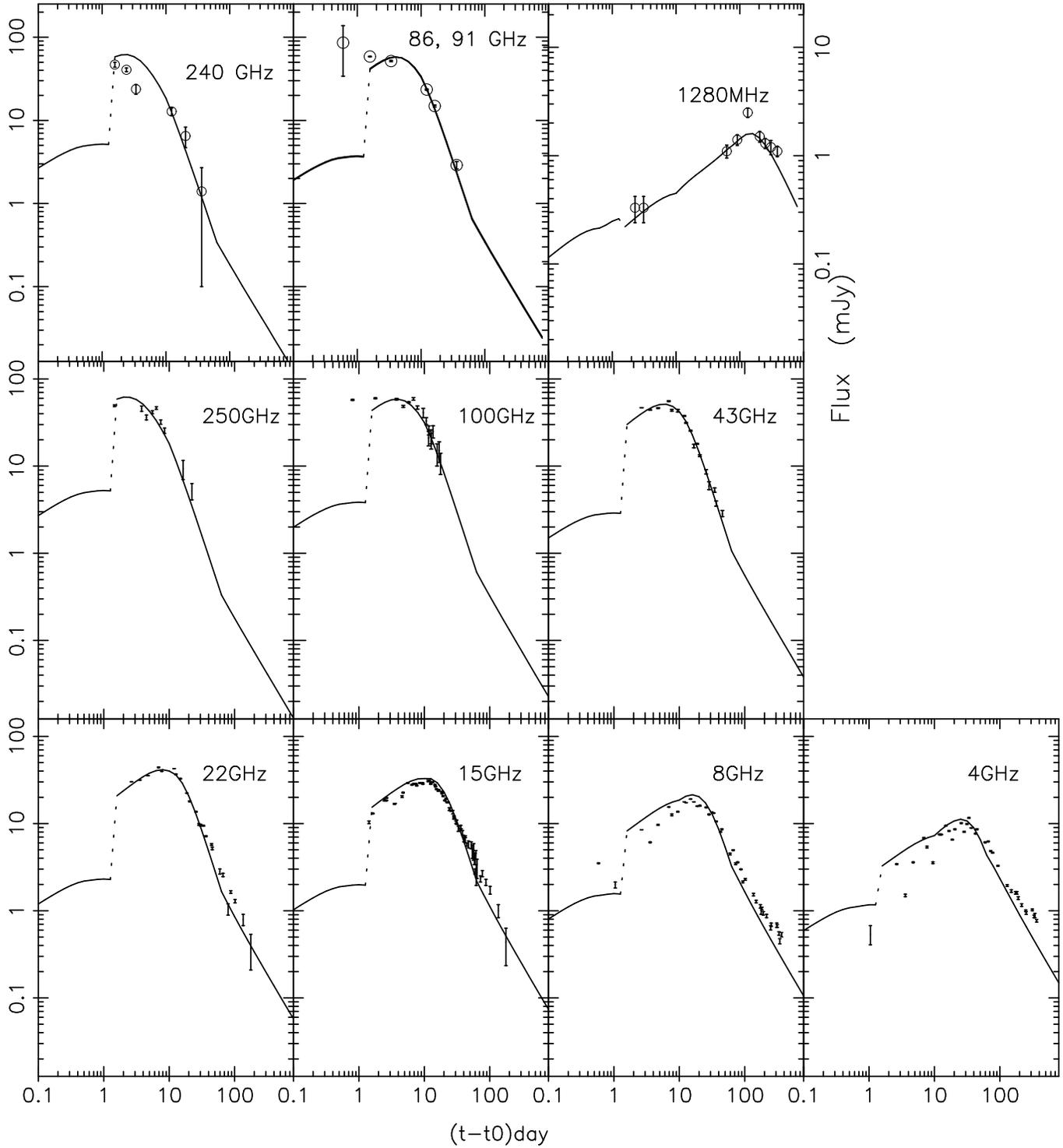
**Fig. 1.** (a) The Optical lightcurve of the afterglow of GRB030329. Open circles represent the data presented in this paper and filled circles are those from the literature (Lipkin et al. 2004). The solid line shows the total flux predicted by model 1 discussed in the text, which has two jets as in Berger et al. (2003). The dashed line shows the contribution of the narrow jet alone and the dotted line that of the wide jet in each band. (b) X-ray observations reported by Tiengo et al. 2003 and 2004, shown along with the prediction of model 1. Contribution of the narrow jet and the wide jet are shown separately as the dashed and the dotted line respectively. The total flux is shown as the solid curve.



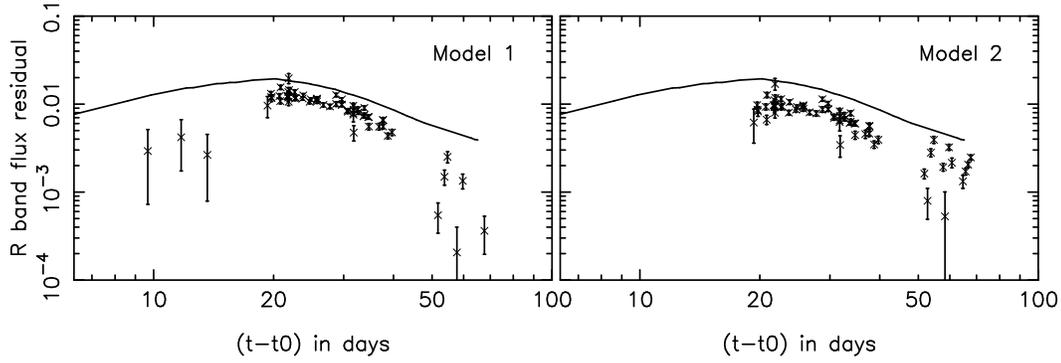
**Fig. 2.** Millimeter and radio observations of GRB030329 afterglow along with the predictions of model 1 (two jets). Observations reported in this paper are represented by open circles, crosses are data points from Berger et al. (2003) and Sheth et al. (2003). The dashed and the dotted lines represent contributions of the narrow and the wide jet respectively, the solid line shows the sum. The dash-dot-dot line shows a model fit for an assumed stellar-wind density profile for the ambient medium.



**Fig. 3.** (a) The Optical lightcurve of the afterglow of GRB030329, shown with the prediction of model 2 (solid line), which assumes a transition of an initially narrow jet to a wider jet at  $\sim 1.5$  days. (b) X-ray observations reported by Tiengo et al. 2003 and 2004, with predictions of model 2. The flattening seen at late times is due to the transition into non-relativistic regime at  $\sim 63$  days.



**Fig. 4.** Millimeter and radio observations of GRB030329, along with the predictions of model 2 (refreshed jet) shown as the solid line. Open circles are data presented in this paper, crosses represent data taken from Berger et al (2003) and Sheth et al (2003).



**Fig. 5.** *R*-band residuals for epochs beyond  $\sim 7$  days, after subtracting the modeled flux of the afterglow and the contribution of the host galaxy ( $R = 22.6$ , Gorosabel et al. 2005b) from the observed flux. The two models are shown in adjacent panels. This shows the *R*-band contribution needed from the associated supernova SN2003dh to explain the total observed light from the OT. The solid line is the red-shifted *K*-corrected SN1998bw *R*-band lightcurve, shown for comparison.