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UOCS – XI. Study of blue straggler stars in open cluster NGC 7142 using UVIT/AstroSat

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ABSTRACT

We present a study of blue straggler stars (BSSs) of open cluster NGC 7142 using *AstroSat*/UVIT data and other archival data. Using a machine-learning-based algorithm, ML-MOC, on *Gaia* DR3 data, we find 546 sources as cluster members. Based on the location on the *Gaia* colour–magnitude diagram, we identify 10 BSS candidates, also detected in UVIT/F148W filter. We study the variable nature of BSSs by constructing their light curves using the *Transiting Exoplanet Survey Satellite* data. Two BSSs reported as eclipsing binaries in *Gaia* DR3 are confirmed to be eclipsing binaries based on our analysis and also show the presence of hot companions as per the multiwavelength spectral energy distributions (SEDs). The physical parameters of the hot companions of these two BSSs derived by fitting binary models to their light curves and those derived from the SEDs are found to be in good agreement. Additionally, five more BSSs in the cluster shows UV excess, four of which are likely to have a hot companion based on SEDs. The hot companions with the estimated temperatures ~14 000–28 000 K, radii ~0.01–0.05 R_☉, and luminosities ~0.03–0.1 L_☉ are inferred to be extremely low-mass (<0.2 M_☉), low-mass (0.2–0.4 M_☉), normal-mass (0.4–0.6 M_☉), and high-mass (>0.6 M_☉) white dwarfs (WDs). For the first time in an open cluster, we find the entire range of masses in WDs found as hot companions of BSSs. These masses imply that the Case-A/Case-B mass transfer and merger are responsible for the formation of at least 60 per cent of the BSSs of this cluster.

Key words: binaries: general-blue stragglers-white dwarfs-open clusters and associations: individual: NGC 7142ultraviolet: stars.

1 INTRODUCTION

Blue straggler stars (BSSs) are the enigmatic stellar population that appears younger than other stars in their host clusters. They are identified by their location on the colour–magnitude diagram (CMD), where they appear brighter and bluer than the cluster main-sequence turn-off (MSTO). They were first identified in a globular cluster (GC) M3 in 1953 (Sandage 1953). It is believed that BSSs can form through a variety of mechanisms, including stellar collisions (Hills & Day 1976; Chatterjee et al. 2013), mass transfer (MT; McCrea 1964), and merger in hierarchical triple system (Perets & Fabrycky 2009). The collisions can result in the formation of a single, more massive star that appears younger than the surrounding stars in the cluster. It is observed in dense stellar environments such as cores of the GCs (Hurley et al. 2005; Chatterjee et al. 2013; Hypki & Giersz 2013). In the case of the MT formation channel, the BSS progenitor fills its Roche lobe and accrete mass from its companion star. This mechanism can further be classified into three sub-categories depending on the stages of the donor during MT. In Case-A MT, the donor is in the main-sequence phase (Webbink 1976) while transferring the mass. On the other hand, in Case-B MT, the donor is in the red-giant-branch phase during MT and eventually forms a He core white dwarf (WD) (McCrea 1964). Finally, in Case-C MT, the donor transfers the mass while in the asymptotic-giantbranch phase (Chen & Han 2008), and leaves behind a CO WD. When it comes to a merger, the Kozai mechanism and tidal friction may cause the formation of closely bound inner binaries in a hierarchical triple system. Over time, angular momentum loss due to either stellar evolution or a magnetized wind can cause the binaries to merge or undergo MT between each other. This process can lead to the creation of BSS in long binary or triple systems, as discussed in studies by Kiseleva, Eggleton & Mikkola (1998), Fabrycky & Tremaine (2007), and Naoz & Fabrycky (2014).

Open clusters (OCs) are ideal environments for investigating the structure and history of the Galactic disc, as well as for studying the development and evolution of both single and binary star populations. The dynamical interactions among the stellar populations within a cluster can result in the formation of various exotic stellar objects,

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including BSSs (Ahumada & Lapasset 2007). OCs are closer and less populated than GCs; hence, observing the properties of BSSs in OCs can help to narrow down their formation mechanisms. Moreover, by studying the BSSs population in OCs, a better understanding of the age distribution and star formation history of the host cluster can also be achieved (De Marchi et al. 2006). The ultraviolet (UV) wavelengths are crucial in detecting BSSs because they emit a significant amount of flux within the UV range. Furthermore, the UV observations can be used to detect their hotter companions that may be present. The discovery of hot companions such as extremely low-mass (ELM; $M < 0.2 M_{\odot}$) and low-mass (LM; $M < 0.4 M_{\odot}$) WDs suggests that the BSSs were formed via Case-A or Case-B MT in binary systems.

The Ultraviolet Imaging Telescope (UVIT) onboard *AstroSat* is capable of detecting BSSs and their hot companions when combined with other multiwavelength data. Several studies on OCs have been conducted in the past in this context. Subramaniam et al. (2016a), for example, reported a post AGB/horizontal-branch companion of a BSS in the OC NGC 188. Sindhu et al. (2019), Jadhav, Sindhu & Subramaniam (2019), and Pandey, Subramaniam & Jadhav (2021) investigated the OC M67 and discovered ELM WDs as BSS companions. Similarly, Vaidya et al. (2022) discovered ELM WDs as hot companions of BSSs in the intermediate-age OC NGC 7789. In the case of NGC 2506, two ELM WDs and one LM WD were discovered as the BSSs hot companions and normal and high-mass WDs as the hot companions of YSS (Panthi et al. 2022). Recently, Rani et al. (2023) reported A-type sub-dwarfs as hotter companions of YSS in the OC NGC 2818.

NGC 7142 (RA = 326.287 deg, Dec. = 65.775 deg) is an old OC located at the distance of ~ 2 kpc (Sun et al. 2020) and is surrounded by molecular and dust clouds (Straižys et al. 2014). It is present within 0.4° of a well-known nebula, NGC 7129, and has been the subject of several photometric and spectroscopic studies aimed at understanding the properties and evolution of different stellar populations. Hoag (1961) presented the initial UBV photometric investigation of this cluster and showed the CMD with the red giants sequences and several stars near the MSTO. Another significant photometric study was done by Van den Bergh & Heeringa (1970), who showed the resemblance of its CMD with the renowned OCs NGC 188 and M67. They reported the age of the cluster to be 1.5-4 Gyr and [Fe/H] = -0.45 ± 0.2 dex, and the distance modulus $(m - M)_V$ ranging from 11.8 to 13.7 mag. Crinklaw & Talbert (1991) published the first CCD photometry of the cluster using the BV bands in their search for variable stars. They affirmed that variable reddening exists and suggested that its extent is less than 0.1 mag in E(B - V). They concluded that this cluster is older than M67 (4.0 \pm 0.5 Gyr) but younger than NGC 188 (7 \pm 0.5 Gyr old), and reported $(m - M)_V = 11.4$ mag. Later, Carraro, Ng & Portinari (1998) compared the synthetic CMDs generated from synthetic model isochrones and reported an age of 4.9 Gyr. Salaris, Weiss & Percival (2004) derived an age of 4 ± 1 Gyr based on the difference in magnitude between the sub-giant branch and the red giant clump. However, Janes & Hoq (2011) recently estimated an age of 6.9 ± 0.8 Gyr using synthetic CMDs. It is noteworthy that the differential extinction in the cluster (Crinklaw & Talbert 1991; Straižys et al. 2014) and the sparse populations of evolved stars make it challenging to locate the reference points in the CMD and may account largely for the discrepancy in the age determination of this cluster. Recently, Sun et al. (2020) presented UBVRI photometry of this cluster using WIYN 0.9 m telescope and reported E(B - E)V) = 0.338 ± 0.031 mag, [Fe/H] = 0.0 ± 0.1 dex, and age = 4.0^{-0.7}_{±13} Gyr. The initial variability study of NGC 7142 had been carried out by Crinklaw & Talbert (1991) followed by Rose & Hintz (2007). The former discovered one short-period variable star near the cluster MSTO, whereas the latter found eight low-amplitude variable stars. Later, Sandquist, Serio & Shetrone (2011) presented the discovery of eight contact or near-contact eclipsing binaries, one long-period variable, and four detached eclipsing binaries in the cluster.

Although NGC 7142 has been extensively researched, the BSSs in this cluster has not been subjected to detailed studies. Furthermore, the variability of the BSSs in this cluster has not been comprehensively explored. As mentioned above, NGC 7142 resembles in properties with the thoroughly examined old clusters NGC 188 and M67; therefore, it is important to explore this cluster as well. Given the importance of UV wavelengths to study BSSs, we present the UVbased analysis of BSSs of NGC 7142 for the first time. We also search for variability signatures of BSSs using *Transiting Exoplanet Survey Satellite (TESS)* data. Recently, a similar analysis was presented by Vernekar et al. (2023), where they studied the light curves (LCs) of five BSSs that are known to be SB1s in the OC M67 using Kepler and *TESS* data. From the LCs, they provided supporting evidence of the MT as the formation mechanism of BSSs to the earlier multiwavelength spectral energy distribution (SED) analysis.

This paper is structured as follows: In Section 2, the observational data and the analysis are described. The data analysis is presented in Section 3. The results and discussions are given in Section 4. Finally, the summary and conclusions of the work are presented in Section 5.

2 OBSERVATIONAL DATA AND REDUCTION

2.1 Ultraviolet Imaging Telescope

The observations of NGC 7142 were carried out using UVIT in one FUV filter, F148W, on 2018 June 13 under the proposal ID A04-075 with the exposure time of 2280.717 s. UVIT is an imaging instrument onboard AstroSat, which is made up of two Ritchey-Chrétien telescopes with 38 cm apertures. In UVIT, one telescope observes in FUV (130-180 nm), and the other in NUV (200-300 nm) and visible wavelengths (320-550 nm). However, the NUV channel developed a technical issue in 2018 and has not been functional since then. UVIT produces images with a full width at half-maximum (FWHM) of \sim 1.5 arcsec and a circular field of view of 28 arcmin in diameter. There are various filters available in the FUV, NUV, and visible bands. The UV detectors operate in photon-counting mode, whereas the visible detectors operate in integration mode. For the details on the effective area curves of UVIT filters and other instrument-related information, readers are referred to Kumar et al. (2012), Subramaniam et al. (2016b), and Tandon et al. (2017). We use CCDLAB (Postma & Leahy 2017, 2021), a customized software package, to correct for geometric distortion, flat field, spacecraft drift, astrometry and create images for each orbit. The orbit-wise images were then co-aligned and combined to generate science-ready images.

2.2 Transiting Exoplanet Survey Satellite

TESS is an all-sky broad-band photometric survey mission covering upto 85 per cent of the sky in its first 2 yr. Each hemisphere is divided into 13 sectors, with one sector being observed for 27.4 d. It provided the data in two cadences, 2 and 30 min, in the 2-yr nominal mission, but was later augmented to 20 and 200 s, in the extended mission. It has four identical cameras, each with a field of view of 24×24 deg. Each of the four cameras is equipped with a $2K \times 2K$ CCD, with each pixel covering 21 arcsec. The cameras are arranged so that all four can cover a 24×96 square degree strip of sky. The detectors are sensitive from 600 to 1000 nm, i.e. within blue to the near-IR wavelengths.

TESS LCs of all our BSS samples are observed using different data projects. These include Quick Look Pipeline (QLP; Huang et al. 2020; Kunimoto et al. 2022), Cluster Difference Imaging Photometric Survey (CDIPS; Bouma et al. 2019), and a point spread function (PSF)-based Approach to TESS High-Quality Data Of Stellar Clusters (PATHOS; Nardiello et al. 2019, 2021). The OLP project generates LCs from a group of stars and other stationary luminous objects with a limited G magnitude of 13.5. For the details on the data reduction process, readers are referred to Huang et al. (2020). On the other hand, the CDIPS project creates the LCs for stars that are candidate members of OCs and moving groups. Each LC corresponds to 20 to 25 d of observations of a star brighter than 16 in G magnitude. The PATHOS project generates a library of high-precision LCs for cluster members using photometry extracted from TESS Full Frame Pictures. The PSF-based technique reduces dilution effects in crowded situations, allowing for the extraction of high-precision photometry for faint stars. The PATHOS data release X has 770 LCs of 168 star clusters, including NGC 7142 (Messina et al. 2022).

3 DATA ANALYSIS

3.1 Membership determination

In the case of OCs, the contamination due to field stars is a big challenge. Hence, obtaining secure membership is of utmost importance to do any further analysis. We used the ML-MOC algorithm (Agarwal et al. 2021) on Gaia DR3 (Vallenari, Brown & Prusti 2022) data in order to determine the cluster members. This is a machine-learning-based algorithm that does not require any prior knowledge of a cluster and employs the k-nearest neighbour (kNN, Cover & Hart 1967) and the Gaussian mixture model (GMM, Peel & McLachlan 2000) algorithms. The steps that we followed to determine cluster members using ML-MOC are described briefly here. First, we downloaded the sources within 30 arcmin radius of the cluster centre that have five astrometric parameters (positions, proper motions, and parallax), appropriate measurements in three Gaia photometric passbands G, G_{BP} , and G_{BP} , non-negative parallaxes, and error in G-mag less than 0.005. These are called All sources. Then, we estimated the mean proper motions and mean parallaxes of the cluster using the kNN algorithm to determine the ranges of proper motions and parallaxes that enclose all of the likely cluster members. These are termed as Sample sources. Next, we apply the GMM, an unsupervised clustering algorithm, to the Sample sources in order to separate the likely cluster members from the field stars by fitting two Gaussian distributions in the proper motions and parallax space of the Sample sources. The fitting of GMM is done by maximizing the likelihood of the distribution parameters using the expectation maximization algorithm (Dempster, Laird & Rubin 1977). The Kcomponent GMM, for N data points in an M-dimensional parameter space, is defined as

$$P(x) = \sum_{i=1}^{K} \omega_i G(x | \mu_i, \Sigma_i), \tag{1}$$

where P(x) represents the probability distribution of data points x and ω_i is the mixture weight of the *i*th Gaussian component, G(x| μ_i ,

 Σ_i). This Gaussian component is defined as

$$G(x|\mu_i, \Sigma_i) = \frac{\exp\left[-\frac{1}{2}(x-\mu_i)^T \Sigma_i^{-1}(x-\mu_i)\right]}{(2\pi)^{M/2} \sqrt{|\Sigma_i|}}$$
(2)

In both the above equations, μ_i is the mean vector, and Σ_i is the full covariance matrix of the ith Gaussian component. The GMM calculates a probability that designates the likelihood of each data point belonging to the cluster. We use a two-component, i.e. cluster and field, GMM on the normalized three-dimensional parameter space $(\mu\alpha, \mu\delta, \omega)$ assuming that the proper motions and parallaxes of the Sample sources follow a two-component Gaussian distribution. Our member selection criteria considers sources with high membership probabilities exceeding 0.6 as members. Subsequently, we expand this list to include Sample sources whose parallax values lie within the range of parallaxes of even higher probability members (>0.8). The members added later are low-probability members with probabilities ranging between 0.2 and 0.6. Out of the 546 sources identified as cluster members, only six are with probabilities ranging between 0.2 and 0.6; the rest are high probability members. The proper motion, spatial, and parallax distributions of sample sources and the identified cluster members are shown in Fig. 1.

3.2 Photometry and CMD

We used the DAOPHOT package of IRAF (Stetson 1987) in order to perform the PSF photometry and obtained the magnitudes of the detected stars in FUV image, F148W. To accomplish this task, we employed the following steps. First, we estimated the background counts and mean FWHM of the image using *imexamine* task. Then, we detected the sources above a certain threshold limit using *daofind* task in IRAF. We performed the aperture photometry of the sources using the phot task. In order to perform PSF photometry, we selected the bright and isolated sources using pstselect and created a PSF model using *psf* task. Finally, PSF magnitudes of all the stars were estimated simultaneously using multiple iterations of PSF fitting to the group of stars using allstar task. We also applied the PSF corrections in order to eliminate the systematic difference between allstar magnitudes and phot magnitudes, followed by the aperture correction, which was determined using the curve-of-growth technique. Moreover, since the FUV detector of UVIT works in photon-counting mode, it may be subjected to counts more than one photon per frame. In order to eliminate this issue, we applied the saturation corrections following Tandon et al. (2020).

To quantify the differential reddening and correct it, we used the method adopted in Della Croce et al. (2023). First, we select the likely main-sequence stars (G > 15 mag or $G_{BP} - G_{RP} < 1.5$ mag) from the members determined using ML-MOC. Then we find counterparts of these members in Pan-STARRS (Chambers et al. 2016) photometric bands r, i, and z by searching within 1 arcsec of the Gaia coordinates. Next, we plot a colour–colour diagram, G-rversus i-z. This diagram is a suitable choice as the evolutionary sequences are generally orthogonal to the reddening vector in these colours (Della Croce et al. 2023). We refer to the median of this colour-colour distribution as our reference point. Afterwards we obtain the median colours of the closest 50 stars of each mainsequence star and estimated the distance of this median value to the reference system along the reddening vector using the coefficients given in Cardelli, Clayton & Mathis (1989). For the remaining stars which are not in the main sequence, we assign the median reddening of the closest 50 neighbours (Della Croce et al. 2023). At last, we apply the extinction value corresponding to the derived distance to member stars and plotted the reddening map, as shown in the



Figure 1. The spatial distribution, proper motion, and parallax distributions of the members and sample sources determined using the ML-MOC algorithm.

left panel of Fig. 2. It is noteworthy that the reddening is more severe in the north-west field; however in the centre as well as in the southern regions, the reddening is low and does not vary significantly. The significantly high reddening in the north-west direction can be attributed to the presence of the NGC 7129 nebula in the same direction. The right panel of Fig. 2 shows the optical CMD of the cluster members determined using ML-MOC before and after the differential reddening and extinction corrections.

Fig. 3 shows the differential reddening and extinction-corrected optical CMD of the cluster. The cluster members determined using the machine-learning-based algorithm, ML-MOC, are shown in red dots along with their membership probabilities. The 10 BSSs identified as members and with counterparts in the UVIT/F148W are shown as blue open circles, labelled according to the coordinates given in Table 1. There are four additional BSSs candidates which are cluster members but have no UV counterparts. We only study the 10 BSSs, which are members of the cluster and detected in the UVIT filter in this work. It is noteworthy that all the BSSs are highly probable cluster members with a membership probability of ~ 0.8 -0.9. Furthermore, the variable stars identified in a study by Sandquist, Serio & Shetrone (2011), which are cluster members as per our criteria, are shown as green-filled circles. BSS 7 is one of the five variable stars from their study. The remaining four variables do not show variable nature using Gaia or TESS data. A PARSEC isochrone (Bressan et al. 2012) of age = 4.0 Gyr, distance = 2368 pc, and [Fe/H] = 0 (Sun et al. 2020) is plotted after applying the extinction correction of $A_G = 0.09$ and reddening of $E(B_P - R_P) = 0.03$. Furthermore, we have also shown a binary sequence isochrone as the blue dashed curve for equal-mass binaries that have a G-band magnitude brighter than 0.75 mag than the MS stars. It is evident from the CMD that G-band magnitudes of BSSs range from ~ 0.4 to \sim 1.4 above the MSTO.

3.3 Light-curve analysis

Variable stars can provide a wealth of new information, allowing us to confirm current theories that seek to predict cluster stellar properties from basic principles (Sandquist, Serio & Shetrone 2011). For example, pulsating stars can reveal information about stellar interior structure and mass. Short-period semidetached and contact binaries may carry encoded information about the dynamical history of the cluster. Detached eclipsing binaries can be particularly valuable since they can offer high-precision measurements of masses and radii for individual stars with minimal theoretical interpretation. In order to check if any of the BSS is variable in nature, we used the *TESS*

(Ricker et al. 2015) data. We searched for the TESS LCs of all our BSS samples in different data releases, as mentioned above. In this work, we have presented the LCs of PATHOS data release. We used Lightkurve (Lightkurve Collaboration 2018), a python based package, to construct and analyse the LC of the BSSs. We found that two of our 10 BSSs candidates (BSS 1 and BSS 7) showed some signatures of variability in sectors 16, 18, and 24, using all the three pipelines mentioned above. The LCs of these BSSs of sector 16, generated using the PATHOS pipeline, are shown in Fig. 4. Then, we computed the power spectrum using Fourier transform from the LC in order to determine the period as shown in Fig. 5. We take all the significant frequencies from the periodogram. A frequency is considered significant if the amplitude or power of that frequency is four times the noise (Breger et al. 1993). After doing it for all the frequencies, the orbital frequency will give a phase-folded curve with two dips resembling two eclipses, whereas all others will give sinusoidal variations as shown in Fig. 6. It can be noted that the orbital period of BSS 1 is \sim 1 d, whereas the period of BSS 7 is ~ 0.5 d. Sandquist, Serio & Shetrone (2011) categorized BSS 7 as an eclipsing contact or near-contact binary with an orbital period of ~ 0.58 d, using high-resolution spectrograph at Hobby–Eberly Telescope. However, none of their other variables were found to be variable in our analysis.

4 RESULTS AND DISCUSSION

4.1 Spectral energy distributions

Stars emit electromagnetic radiation across the entire frequency or wavelength range. The SED describes the distribution of this radiation over wavelengths. Analysing the SEDs of stars allows us to estimate their fundamental parameters and thus gain a better understanding of their formation and evolution. Furthermore, the excess in UV fluxes can indicate the likelihood of a hotter companion of BSSs, hence constraining their formation mechanism. We used the virtual observatory SED analyser (VOSA, Bayo, Rodrigo & y Navascus 2008) to create the SEDs. We found that three of the 10 BSS candidates (BSS 4, BSS 9, and BSS 10) in this cluster have a nearby source within 3 arcsec, and hence we did not construct their SEDs.

We obtained photometric fluxes of sources in NUV from GALEX (Martin et al. 2005), optical from *Gaia* DR3 (Babusiaux et al. 2023) and PAN-STARRS (Chambers et al. 2016), near-IR from Two Micron All Sky Survey (2MASS, Cohen, Wheaton & Megeath 2003), and far-IR from Wide-field Infrared Survey (WISE, Wright



Figure 2. The left panel shows the reddening map of the cluster members. The right panel shows the observed *Gaia* DR3 CMD as blue dots and differential reddening and extinction-corrected CMD as orange dots.



Figure 3. The differential reddening and extinction-corrected optical CMD of NGC 7142 showing members determined using ML-MOC, BSSs detected in UVIT/*F*148*W* filter, a PARSEC isochrone, and an equal mass binary sequence.

et al. 2010), using VOSA. The photometric fluxes were corrected for extinction by VOSA according to the extinction law by Fitzpatrick (1999) and Indebetouw et al. (2005) using the value of extinction A_v = 1.04 ± 0.09 provided by us, which is taken from Sun et al. (2020). Table 1 shows the values of the BSS fluxes in various filters.

VOSA calculates synthetic photometry for selected theoretical models using filter transmission curves and performs a χ^2 minimization test by comparing the synthetic photometry with the extinction-corrected observed fluxes to determine the best-fitting SED parameters. The following formula is used to calculate the reduced χ^2_r :

$$\chi_r^2 = \frac{1}{N - N_f} \sum_{i=1}^N \frac{(F_{o,i} - M_d F_{m,i})^2}{\sigma_{o,i}^2},$$
(3)

where N is the number of photometric data points used and N_f is the number of model-free parameters. The observed and model fluxes of the star are denoted by $F_{o,i}$ and $F_{m,i}$, respectively. M_d is the scaling factor that must be multiplied by the model to obtain the fit, and it is given by $(R/D)^2$, where R is the radius of the star, D is the distance to the star, and $\sigma_{o,i}$ is the error in the observed flux. We

used Kurucz stellar models (Castelli, Gratton & Kurucz 1997) to fit the SEDs to the BSSs. We kept T_{eff} and log g as free parameters, with ranges of 3500-50 000 K and 3-5, respectively. We set the metallicity ([Fe/H]) value to zero, which is the cluster metallicity reported by Sun et al. (2020). We first removed the UV data points from the SEDs and checked whether the optical and IR data points fit satisfactorily with the model flux. We carefully examined each UV filter to see if there was any excess to the UV and/or IR data points. We consider sources with UV excess greater than 50 per cent with respect to the single-component model for the binary component SED fits. It is interesting to note that all the seven BSSs showed the excess in UV fluxes greater than 50 per cent, and therefore we attempted the binary-component SED fits to them. Since we select stars based on a 50 per cent flux excess, it is not surprising that we find that all candidates have excess larger than this threshold. BSS 6 could not be fitted with the binary component SED since all the models fitting the data points showed the highest temperatures and hence were unreliable. The single-component SED of BSS 6 is shown in Fig. 7.

In order to fit the binary component SEDs, we used a python code BINARY SED FITTING,¹ which is based on χ_r^2 minimization technique (Jadhav et al. 2021b). We used the Koester model (Koester 2010) for this purpose, as this gives a temperature range of 5000-80 000 K and log g range of 6.5-9.5, thus making it suitable to fit a compact hotter companion. The double-component SEDs for all the BSSs are shown in Fig. 8. The top panel for each BSSs shows the fitted SED and the bottom panel shows the residual for single and composite fit in each filter. We note that the residuals come out to be nearly zero on fitting the binary component SED since the excess in UV fluxes have reduced significantly. The parameters of all the BSSs are listed in Table 2, where the parameters of the cooler components are taken from VOSA, and that of the hotter companions are taken from the python code mentioned above. The errors in the parameters are determined by following the statistical approach mentioned in Jadhav et al. (2021a). The χ^2 of the fits are found to be large even when the SED fits are visually good, mostly due to very small observational flux errors (Rebassa-Mansergas et al. 2021). We estimate modified reduced χ^2 , visual goodness of fit (vgf_b), determined by VOSA by forcing the observational errors to be at least 10 per cent of the observed flux. The vgf_b values <15 imply that the

¹https://github.com/jikrant3/Binary SED Fitting

Table 1. Coordinates of all BSSs in columns 2 and 3, UVIT *F*148*W* in column 4, GALEX NUV flux in column 5, *Gaia* DR3 Gbp, G, and Grp fluxes in columns 6–8, 2MASS J, H, and Ks fluxes in columns 9–11, and WISE W1, W2, W3, and W4 fluxes in columns 12–15. All flux values are extinction-corrected and listed in the unit of erg s⁻¹ cm⁻² Å⁻¹.

Name	RA	Dec.	UVIT.F148W±err	GALEX.NUV±err	
PS1.g±err GAIA3.Gbp±err		GAIA3.G±err	PS1.r±err	PS1.i±err	
GAIA3.Grp±err	PS1.z±err	PS1.y±err	2MASS.J±err	2MASS.H±err	
2MASS.Ks±err	WISE.W1±err	WISE.W2±err	WISE.W3±err	WISE.W4±err	
BSS 1	325.9511	65.8598	2.361e-17 ± 7.643e-21	3.631e-16 ± 2.728e-17	
$8.675e-15 \pm 4.939e-17$	$8.129e-15 \pm 3.039e-17$	$7.291e-15 \pm 1.941e-17$	$8.142e-15 \pm 1.010e-17$	$6.525e-15 \pm 2.584e-16$	
$6.370e-15 \pm 2.704e-17$	$5.725e-15 \pm 2.515e-17$	5.114 ± 1.917 e-17	$3.204e-15 \pm 7.969e-17$	$1.546e-15 \pm 4.273e-17$	
$6.225e-16 \pm 1.089e-17$	$1.262e-16 \pm 2.673e-18$	$3.672e-17 \pm 7.441e-19$	$5.706\text{e-}19 \pm 1.934\text{e-}19$	$7.371e-19 \pm 2.748e-19$	
BSS 2	325.9342	65.8757	$1.679e-17 \pm 5.301e-21$	$1.141e-16 \pm 2.201e-17$	
$3.188e-15 \pm 3.971e-18$	$2.975e-15 \pm 8.663e-18$	$2.659e-15 \pm 6.804e-18$	$3.016e-15 \pm 8.711e-18$	$2.477e-15 \pm 1.662e-17$	
$2.311e-15 \pm 8.389e-18$	$2.062e-15 \pm 5.176e-18$	$1.838e-15 \pm 1.327e-18$	$1.156e-15 \pm 2.875e-17$	$4.806e-16 \pm 1.682e-17$	
$2.076e16 \pm 7.459e18$	$4.288e-17 \pm 9.479e-19$	$1.213e-17 \pm 2.906e-19$	$3.508e19 \pm 1.460e19$	$1.424e-18 \pm 6.521e-19$	
BSS 3	326.3274	65.7773	$2.365e-17 \pm 7.672e-21$	$1.627e-16 \pm 2.349e-17$	
$3.662e-15 \pm 6.643e-18$	$3.332e-15 \pm 9.242e-18$	$2.899e-15 \pm 7.439e-18$	3.306e-15 9.287e-18	$2.644e-15 \pm 1.322e-17$	
$2.454e-15 \pm 9.100e-18$	$2.161e-15 \pm 8.379e-18$	$1.920e-15 \pm 1.180e-17$	1.128e-15 3.846e-17	$5.150e-16 \pm 1.707e-17$	
$2.123e-16 \pm 7.625e-18$	$4.304e-17 \pm 9.117e-19$	$1.238e-17 \pm 2.737e-19$	$5.034\text{e-}19 \pm 7.996\text{e-}20$	$1.205e-18 \pm 1.914e-19$	
BSS 5	326.1820	65.7889	$2.463e-17 \pm 9.270e-21$	$3.198e-16 \pm 4.833e-17$	
$4.229e-15 \pm 4.538e-18$	$3.836e-15 \pm 1.073e-17$	$3.196e-15 \pm 8.152e-18$	$3.637e-15 \pm 8.130e-18$	$2.756e-15 \pm 7.566e-18$	
$2.563e-15 \pm 9.162e-18$	$2.184e-15 \pm 3.453e-18$	$1.920e-15 \pm 4.722e-18$	$1.032e-15 \pm 3.327e-17$	$4.775e-16 \pm 1.363e-17$	
$1.988e16 \pm 6.410e18$	$3.766e-17 \pm 7.977e-19$	$1.025e-17 \pm 2.455e-19$	$4.969\text{e-}19\pm8.260\text{e-}20$	$7.833e-19 \pm 1.302e-19$	
BSS 7	326.3131	65.823	$3.050e-17 \pm 1.509e-20$	$2.531e-16 \pm 5.967e-17$	
$2.600e-15 \pm 8.428e-18$	$2.396e-15 \pm 1.831e-17$	$2.155e-15 \pm 6.738e-18$	$2.490e-15 \pm 3.566e-18$	$2.085e-15 \pm 3.831e-18$	
$1.922e-15 \pm 1.337e-17$	$1.795e-15 \pm 1.936e-18$	$1.517e-15 \pm 1.203e-17$	$8.908e-16 \pm 2.379e-17$	$4.252e-16 \pm 1.409e-17$	
$1.821e-16 \pm 6.209e-18$	$4.005e-17 \pm 8.854e-19$	$1.105e-17 \pm 2.546e-19$	$6.942\text{e-}19 \pm 1.800\text{e-}19$	$6.809e-19 \pm 1.766e-19$	
BSS 8	326.3476	65.80431	$2.387e-17 \pm 8.150e-21$	$2.745e-16 \pm 2.814e-17$	
$5.206e-15 \pm 8.971e-18$	$4.886e-15 \pm 1.423e-17$	$4.391e-15 \pm 1.122e-17$	$4.974e-15 \pm 8.664e-18$	$4.130e-15 \pm 2.978e-18$	
$3.826e\text{-}15 \pm 1.364e\text{-}17$	$3.415e-15 \pm 3.736e-18$	$3.044e-15 \pm 2.063e-17$	$1.888e-15 \pm 5.045e-17$	$8.602e-16 \pm 3.169e-17$	
$3.433e-16 \pm 1.043e-17$	$7.249e-17 \pm 1.735e-18$	$2.051e-17 \pm 4.913e-19$	$8.927e-19 \pm 1.006e-19$	$1.240e-18 \pm 1.398e-19$	



Figure 4. The PATHOS LC of BSS 1 and BSS 7, respectively.



Figure 5. The PATHOS periodograms of BSS 1 and BSS 7. For each of them, the first peak represents the orbital period, whereas the second peak represents the first harmonics.



Figure 6. Phase-folded LCs of BSS 1 and BSS 7. The LCs are folded using the period determined in the periodograms shown in Fig. 5.



Figure 7. Single-component SED fitting of BSS 6 showing UV excess. In the top panel, the blue data points show the extinction-corrected observed fluxes, whereas the black error bars represent the errors in the observed fluxes, and the orange curve represents the Kurucz stellar model fit. The bottom panel shows the residual between extinction-corrected observed fluxes and the model fluxes across the filters from UV to IR wavelengths.

fits are good (Jiménez-Esteban et al. 2018; Rebassa-Mansergas et al. 2021).

4.2 Binary modelling of the LCs

Binary modelling is an excellent technique to obtain constraints on the stellar orbital parameters of a binary system. The additional constraints help us to understand the possible formation mechanism of the BSSs. In this study, we used PHYSICS OF ECLIPSING BINARIES (PHOEBE) v0.32 (Prša & Zwitter 2005) to model the two BSSs (BSS 1 and BSS 7) systems. This software package is an open-source eclipsing binary modelling code that theoretically calculates the stellar and orbital parameters of an eclipsing binary system using a synthetic LC and a radial velocity curve. This is based on the Wilson–Devinney code (Wilson & Devinney 1971) and accepts an observed LC and radial velocity measurements as inputs, as well as initial guesses for parameters such as effective temperature, surface gravity, inclination, semimajor axis, luminosity, and surface potential.

The modelling of the two systems was performed in the following steps. Initially, we manually changed the parameters to obtain a close solution to the LC. However, due to the unavailability of the radial velocity data for both systems, we used the Q-search method to determine the mass ratios. In the Q-search method, we performed multiple modelling runs for the same system with different values of the mass ratio. The value corresponding to the lowest sum of the squared deviations is taken as the mass ratio of the system. Once the mass ratio and an initial reasonable LC are obtained, we carry out iterative minimization to find the best-fitting model. The differential correction routine, available in PHOEBE, is performed repeatedly, where the next iteration is performed using the values of the previous cycle. This process is continued till the convergence of the cost function of the fitting. The parameters of BSS 1 and BSS 7 derived by fitting the binary model to the LCs using PHOEBE are listed in Table 3.

It is noteworthy that only two BSSs showed the variability signature in *TESS*, despite the fact that many may also be binary systems with hot companions and variables found in the literature. This could be due to several reasons. A binary system becomes an eclipsing binary only if the star is able to occult the light of the other, which is possible either if the system has high inclination or the radii of the companions are large. The BSSs in this cluster are found to have compact objects like WDs as companions, which have a low probability of exhibiting eclipses. A similar system is shown in Vernekar et al. (2023). Furthermore, even if many more systems are eclipsing, their detection is limited by the amount of data available using *TESS*. Therefore, the BSSs system with longer periods cannot be probed using this method.

4.3 Properties of BSSs

On fitting the SED of the BSSs using the method described above, we found that the temperatures of the BSSs vary from 6500 to 8000 K, which is consistent with the age of the cluster. Furthermore, the luminosities vary from 8.72 to 26.91 L_{\odot} , and radii vary from 2.04 to 4.09 R_{\odot} . Upon comparing temperatures derived from fitting the SED with the temperatures based on BP/RP spectra-based from *Gaia* DR3 (Babusiaux et al. 2023), we found that the SED-based temperatures are within 400 K with the low-resolution spectra-based temperatures. We also find the rough estimation of masses of the BSSs by comparing it with the zero-age main sequence (ZAMS) of age 8.5 Gyr. The masses of the BSSs vary from 1.60 M_{\odot} to 2.15 M_{\odot} . Since the turn-off mass of the cluster is 1.55 M_{\odot} , we suggest that the BSSs are likely to have gained ~0.05 to 0.60 M_{\odot} via MT/merger channel.

We performed binary modelling of the LCs by fixing the temperatures of the BSSs generated from the SED. The binary LCs of BSS 1 and BSS 7 fitted with the model using PHOEBE are shown in Fig. 9. The observed and synthetic LCs are in agreement with each other, which is well conveyed by the residuals shown in the lower panels. This implies that the temperatures determined by constructing the SED are much accurate. In the case of BSS 1, the radius derived using SED fitting is $4.09 \pm 0.01 \text{ R}_{\odot}$ and that from binary modelling of LC is $3.55 \pm 0.12 \text{ R}_{\odot}$, which are consistent with each other. However, in the case of BSS 7, the radius of the BSSs derived by constructing the SED is $2.04 \pm 0.04 \text{ R}_{\odot}$ and that from the binary modelling of LC is $5.69 \pm 0.29 \text{ R}_{\odot}$. The parameters of both the BSSs derived from fitting the LC are listed in Table 3. Fig. 10 shows



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Figure 8. The binary-component SED of BSSs. The top panel displays the double-component SED of each BSS, with the blue data points representing the extinction-corrected flux values and the black error bars showing flux errors. The green dashed line represents the cool (A) component fit, while the brown dashed line represents the hot (B) component fit, with the pink curves representing iteration residuals. The composite fit is depicted as an orange curve, with the grey data points denoting the data points that were not included in the fits. The fractional residual for single (green) and composite (orange) fits is shown in the bottom panel. On the *x*-axis, the black error bars represent fractional errors. The cool and hot component parameters derived from the SED, as well as the estimated errors, are listed at the top of the plots.

the mesh plot of both the BSSs systems in different phases. The deformation of the blue straggler is indeed dominating the entire LC with a small effect from eclipses.

We also checked for the information on the variability of BSSs in *Gaia* DR3 (Eyer et al. 2023), where BSS 1 and BSS 7 are reported as eclipsing binaries with the period of \sim 1.07 and 0.58 d. On the other hand, BSS 3 and BSS 8 are found to be short time-scale (<0.5–1 d) MS-type oscillators (GDOR| DSCTU| SXPHE). However, for BSS 3 and BSS 8, we did not find any signatures of variability using *TESS* data.

4.4 Properties of hot companions

In order to gain a complete understanding of how the BSSs have formed, examining the nature of their hotter companions is crucial. In order to achieve this, we constructed the Hertzsprung–Russell (H-R) diagram as shown in Fig. 11. The locations of the hotter companions of BSS 1 and BSS 8 in the left panel of this figure indicate that they are ELM WDs (Althaus, Bertolami & Córsico 2013) of mass 0.18 M_☉. The stellar remnants known as ELM WDs do not ignite helium in their cores (Brown et al. 2011). The hotter companion of BSS 7, on the other hand, lies on the WD cooling curve with a mass

Table 2. The best-fitting parameters of BSSs fitted with the double-component SEDs. For each of them, whether cooler (A) or hotter (B) companion in column 2, luminosity, temperature, and radius in columns 3–5, the reduced χ_r^2 values in column 6 (the χ_r^2 values of the single fits are given in the brackets), scaling factor in column 7, number of data points used to fit the SEDs is given in column 8, and the values of vgf_b parameter in column 9 (the vgf_b values of the single fits are given in the brackets).

Name	Component	Luminosity $[L_{\odot}]$	T _{eff} [K]	Radius $[R_{\odot}]$	χ_r^2	Scaling factor	N _{fit}	vgf _b
BSS 1	А	26.97 ± 0.03	6500 ± 125	4.09 ± 0.01	25.11 (211.7)	1.51E-21	15	0.55 (1.94)
	В	$0.08^{+0.02}_{-0.01}$	15000 ± 125	0.04 ± 0.0	_		_	_
BSS 2	А	10.11 ± 0.01	6750 ± 125	2.35 ± 0.00	267.90 (3647.01)	5.03E-22	15	0.24 (1.42)
	В	0.05 ± 0.02	28000 ± 1000	0.01 ± 0.0	-		-	-
BSS 3	А	11.00 ± 0.01	7000 ± 125	2.25 ± 0.00	41.74 (111.4)	4.62E-22	15	0.15 (5.11)
	В	$0.06^{+0.02}_{-0.01}$	22000 ± 1000	0.02 ± 0.00	_		_	-
BSS 5	А	12.27 ± 0.01	7500 ± 125	2.08 ± 0.00	159.45 (439.1)	3.93E-22	15	0.25 (2.61)
	В	0.03 ± 0.01	21000 ± 1000	0.01 ± 0.00	-		_	-
BSS 7	А	8.72 ± 0.01	7000 ± 125	2.04 ± 0.07	78.92 (376.08)	3.77E-22	15	0.57 (0.94)
	В	0.1 ± 0.02	19750 ± 250	0.03 ± 0.0	-	-	-	
BSS 8	А	16.87 ± 0.02	6750 ± 125	3.01 ± 0.11	38.23 (302.1)	8.22E-22	15	0.12 (2.75)
	В	0.09 ± 0.02	14000 ± 125	0.05 ± 0.00	-		_	-

Table 3. Parameters of BSS 1 and BSS 7 obtained by fitting the binary model to the LCs using PHOEBE. For both of them, the starting time of the LC in row 1, orbital periods in row 2, the phase shift in row 3, the mass ratio in row 4, the inclination of the orbits in row 5, eccentricity in row 6, temperatures of BSSs (kept fixed) in row 7, temperatures of the hotter companions in row 8, the radius of BSSs in row 9, and radius of the hotter companions in row 10.

Parameters	BSS 1	BSS 7
HJD0 (d)	2458738.18 ± 0.015	2451738.319 ± 0.015
Period (d)	1.07066	0.5807
Phase shift	0.01456 ± 0.014	0.0116 ± 0.03
Mass ratio	0.0257 ± 0.0009	0.1142 ± 0.0057
Inclination	58.35 ± 0.05	75.62 ± 0.64
Eccentricity	0.005 ± 0.002	0.001 ± 0.002
T _{eff} (cooler)(K)	6000	7000
T _{eff} (hotter) (K)	15000 ± 4500	20000 ± 7000
Radius (cooler) (R_{\odot})	3.55 ± 0.12	5.69 ± 0.29
Radius (hotter) (R_{\odot})	0.038 ± 0.011	0.037 ± 0.013

of 0.20 M_{\odot} (Panei et al. 2007), indicating that it is an LM WD. The discovery of ELM/LM WD companions supports the Case-A/Case-B MT formation mechanism of BSS because ELM WDs and LM WDs cannot form from a single star evolution within Hubble time (Brown et al. 2011). This implies that these WDs must have lost mass during their evolution. Furthermore, the hotter companions of BSS 3 and BSS 5 are likely to be normal-mass WDs, suggesting Case-C MT as the possible formation channel. The hotter companion of BSS 2 is lying on the WD cooling curve of a high-mass WD, suggesting merger in a hierarchical triple system as the formation mechanism of this BSS. Moreover, the radii of the hotter companions of BSS 1 and BSS 7 derived from fitting the SED and fitting the LCs are in good agreement. In the case of BSS 1, the SED-based radius of the hotter companion is 0.04 \pm 0.00 R_{\odot}, whereas the LC-based radius is 0.038 \pm 0.011 $R_{\odot}.$ Similarly, in the case of BSS 7, the SEDbased radius is 0.03 \pm 0.00 R_{\odot} and that derived from LC fitting is 0.037 ± 0.013 R_{\odot}. The reason for such a large error in the parameters of the secondary companion is that the whole LC is dominated by the luminosity variation of the deformed primary. Only a small portion of the light variation arises due to eclipse. Since the parameters such as T_{eff} and radius are measured from the width and depth of the eclipses, the ability to obtain good constraints is not possible, leading to large errors.



Figure 9. The binary model fitted LCs of BSS 1 and BSS 7. The upper panels show the LC, where the black dots represent the observed LC and the red line represents the synthetic LC. The lower panels show the residual between observed and the synthetic LCs, which shows a good agreement with each other.

It is evident from the right panel of Fig. 11 that the log age of hot companions of BSS 1 and BSS 8 are \sim 8.3, implying that these BSSs are formed via MT almost 0.2 Gyr ago. Moreover, the log age of BSS 3 and BSS 7 WDs are \sim 7.6, which suggests that these objects were formed 30 Myr ago. On the other hand, the hot companions of BSS 2 and BSS 5 are lying on the WD cooling curves of log age \sim 7.9, which suggests that these BSSs were formed \sim 70 Myr ago.

This is the first-ever OC studied using *AstroSatl* UVIT data that has hotter companions of BSSs with masses varying from 0.18 to 1.0 M_{\odot}. The most comprehensive study of a BSS population in an OC is NGC 188, where 80 per cent of the BSS are found to be SB1s with typical orbital periods, $P \sim 1000$ d (Mathieu & Geller 2009; Gosnell et al. 2015). Additionally, Geller & Mathieu (2011) found that the statistical mass distribution of the companions to



Figure 10. The mesh plot showing the binary systems of BSS 1 in the left panel and of BSS 7 in the right panel shown in different phases. In each of them, the bigger object represents the BSS, whereas the smaller object represents the WD companion.



Figure 11. The left panel contains a PARSEC isochrone with an age of 4.0 Gyr and distance of 2368 pc along with a ZAMS, which is constructed by stitching ZAMS of ages 8.0, 8.5, and 9.0 Gyr. We have shown the cooler and hotter companions of NGC 7142 BSSs, along with the hot BSSs companions of other OCs. The WD cooling curves of different masses which are taken from Althaus, Bertolami & Córsico (2013) in the case of ELM WDs, Panei et al. (2007) in the case of LM WDs, and Tremblay & Bergeron (2009) in the case of normal to high-mass WDs are also shown. The right panel depicts the hot companions of BSSs lying on the WD cooling curves of different masses, indicating their approximate cooling ages.

these BSSs peaks around 0.5 M_{\odot} , strongly indicating the presence of WD secondary stars. Subsequent to that, *Hubble Space Telescope* UV photometry further confirmed that seven out of the 20 BSSs in NGC 188 possess hot WD companions with an age less than 400 Myr. This suggests recent MT events from RGB or AGB companions (Gosnell et al. 2014, 2015). Various studies based on UVIT/*AstroSat* mentioned in Section 1, including this work, also suggest MT as the predominant mechanism for the formation of BSSs in OCs. Even in the case of GCs, the absence of a correlation between cluster density and BSSs frequency implies that internal binary evolution, including MT, is the primary formation pathway, as opposed to dynamical collisions (Knigge, Leigh & Sills 2009). Furthermore, the presence of a significant population of BSs and other AFG-type post-MT binaries in the field suggests that MT is likely a prevalent formation mechanism across various environments (Carney et al. 2001; Murphy et al. 2018; Escorza et al. 2019; Panthi et al. 2023). Collectively, these populations provide valuable insights into the outcomes of post-MT events in solar-type binaries. However, detailed investigations of BSSs populations, particularly across diverse populations encompassing different cluster ages, densities, and metallicities, remain highly relevant.

5 SUMMARY AND CONCLUSIONS

The work presented in this paper can be summarized as follows:

(i) We use a machine-learning-based algorithm, ML-MOC, on *Gaia* DR3 data, to identify 546 cluster members in NGC 7142,

including 10 BSSs. We analyse these BSSs using AstroSat/UVIT data and complementary data from other archival data in a single FUV filter, F148W.

(ii) We construct the multiwavelength SED of seven out of 10 BSSs. Three BSSs had a neighbouring source within 3 arcsec, and hence their SEDs are not reliable. It is interesting to note that all the seven BSSs showed excess in UV wavelengths and, hence, likely to have the hotter companions associated with them. Out of these seven BSSs, we constructed the binary-component SED of six of them. In the case of the remaining one BSS, the hotter companion was fitting the highest possible temperature; hence, we show only the singlecomponent fit of this star. This work reports the properties of these six BSSs and their hotter companions.

(iii) An ELM WD with mass $\sim 0.18 \text{ M}_{\odot}$, temperature $\sim 15\,000 \text{ K}$, luminosity $\sim 0.08 L_{\odot}$, and radius $\sim 0.04 R_{\odot}$ is found to be the hotter companion of BSS 1. Moreover, we discover an LM WD of mass ${\sim}0.30~{\rm M}_{\odot}$ with $T_{\rm eff}\,{\sim}19\,750$ K, $L\,{\sim}\,0.10~{\rm L}_{\odot},$ and $R\,{\sim}\,0.03~{\rm R}_{\odot}$ as a companion of BSS 7. The discovery of ELM and LM WDs suggests that these BSS are formed via Case-A/Case-B MT. We construct the TESS LCs of 10 BSSs candidates of this cluster, out of which BSS 1 and BSS 7 showed the signatures of variability. The temperatures and radii of BSS 1 and BSS 7 as well as their hotter companions obtained by constructing the SED and fitting the binary models to the TESS LCs are in good agreement with each other. We confirm that BSS 1 and BSS7 are indeed eclipsing binaries using the TESS LCs along with the estimation of their orbital parameters.

(iv) The second ELM WD of mass $\sim 0.18~M_{\odot}$, temperature ${\sim}15\,000$ K, luminosity ${\sim}0.09$ $L_{\odot},$ and radius ${\sim}0.10$ R_{\odot} is found as a companion of BSS 8. Furthermore, a normal-mass WD with a mass of ~0.5 M_{\odot}, temperature of 22000 K, $L \sim 0.06 L_{\odot}$, and $R \sim 0.02 \ {
m R}_{\odot}$ is found to be the companion of BSS 3. Another BSS, BSS 5, has a normal-mass WD of mass ${\sim}0.6~M_{\odot}$ with the estimated temperature of 21 000 K, luminosity of ${\sim}0.03~L_{\odot},$ and a radius of 0.01R_o, as a hot companion. A high-mass WD of mass $\sim 1.0 \text{ M}_{\odot}$ is discovered as the hot companion of BSS 2, with an estimated temperature of 28 000 K, luminosity $\sim 0.05 L_{\odot}$, and radius \sim 0.01 R_{\odot}. These normal/high-mass WDs point towards the Case-C MT/merger as possible formation channels.

(v) There is even a possibility that BSS 6, which was not successfully fitted with a second companion, may as well be a binary as it showed UV excess. We confirm that out of 10 BSSs, at least six have WDs as the hot companions, and at least three BSSs are formed via the Case-A/Case-B MT channel. Moreover, the remaining three are formed via either Case-C MT or merger formation channels.

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DATA AVAILABILITY

The data underlying this article are publicly available at https://as trobrowse.issdc.gov.in/astro_archive/archive/Home.jsp The derived data generated in this research will be shared on reasonable request to the corresponding author.

REFERENCES

- Agarwal M., Rao K. K., Vaidya K., Bhattacharya S., 2021, MNRAS, 502, 2582
- Ahumada J., Lapasset E., 2007, A&A, 463, 789
- Althaus L. G., Bertolami M. M. M., Córsico A. H., 2013, A&A, 557, A19 Babusiaux C. et al., 2023, A&A, 674, A32
- Bayo A., Rodrigo C., Navascus Y., 2008, A&A, 492, 277
- Bouma L., Hartman J., Bhatti W., Winn J., Bakos G., 2019, ApJS, 245, 13
- Breger M. et al., 1993, A&A, 271, 482
- Bressan A., Marigo P., Girardi L., Salasnich B., Dal Cero C., Rubele S., Nanni A., 2012, MNRAS, 427, 127
- Brown J. M., Kilic M., Brown W. R., Kenyon S. J., 2011, ApJ, 730, 67
- Cardelli J. A., Clayton G. C., Mathis J. S., 1989, ApJ, 345, 245
- Carney B. W., Latham D. W., Laird J. B., Grant C. E., Morse J. A., 2001, AJ, 122.3419
- Carraro G., Ng Y. K., Portinari L., 1998, MNRAS, 296, 1045
- Castelli F., Gratton R., Kurucz R., 1997, A&A, 318, 841
- Chambers K. C. et al., 2016, preprint (arXiv:1612.05560)
- Chatterjee S., Rasio F. A., Sills A., Glebbeek E., 2013, ApJ, 777, 106
- Chen X., Han Z., 2008, MNRAS, 384, 1263
- Cohen M., Wheaton W. A., Megeath S., 2003, AJ, 126, 1090
- Cover T., Hart P., 1967, IEEE Trans. Inf. Theory, 13, 21
- Crinklaw G., Talbert F., 1991, PASP, 103, 536
- De Marchi F., De Angeli F., Piotto G., Carraro G., Davies M. B., 2006, A&A, 459, 489
- Della Croce A. et al., 2023, A&A, 674, A93
- Dempster A. P., Laird N. M., Rubin D. B., 1977, J. R. Stat. Soc. B, 39, 1
- Escorza A. et al., 2019, A&A, 626, A128
- Ever L. et al., 2023, A&A, 674, A13
- Fabrycky D., Tremaine S., 2007, ApJ, 669, 1298
- Fitzpatrick E. L., 1999, PASP, 111, 63
- Geller A. M., Mathieu R. D., 2011, Nature, 478, 356
- Gosnell N. M., Mathieu R. D., Geller A. M., Sills A., Leigh N., Knigge C., 2014, ApJ, 783, L8
- Gosnell N. M., Mathieu R. D., Geller A. M., Sills A., Leigh N., Knigge C., 2015, ApJ, 814, 163
- Hills J., Day C., 1976, Astrophys. Lett., 17, 87
- Hoag A., 1961, Mitchell, RI, Hallman, KL and Sharpless, S. Publ. of the US Naval Obs. 17
- Huang C. X. et al., 2020, Res. Notes AAS, 4, 204
- Hurley J. R., Pols O. R., Aarseth S. J., Tout C. A., 2005, MNRAS, 363, 293
- Hypki A., Giersz M., 2013, MNRAS, 429, 1221
- Indebetouw R. et al., 2005, ApJ, 619, 931
- Jadhav V. V., Pandey S., Subramaniam A., Sagar R., 2021a, JA&A, 42, 89
- Jadhav V. V., Pennock C. M., Subramaniam A., Sagar R., Nayak P. K., 2021b, MNRAS, 503, 236
- Jadhav V. V., Sindhu N., Subramaniam A., 2019, ApJ, 886, 13
- Janes K. A., Hoq S., 2011, AJ, 141, 92
- Jiménez-Esteban F., Torres S., Rebassa-Mansergas A., Skorobogatov G., Solano E., Cantero C., Rodrigo C., 2018, MNRAS, 480, 4505
- Kiseleva L., Eggleton P., Mikkola S., 1998, MNRAS, 300, 292
- Knigge C., Leigh N., Sills A., 2009, Nature, 457, 288
- Koester D., 2010, Mem. Soc. Astron. Ital., 81, 921

Kumar A. et al., 2012, in Takahashi T., Murray S. S., Herder J.-W. A. den, eds, Proc.SPIE Conf. Ser. Vol. 8443, Space Telescopes and Instrumentation 2012: Ultraviolet to Gamma Ray. SPIE, Bellingham, p. 84431N Kunimoto M. et al., 2022, ApJS, 259, 33

- Lightkurve Collaboration, 2018, Astrophysics Source Code Library, record ascl:1812.013

- Martin D. C. et al., 2005, ApJ, 619, L1
- Mathieu R. D., Geller A. M., 2009, Nature, 462, 1032
- McCrea W., 1964, MNRAS, 128, 147
- Messina S., Nardiello D., Desidera S., Baratella M., Benatti S., Biazzo K., D'Orazi V., 2022, A&A, 657, L3
- Murphy S. J., Moe M., Kurtz D. W., Bedding T. R., Shibahashi H., Boffin H. M., 2018, MNRAS, 474, 4322
- Naoz S., Fabrycky D. C., 2014, ApJ, 793, 137
- Nardiello D. et al., 2019, MNRAS, 490, 3806
- Nardiello D. et al., 2021, MNRAS, 505, 3767
- Pandey S., Subramaniam A., Jadhav V. V., 2021, MNRAS, 507, 2373
- Panei J. A., Althaus L. G., Chen X., Han Z., 2007, MNRAS, 382, 779
- Panthi A., Subramaniam A., Vaidya K., Jadhav V., Rani S., Thirupathi S., Pandey S., 2023, MNRAS, 525, 1311
- Panthi A., Vaidya K., Jadhav V., Rao K. K., Subramaniam A., Agarwal M., Pandey S., 2022, MNRAS, 516, 5318
- Peel D., McLachlan G. J., 2000, Stat. Comput., 10, 339
- Perets H. B., Fabrycky D. C., 2009, ApJ, 697, 1048
- Postma J. E., Leahy D., 2017, PASP, 129, 115002
- Postma J. E., Leahy D., 2021, JA&A, 42, 1
- Prša A., Zwitter T., 2005, ApJ, 628, 426
- Rani S., Pandey G., Subramaniam A., Rao N. K., 2023, ApJ, 945, 11
- Rebassa-Mansergas A. et al., 2021, MNRAS, 506, 5201
- Ricker G. R. et al., 2015, J. Astron. Telesc. Instrum. Syst., 1, 014003
- Rose M. B., Hintz E. G., 2007, AJ, 134, 2067
- Salaris M., Weiss A., Percival S. M., 2004, A&A, 414, 163
- Sandage A., 1953, AJ, 58, 61
- Sandquist E. L., Serio A. W., Shetrone M., 2011, AJ, 142, 194

- Sindhu N. et al., 2019, ApJ, 882, 43
- Stetson P. B., 1987, PASP, 99, 191
- Straižys V., Maskoliūnas M., Boyle R., Zdanavičius K., Zdanavičius J., Laugalys V., Kazlauskas A., 2014, MNRAS, 437, 1628
- Subramaniam A. et al., 2016a, ApJ, 833, L27
- Subramaniam A. et al., 2016b, in Herder J.-W. A. den, Takahashi T., Bautz M., eds, Proc.SPIE Conf. Ser. Vol. 9905, Space Telescopes and Instrumentation 2016: Ultraviolet to Gamma Ray. SPIE, Bellingham, p. 99051F
- Sun Q., Deliyannis C. P., Twarog B. A., Anthony-Twarog B. J., Steinhauer A., 2020, AJ, 159, 246
- Tandon S. et al., 2017, AJ, 154, 128
- Tandon S. et al., 2020, AJ, 159, 158
- Tremblay P.-E., Bergeron P., 2009, ApJ, 696, 1755
- Vaidya K., Panthi A., Agarwal M., Pandey S., Rao K. K., Jadhav V., Subramaniam A., 2022, MNRAS, 511, 2274
- Vallenari A., Brown A., Prusti T., 2022, A&A, 674, A5
- Van den Bergh S., Heeringa R., 1970, A&A, 9, 209
- Vernekar N., Subramaniam A., Jadhav V. V., Bowman D. M., 2023, MNRAS, 524, 1360
- Webbink R., 1976, ApJ, 209, 829
- Wilson R. E., Devinney E. J., 1971, ApJ, 166, 605
- Wright E. L. et al., 2010, AJ, 140, 1868

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