

HST and Optical Observations of Three Pulsating Accreting White Dwarfs in Cataclysmic Variables¹

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ABSTRACT

Ultraviolet observations using the Solar Blind Channel on the Hubble Space Telescope provide light curves and low resolution spectra of three pulsating white dwarfs in the cataclysmic variables SDSS013132.39-090122.3, SDSSJ161033.64-010223.3 and SDSSJ220553.98+115553.7. The UV light curves show enhanced pulsation amplitudes over those from simultaneous and previous optical photometry, while the UV-optical spectra are fit with white dwarf temperatures near 15,000K. These temperatures place the accreting white dwarfs outside the instability zone for non-interacting DAV white dwarfs and show that the instability strip is complex for accreting white dwarfs.

Subject headings: cataclysmic variables — stars: individual(SDSSJ013132.39-090122.3, SDSSJ161033.64-010223.3, SDSSJ220553.98+115553.7) — stars: oscillations — ultraviolet:stars

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1. Introduction

Since the discovery of the first pulsating white dwarf in an accreting close binary cataclysmic variable (GW Lib; Warner & van Zyl 1998, van Zyl et al. 2000, 2004), photometry of systems with similar orbital periods and optical spectra has revealed at least 10 more (Warner & Woudt 2004, Woudt & Warner 2004, Araujo-Betancor et al. 2005, Vanlandingham et al. 2005, Patterson et al. 2005a,b, Mukadam et al. 2006, Gänsicke et al. 2006, Nilsson et al. 2006). Seven of these pulsating accreting white dwarfs were found by followup work on cataclysmic variables (CVs) that were discovered in the Sloan Digital Sky Survey (SDSS; York et al. 2000) by Szkody et al. (2002b,2003,2004,2005,2006). The optical clue for candidate pulsating white dwarfs is the presence of broad absorption lines (originating from the white dwarf) flanking the Balmer emission lines (originating from the disk). In order to be able to view the white dwarf, the accretion disk contribution to the optical light has to be small, a situation that occurs at very low mass transfer rates, and hence short orbital periods (near 80 min). Since the SDSS provides spectra that are about 2 magnitudes fainter than previous surveys, up to 25% of SDSS discovered CVs show the white dwarf presence (Szkody et al. 2005). As the analysis of pulsation frequencies can lead to mass determination and other aspects of the internal structure of the white dwarf via the technique of asteroseismology, the compilation of a significant number of pulsating CVs provides a valuable database for probing the effects of mass transfer and accretion on the long term evolution of white dwarfs as well as the effect of external heat input, He enriched envelopes and fast rotation on the location of the instability strip.

Non-interacting hydrogen atmosphere (DA) pulsating white dwarfs (DAVs or ZZ Ceti stars) show typical non-radial g-mode pulsations with periods of 50-1400s and have temperatures in the narrow range of 10800-12,300K and $\log g \sim 8$ (Bergeron et al. 1995,2004, Koester & Allard 2000, Mukadam et al. 2004, Gianninas et al. 2005). While the pulsation frequencies and amplitudes of CV pulsators appear to be similar to those of ZZ Ceti systems (van Zyl et al. 2004, Warner & Woudt 2004, Woudt & Warner 2004, Mukadam et al. 2006), the temperatures of the underlying white dwarfs may be quite different. UV observation of GW Lib (Szkody et al. 2002a) with the Space Telescope Imaging Spectrograph (STIS) revealed a white dwarf with a temperature of 14,700K (or a best fit with 63% of the white dwarf having a temperature of 13,300K and the remainder at 17,100K), temperatures far outside the empirical instability strip for non-interacting ZZ Ceti stars. Townsley, Arras & Bildsten (2004) could match the pulsation periods and high temperature if they assumed $\ell=1$ g-modes and a $1M_{\odot}$ white dwarf with an accreted H-rich layer of $3 \times 10^{-5}M_{\odot}$. On the other hand, Araujo-Betancor et al. (2005) used a snapshot STIS spectrum to show that the CV pulsator HS2331+3905 could be fit with a 10,500K, $0.6M_{\odot}$ white dwarf, within the range of normal ZZ Ceti stars. Of course, pulsating CV white dwarfs are rotating very rapidly compared to

non-interacting ZZ Ceti stars and the thermal profiles are quite different because the surface layer is heated from above by compression and irradiation due to accretion. Although the accretion rates for the CV pulsating white dwarfs are small ($\sim 10^{-13} M_{\odot} \text{ yr}^{-1}$), the thermal profiles may still be altered and the effective temperature might not be the best parameter to locate the instability strip. This does not imply that the temperature in the driving regions is not important in locating the instability strip, but only that the effective temperature may not be reflecting the temperature in the driving regions due to accretion heating in the surface layers. Arras et al. (2006) show that the temperature in the driving zone, $\log g$ and the He abundance are all important for accreting white dwarfs.

In order to explore the physical parameter space occupied by accreting pulsators (Arras et al. 2006), and to attempt mode identification through comparison of UV to optical amplitudes (Robinson et al. 1995), we proposed further study of three additional known pulsators with STIS. After the proposal was accepted, STIS was taken out of service but we were able to switch our program to the Solar Blind Channel (SBC) which has the same wavelength coverage, albeit much lower resolution and no TIME-TAG mode of photon counting. The three objects are SDSSJ013132.39-090122.3 (Szkody et al. 2003; Warner & Woudt 2004), SDSS J161033.64-010223.3 (Szkody et al. 2002; Woudt & Warner 2004) and SDSS J220553.98+115553.7 (Szkody et al. 2003; Warner & Woudt 2004). Throughout the rest of this paper, we will refer to these objects as SDSSJ0131, SDSSJ1610 and SDSSJ2205. A summary of the orbital periods, visual brightness and optical pulsation periods is provided for these 3 systems in Table 1. The SDSS spectra are shown in Figure 1.

2. Observations and Reductions

2.1. HST Ultraviolet Observations

The HST UV data were acquired using the prism PR110L and the Solar Blind Channel on the Advanced Camera for Surveys (ACS) to provide UV spectra from approximately 1245 to 1800Å. Due to the prism, the dispersion varied from about $1.5\text{\AA} \text{ pxl}^{-1}$ at the far UV end to about $25\text{\AA} \text{ pxl}^{-1}$ at the longest wavelengths. Each system was observed for 5 HST orbits using 61s integrations in ACCUM mode. The dead time between observations was 40s so the time resolution of the spectra are 101s. The first orbit had 26 integrations (due to the initial overhead of setting up on a target) and the remaining 4 orbits contained 29 integrations. Thus, the total time on each source consisted of 142 integrations of 61s each or 8662s.

The UV data were analyzed using tools available under STSDAS and pyraf (aXe14). The primary target was extracted using a variety of extraction widths to optimize the S/N

of the resulting spectra and light curves. For the spectra, a wide extraction of 17 pixels was used and the resulting 142 spectra were co-added to produce a final spectrum. For the light curves, a smaller extraction width produced the optimum (largest) pulse signal. This optimum extraction width was determined to be 9, 12 and 4 pixels for SDSSJ0131, 1610 and 2205 respectively. The fluxes were then added across wavelength for each individual spectrum to produce photometric points throughout the 5 orbits for a light curve that could be analyzed for periodicity using Discrete Fourier transforms. For the 2 brightest systems (SDSSJ0131 and 1610), the spectra at the peaks and troughs of the resulting light curves were then individually combined into peak and trough summed spectra, as any difference in temperature between these spectra would help constrain the change in temperature during a pulsation cycle of the dominant mode.

2.2. Optical Observations

In order to prevent excessive UV light from entering the ACS detectors, STScI required continual ground observations prior to and during the HST observations to insure that the objects were not going into a dwarf nova outburst. These monitoring observations were kindly provided by a large group of worldwide amateurs and professionals (from AAVSO and ROTSE and scientists at observatories for other programs). All 3 systems remained at quiescence throughout the planning and execution of the observations.

Coordinated simultaneous optical observations were achieved for SDSSJ1610 in order to obtain the amplitude and period of the optical pulsations. The 3.5m telescope at Apache Point Observatory (APO) was used with a B filter on the CCD camera SPIcam along with the Nordic Optical Telescope (NOT) 2.56m telescope and B filter on the Faint Object Spectrograph and Camera (ALFOSC) on the night of June 30/July 01. APO used an integration time of 20s while NOT used 14.1s. Additional data were obtained on the 2m Himalayan Chandra Telescope (HCT) on the nights preceding and following the HST observation. These observations were accomplished with no filter and 30s integration times.

To check the stability of periods for SDSSJ0131, additional APO data were obtained on 2005 Dec 1 and 5 and also on 2006 Jan 30, using the imaging CCD on the Dual Imaging Spectrograph (DIS) with no filters. DIS utilizes a dichroic that splits the light at 5550Å so that the blue and red portions of the beam are then incident on two distinct CCD cameras. The light curves obtained using the blue CCD camera are used in this paper, sensitive in the wavelength range 3500-5550Å. Windowing was used to read a small portion of the CCD in

order to reduce the read out time and obtain a suitable time resolution. A standard IRAF² reduction to extract sky-subtracted light curves from the CCD frames using weighted circular aperture photometry (O’Donoghue et al. 2000) was employed. The optical light curve of the target star was divided by a sum of one or more brighter comparison stars and converted to the same fractional amplitude scale and the times converted to Barycentric Coordinated Time (Standish 1998). A Discrete Fourier Transform (DFT) was then computed for all the optical light curves up to the Nyquist frequency.

Additional optical data for SDSSJ0131 over a longer timescale were obtained from the 1.9m Radcliffe telescope at the Sutherland site of the South African Astronomical Observatory (SAAO), using the University of Cape Town CCD Photometer (UCTCCD; O’Donohue 1995) with no filter. Data were acquired during several nights in 2003 July, August and September; 2004 September and November and 2006 August. Fourier transforms were computed for each night of data, as well as combined FTs for each month.

A summary of the HST and ground-based observations is given in Table 2.

3. Light Curves and Pulsations

The summed photometry from each HST spectrum was analyzed in a similar manner as the DIS data. Figure 2 shows a light curve for a single HST orbit for each system; the pulsation is apparent in each case. The combined data from all 5 orbits were then analyzed with DFT routines to obtain the quantitative periods and amplitudes. Figure 3 shows the amplitude plots and Table 3 summarizes the resulting periods and amplitudes. The APO, NOT, and HCT data were analyzed in the same way and the results shown in Table 3 and Figure 4.

The optimum dataset exists for SDSSJ1610 since we can compare periods and amplitudes from simultaneous UV and optical data observations. Figure 2 shows that the UV pulse is most clearly resolved for this system. Table 3 shows that the observed UV and optical periods are identical, while the UV amplitudes are increased by 2-6 times over those of the optical. Our optical results are consistent in period and amplitude with the 2 independent periods that were evident in the past data of Woudt & Warner (2004) as listed in Table 1. The observed 304s period is a harmonic of the 608s period.

²IRAF (Image Reduction and Analysis Facility) is distributed by the National Optical Astronomy Observatories, which are operated by AURA, Inc., under cooperative agreement with the National Science Foundation.

For pulsating white dwarfs, the identification of the excited non-radial g-modes can provide the total mass of the star as well as the mass of the H layer (Kawaler & Hansen 1989). Each of the eigenmodes that can be excited in the star is described by a set of indices: k is the radial quantum number that gives the number of nodes between the surface and the center of the star, ℓ is the azimuthal quantum number that gives the number of nodes on the surface, and m is the number of nodes along the meridian; it is used to describe the frequency if the spherical symmetry is lost due to rotation or a magnetic field. We use spherical harmonics (Y_m^ℓ) to describe these eigenmodes. This identification of the quantum numbers k , ℓ , and m for each observed period is not easy, as it involves matching observed periods with predictions from models (with many free parameters). Robinson et al. (1995) pioneered a method of determining the ℓ value based on the change in amplitude as a function of wavelength due to limb darkening and the modified geometric cancellation (Robinson, Kepler & Nather 1982). Limb darkening effectively reduces the viewing area of the stellar disk, and this reduction in area depends on wavelength. At UV wavelengths, the increased limb darkening decreases the contribution of zones near the limb, and modes of $\ell=3$ are canceled less effectively in the UV compared to $\ell=1$ or 2. However, this does not hold true for $\ell=4$ modes, which do not show a significant change in amplitude as a function of wavelength.

Robinson et al. (1995), Nitta et al. (2000) and Kepler et al. (2000) calculated amplitudes in the UV for different values of ℓ . To attempt to constrain the ℓ value for the 608s period observed in SDSSJ1610, we determined the amplitudes for this period for four different wavelength regions: 1245-1360Å, 1360-1500Å, 1500-1640Å, and 1640-1800Å. The effective wavelengths of these regions were computed using the best-fit white dwarf model (section 4). We also computed the amplitudes for 20 individual wavelength bins but the resulting error bars were too large to be useful. Figure 5 shows the observed pulse amplitude ratios for our four wavelength regions along with the theoretical amplitude ratio for single DA white dwarf stars using Koester’s atmosphere models (Finley, Koester & Basri 1997). The top panel shows the model calculations for $\ell=1$ to 4 modes for $\log g=8$, $T_{eff}=14,500\text{K}$ (top) and $12,500\text{K}$ (bottom). While neither plot can fit all the wavelengths, the $\ell=1$ mode with $T_{eff}=12,500\text{K}$ is a good fit to all but the shortest wavelength and it appears that high order modes ($\ell=3,4$) can be ruled out. However, since the temperatures for accreting white dwarfs are quite different than for DAVs (see below), the simple non-interacting DA model may not be applicable. In addition, since the disk may cast a shadow on the equatorial regions of the white dwarf, the limb darkening and geometric cancellation may be different than that for single white dwarfs.

It is also interesting that the peaks in the FT of SDSSJ1610 show no splitting (Figure 4), even though the resolution is $\sim 10^{-5}$ Hz ($v\sin i \sim 0.5$ km s $^{-1}$), whereas fitting of line widths

in HST spectra have shown velocities about 50-100 km s⁻¹ (Sion et al. 1994) for accreting white dwarfs. However, there may be some process that is suppressing some of the modes. Long term monitoring of GW Lib (van Zyl et al. 2004) showed clusters of periods separated only by about 1 μ Hz whose origin is not clear. It will take a much longer combined run on SDSSJ1610 to tell if this is also the case in this system.

For the other two objects, where simultaneous optical light curves could not be obtained, we can only compare the observed UV periods and amplitudes with available optical values. For SDSSJ2205, the single observed UV period of 576s matches one of the previously observed optical periods noted by Warner & Woudt (2004), and the UV amplitude is about 6 times the optical value. The strongest period (330s) observed in the optical by Warner & Woudt is not visible in the UV.

The brightest system (SDSSJ0131) shows no UV pulsation at the highest amplitude period (595s) previously evident in the optical data published in Warner & Woudt (2004). Their other 2 optical periods (260s and 335s) are also not apparent in the SBC data. Instead, a shorter period of 213 s near the Nyquist frequency is evident in the HST data with an amplitude near 80 mma. The 213s period is obvious, although just barely resolved, in the SBC light curve (Figure 2). This is the shortest period yet seen in an accreting pulsator. Our extended optical coverage on this object with APO in the months following the HST observation and with SAAO in the years preceding and following the HST (Table 2) provides a handle on the stability of the observed periods. Figure 6 shows the DFT of the 2006 January data while Figure 7 shows all the monthly combined runs from the SAAO (Table 4 lists the observed frequencies and periods from the SAAO runs). It is clear that the pulsation spectra are highly variable, with longer periods most prominent in the runs of 2003 (Warner & Woudt 2004 use data from 2003 September), while the 2004 runs show only the short period seen in the 2005 HST data. Then, in 2006 January (Figure 6) both the long period (near 600s) and the short (213s) period are evident, and in 2006 August (Figure 7), these are both still present, along with an even longer period (1130s) that is similar to one seen in 2003 Aug. While changes in pulsation spectra are common in cool non-interacting ZZ Ceti stars on timescales of months and even days, it is not obvious what is to be expected with these hotter pulsators. Long term monitoring of GW Lib (van Zyl et al. 2004) does show that its pulsation amplitudes are unstable as well. The 2 periods near 650 and 370s remain visible in all runs (including the HST data), while other periods appear and become invisible in other runs. In SDSSJ0131, it appears that all periods show large amplitude variations. Comparing the amplitude of the 213s period in the UV with that seen in the optical gives a ratio of about 6, similar to that of SDSSJ1610 and 2205.

4. Spectral Fits

The UV spectra of low mass transfer rate CVs are a combination of the underlying white dwarf, the accretion disk and a source of emission lines (likely an accretion disk chromosphere). To model the observed SBC spectra, we employed a procedure similar to our analysis of STIS data on similar systems (e.g. Gänsicke et al. 2005). This involved using Hubeny white dwarf LTE models TLUSTY195 and SYNSPEC46 (Hubeny & Lanz 1995) in a grid of temperatures, gravities and abundances to create a best fit to the spectrum. Due to the low resolution of the SBC and the past results which were all consistent with $\log g=8$, we fixed the $\log g$ at this value and used 0.01 solar abundance to fit the best temperature white dwarf. The constraints on the fit include that the white dwarf must provide the best match to the UV continuum shape, especially the broad Ly α line, and its optical flux must fall below the observed optical fluxes. Figure 8 shows the model results for a range of white dwarf temperatures (10,000K to 20,000K) compared to the observed HST and SDSS spectra of the brightest system SDSSJ0131. It is immediately clear that the temperature typical of non-interacting ZZ Cet pulsators (~ 12000 K) is too cool to fit the UV spectrum. A temperature of 14,500K is derived from the best fit. This temperature white dwarf extrapolated to optical wavelengths produces a g magnitude of 19.2, well below the observed SDSS g mag of 18.3. In addition to the white dwarf, there is some continuum source (the disk) that produces about 30% of the light across the spectrum. This is especially evident in the core of Ly α which should be zero for a high gravity white dwarf. As there is no good model of a non-steady state disk at these low mass transfer rates, we added this continuum component as a simple black body (power laws were tried with similar results), used simple Gaussians to fill in the emission lines and redid the fits with the two components. The resulting best fit for SDSSJ0131 is shown in Figure 9 and the results summarized in Table 5. A similar procedure produced the best fits for SDSSJ1610 and SDSSJ2205 (Figures 10 and 11 and Table 5).

The spectra that produced the largest amplitude pulse were also added and fit with this procedure. These spectra had higher S/N but less overall flux. Thus, the shape of the spectrum may be better defined but the total fluxes are not correct for distances, etc. We also list these temperature fits as T_{optext} in Table 5 as a measure of the error of our temperature fits.

We also attempted to fit the difference between the peaks of the pulsations and the troughs for the two brightest sources SDSSJ0131 and SDSSJ1610. The spectra of the peak points in the light curves were summed to produce one total peak spectrum and then the troughs were summed to produce a total trough spectrum. The resulting fits to the peak and trough spectra showed no change in the white dwarf temperatures in either case, which

is not surprising given that the uncertainties of our fits are on the order of 1000K.

The temperatures of the white dwarf in all 3 systems are similar, close to 15,000K and much hotter than the temperatures of the DAV pulsators. These temperatures are also very similar to that derived for GW Lib (Szkody et al. 2002). Because we had used a two-temperature white dwarf fit for GW Lib, rather than a white dwarf + BB, we redid the fits to the STIS data of GW Lib using the approach here. The result produces a best fit with $T_{wd}=15,400\text{K}$ and $T_{BB}=12,000\text{K}$ but this fit is not as good as our past fit with 2 white dwarf components of 13,300K and 17,100 covering 63% and 37% respectively of the white dwarf surface. Further confirmation that the two temperature white dwarf is correct for GW Lib comes from inspection of its STIS spectrum. This shows that the core of Ly α does go down to zero flux, indicating that the white dwarf, rather than a second component such as the disk, contributes all the UV flux. A comparison of the single temperature white dwarf fit (14,700K) with that of the WD+BB also shows that adding in a second component only changes the derived temperature of the white dwarf by less than 1000K.

Thus, we now have secure results that four of the accreting ZZ Ceti stars in CVs are hotter than single pulsators. On the other hand, HS2331+3905 is also secure in having a white dwarf temperature of 10,500K (see Fig. 13 in Araujo-Betancor et al. 2005). While the UV emission lines in HS2331+3905 are much stronger than in the four objects that have hot white dwarfs (likely partly due to the weak white dwarf continuum), and there are continuum variations which might imply a higher inclination, the fact remains that the white dwarf pulsations are visible in HS2331+3905 so the white dwarf light is prominent and it is cool. In contrast to our three SDSS sources, all of which have UV fluxes greater than the optical as in Figure 8, HS2331 has larger optical than UV fluxes, and a distinct rise in flux longward of 1600Å which pins down the white dwarf temperature to its cool value.

Since accreting white dwarfs are known to show UV absorption lines of Si, C and other metals, it is clear they do not have pure H atmospheres (Sion 1999). Thus, perhaps it is not unexpected that the instability strip is not the same as for ZZ Ceti stars. For accreting model white dwarfs with a high He abundance (>0.38), Arras, Townsley & Bildsten (2006) find an additional hotter instability strip at $\sim 15,000\text{K}$ for low mass white dwarfs due to He II ionization. Thus, they infer that HS2331+3905 has a low mass white dwarf while GW Lib (and hence SDSSJ0131, SDSSJ1610 and SDSSJ2205) should have highly evolved donor stars and more massive white dwarfs.

5. Conclusions

Our low resolution light curves throughout five HST orbits for our three white dwarf accreting pulsators all have increased pulsation amplitudes in the UV compared to the optical, consistent with $\ell=1$ modes. The simultaneous optical/UV data for SDSSJ1610 showed identical periods in both wavelength regions, while SDSSJ0131 and SDSSJ2205 showed one of three periods evident in the optical at other times. The long term monitoring of SDS0131 shows that all pulsation periods come and go, indicating a high level of instability for this system.

The summed UV spectra reveal hot white dwarfs in all three systems, similar to past results on GW Lib but contrary to the lone cool accreting white dwarf pulsator HS2331+3905. Based on just these five systems with well-determined white dwarf temperatures from the UV, it appears that there is a wide range in the instability strip for accreting systems, with most accreting pulsators in a regime that is hotter than for non-accreting ZZ Ceti stars. Whether this is due to a difference in the white dwarf mass and/or He composition provided by an evolved donor as proposed by Arras, Townsley & Bildsten 2006 can be confirmed if the white dwarf masses and/or the composition of the donor can be determined. As there is no evidence of a donor star out to 9000\AA in the SDSS spectra and the systems are generally faint, the donor will be difficult to characterize. Since the absorption lines of the white dwarf are difficult to measure in the optical with the contaminating Balmer emission, the best approach may be the use of high resolution UV (the Cosmic Origins Spectrograph) time-resolved data in the future to obtain the mass of the white dwarf, although the unknown inclinations will create some uncertainty.

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Table 1. Source Characteristics

SDSS J	$g(\text{mag})$	$P_{orb}(\text{min})$	Opt Pulse P (sec)	Amp (mma)	Ref
0131	18	98	260, 335, 595	4, 9, 17	Warner & Woudt 2004
1610	19	81	345, 607 ^a	7, 27	Woudt & Warner 2004
2205	20	–	330, 475, 575	10, 7, 8	Warner & Woudt 2004

^aSeveral combinations of these 2 independent periods are also evident.

Table 2. Summary of Observations

SDSSJ	Obs	Instr	Filter	Int(s)	Time
0131	HST	SBC	PR110L	61	2005 Jun 18 1:35:49 - 8:42:35
0131	APO	DIS	none	15, 25	2005 Dec 01 07:54:56 - 08:42:22
0131	APO	DIS	none	15, 30	2005 Dec 05 07:25:21 - 08:19:53
0131	APO	DIS	none	15	2006 Jan 30 01:41:23 - 03:12:27
0131	SAAO	UCTCCD	none	30	2003 July 2 (2.18h), 3 (2.23h), 5 (2.17h)
0131	SAAO	UCTCCD	none	120, 100	2003 July 25 (3.34h), 28 (3.47h)
0131	SAAO	UCTCCD	none	20, 60	2003 Aug 20 (2.32h), 22 (1.53h)
0131	SAAO	UCTCCD	none	10-90	2003 Sep 17 (3.37h), 20 (1.47h), 21 (4.65h), 22 (6.16h), 24 (1.28h), 25 (2.20h)
0131	SAAO	UCTCCD	none	60	2004 Sep 14 (2.88h), 15 (3.48h), 16 (5.10h) 19 (3.13h), 20 (5.11h), 21 (0.83h)
0131	SAAO	UCTCCD	none	30-60	2004 Nov 3 (5.43h), 4 (1.53h), 5 (3.21h), 6 (4.57h)
0131	SAAO	UCTCCD	none	20	2006 Aug 25 (4.03h), 26 (4.78h), 27 (5.05h)
1610	HCT	CCD	none	30	2005 Jun 30 15:21:21 - 16:22:56
1610	HST	SBC	PR110L	61	2005 Jun 30 23:13:19 - 2005 Jul 01 6:18:13
1610	NOT	ALFOSC	B	14.1	2005 Jun 30 22:16:59 - 2005 Jul 01 03:29:39
1610	APO	SPIcam	B	20	2005 Jul 01 03:33:33 - 07:58:28
1610	HCT	CCD	none	30	2005 Jul 01 17:01:53 - 17:56:34
2205	HST	SBC	PR110L	61	2005 May 23 16:05:03 - 23:07:33

Table 3. Summary of Observed Periods (HST, APO, NOT, HCT)

SDSSJ	Wavelength (Å)	Period (s)	Amp (mma)
0131	1255-1800	213.72±0.10	78.0±8.8
0131	optical	581.3±3.8; 211.8±1.2; 79.3±0.2	31.3±3.3; 13.6±3.3; 9.5±3.3
1610	1245-1800	608.22±0.30; 304.10±0.29 ^a ; 220.81±0.31	186.1±7.5; 48.3±7.5; 23.4±7.4
1610	optical	608.20±0.28; 304.16±0.20 ^a ; 220.68±0.09	30.6±2.2; 10.5±2.2; 12.2±2.2
2205	1202-1800	576.2±1.6	46±11

^a304s period is a harmonic of the 608s period.

Table 4. Summary of SAAO Periods for SDSSJ0131

Run	Freq (μHz)	Period (s)	Amp (mma)	Comment
2003 July 2-5	1664.1 ± 0.2	600.9	20.9 ± 2.1	f1 ^a
	2989.5 ± 0.3	334.5	12.1 ± 2.2	f2
	4629.5 ± 0.4	216.0	8.0 ± 2.2	f3 ^a
2003 July 25-28	1690.0 ± 0.2	591.7	20.6 ± 3.0	f1
	1603.1 ± 0.4	628.8	10.4 ± 3.0	
	2970.9 ± 0.4	336.6	9.2 ± 3.0	f2
2003 Aug 20-22	861.7 ± 0.4	1160.5	11.5 ± 2.3	f5
	1680.8 ± 0.2	595.0	24.6 ± 2.4	f1
	2946.2 ± 0.4	339.4	12.2 ± 2.3	f2
	3777.8 ± 0.5	264.7	9.9 ± 2.4	f4
2003 Sep 17-25	1687.5 ± 0.1	592.6	13.1 ± 1.1	f1
	2965.5 ± 0.1	337.2	8.9 ± 1.1	f2
	3786.2 ± 0.2	264.1	4.1 ± 1.1	f4
	7393.0 ± 0.2	135.3	3.8 ± 1.1	
2004 Sep 14-21	4671.7 ± 0.1	214.1	10.8 ± 1.8	f3
2004 Nov 3-6	4685.3 ± 0.1	213.4	12.7 ± 1.3	f3
2006 Aug 25-27	885.2 ± 0.2	1129.6	16.2 ± 1.1	f5?
	1758.9 ± 0.2	568.5	14.7 ± 1.1	f1?
	4669.2 ± 0.3	214.2	10.8 ± 1.1	f3

^af3 may be f1+f2?; f1 may be 2f5?

Table 5. White Dwarf Fits

Parameter	SDSSJ0131	SDSSJ1610	SDSSJ2205
T_{wd} (K)	14500	14500	15000
T_{BB} (K)	15000	12000	11000
distance (pc)	360	480	810
g_{wd} (model)	19.3	19.9	21.0
g_{wd} (sdss)	18.3	19.1	20.1
T_{wd} (optext)	14000	14000	14000
T_{BB} (optext)	16000	16000	14000

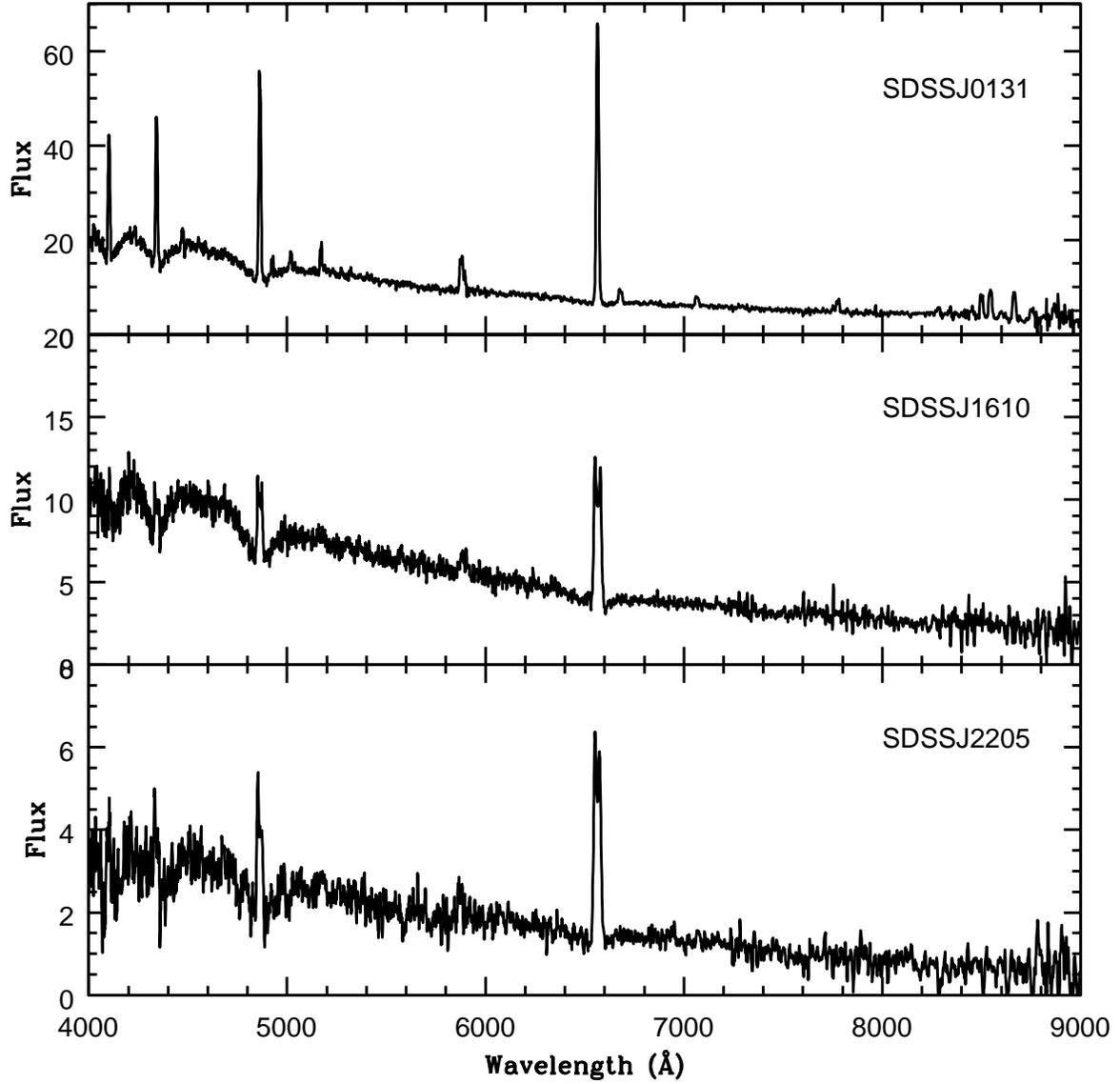


Fig. 1.— SDSS spectra of our 3 sources, showing the absorption troughs surrounding the Balmer emission lines. Units of flux are 10^{-17} ergs cm^{-2} s^{-1} \AA^{-1} .

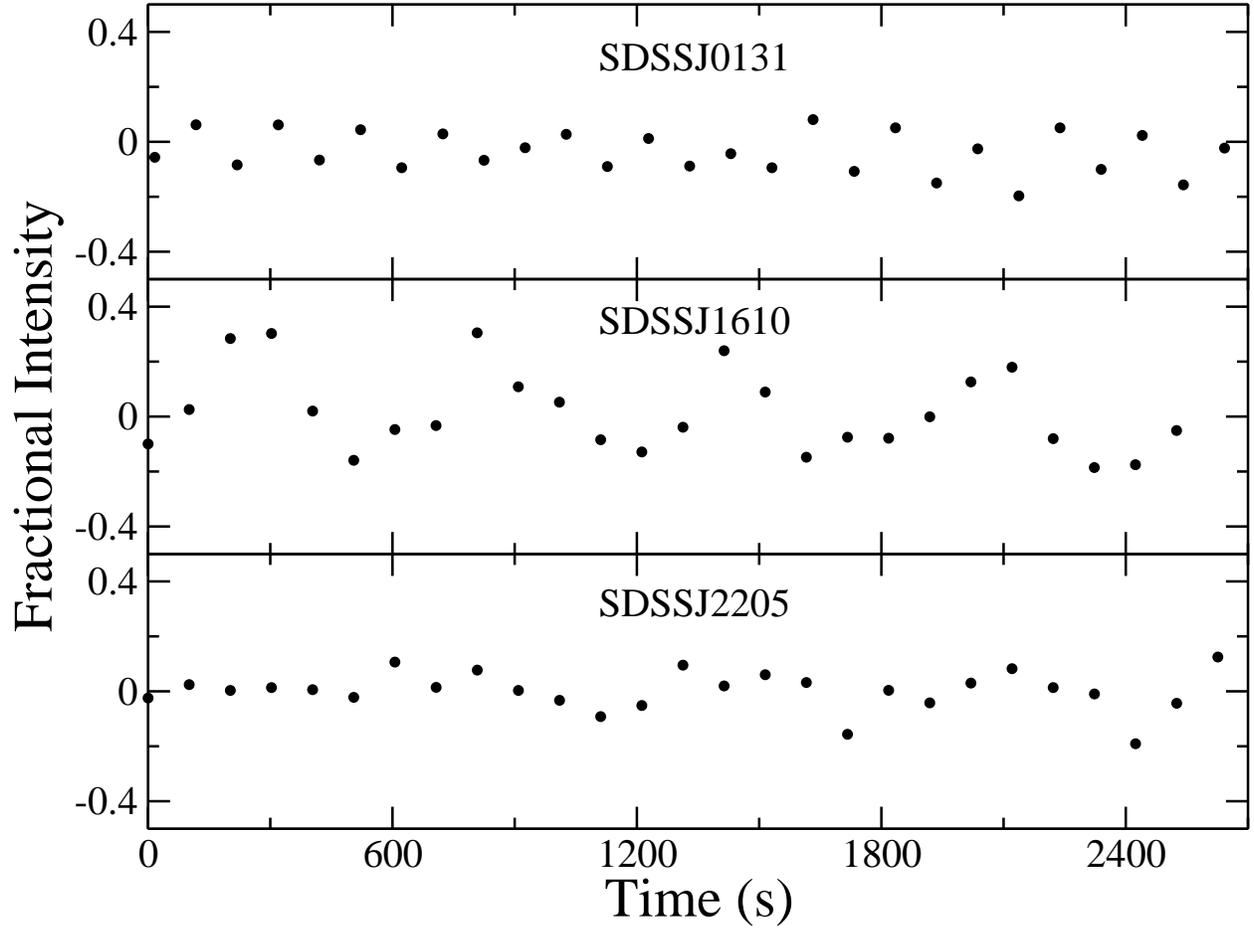


Fig. 2.— Single orbit SBC light curves from the extracted spectra for each of our 3 systems. Each point is one 61s integration.

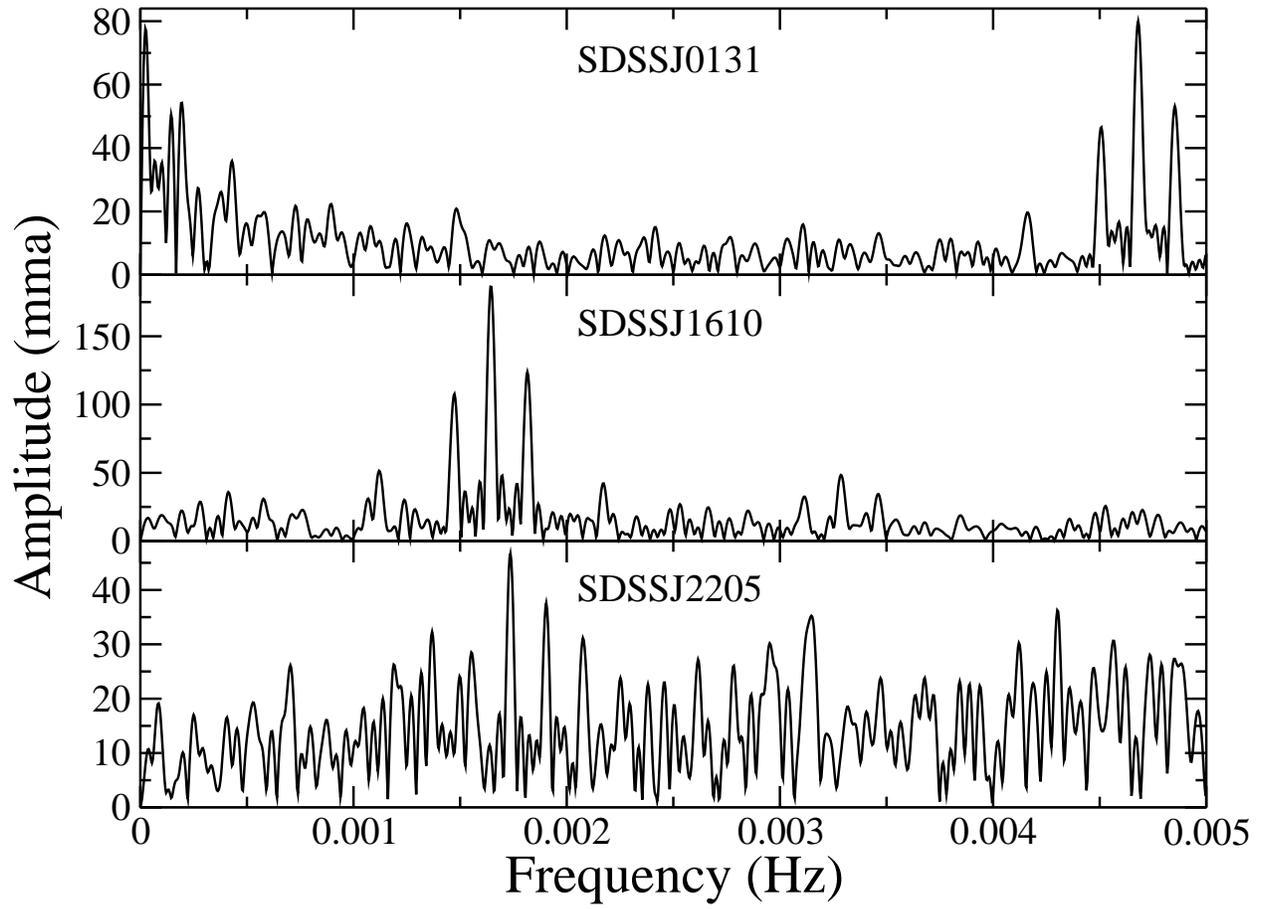


Fig. 3.— DFTs of the combined light curves for the 5 HST orbits of each system.

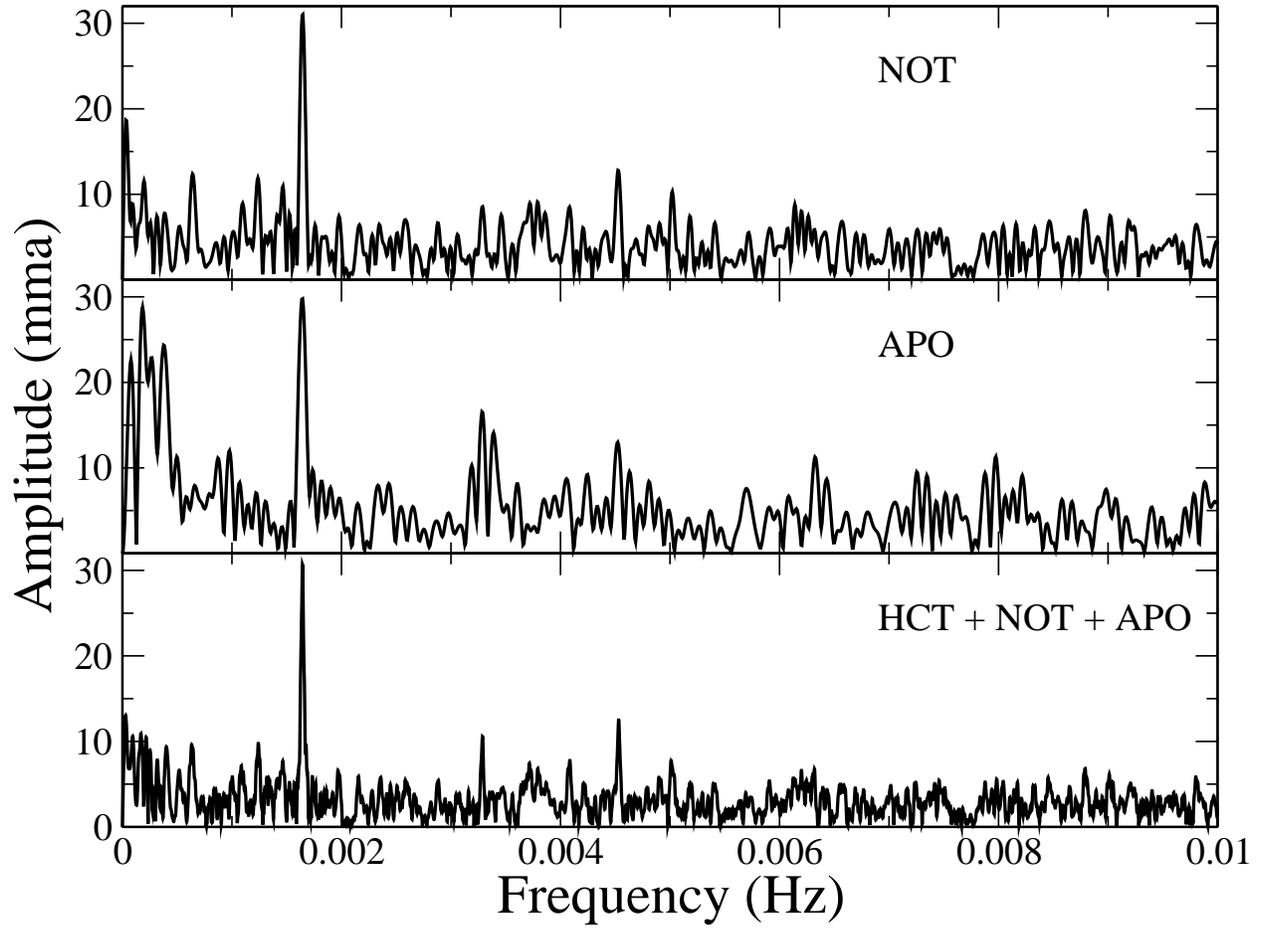


Fig. 4.— DFTs of the APO and NOT data obtained simultaneously with the HST observation of SDSSJ1610 and the combined DFT of the APO, NOT and HCT data.

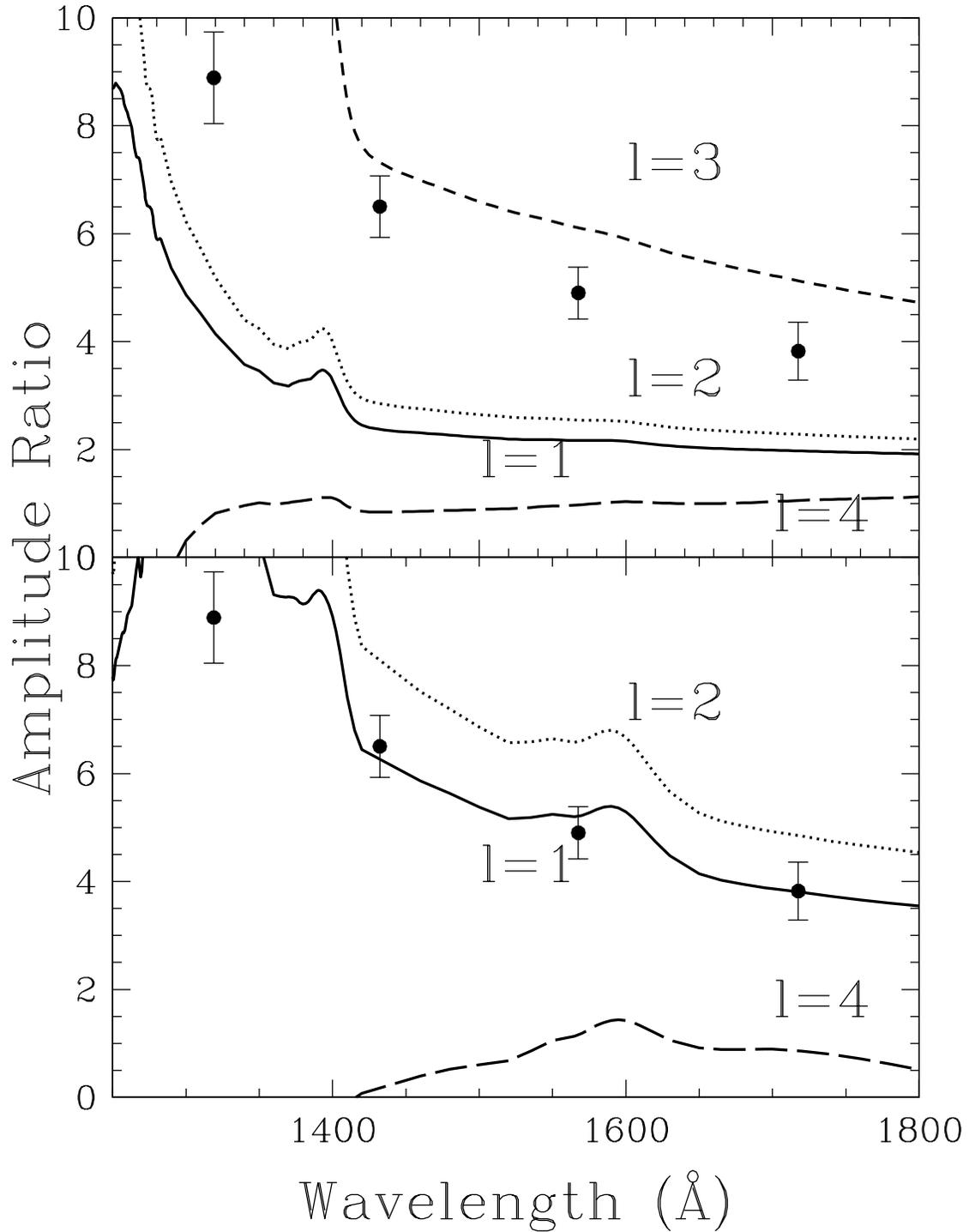


Fig. 5.— Amplitude ratio expected for different modes of the 608s pulsation for DA white dwarfs of $\log g=8$ and 2 different T_{eff} : top, 14,500K and bottom, 12500K. Observed ratios at 4 different UV wavelengths sampled for SDSSJ1610 are superposed.

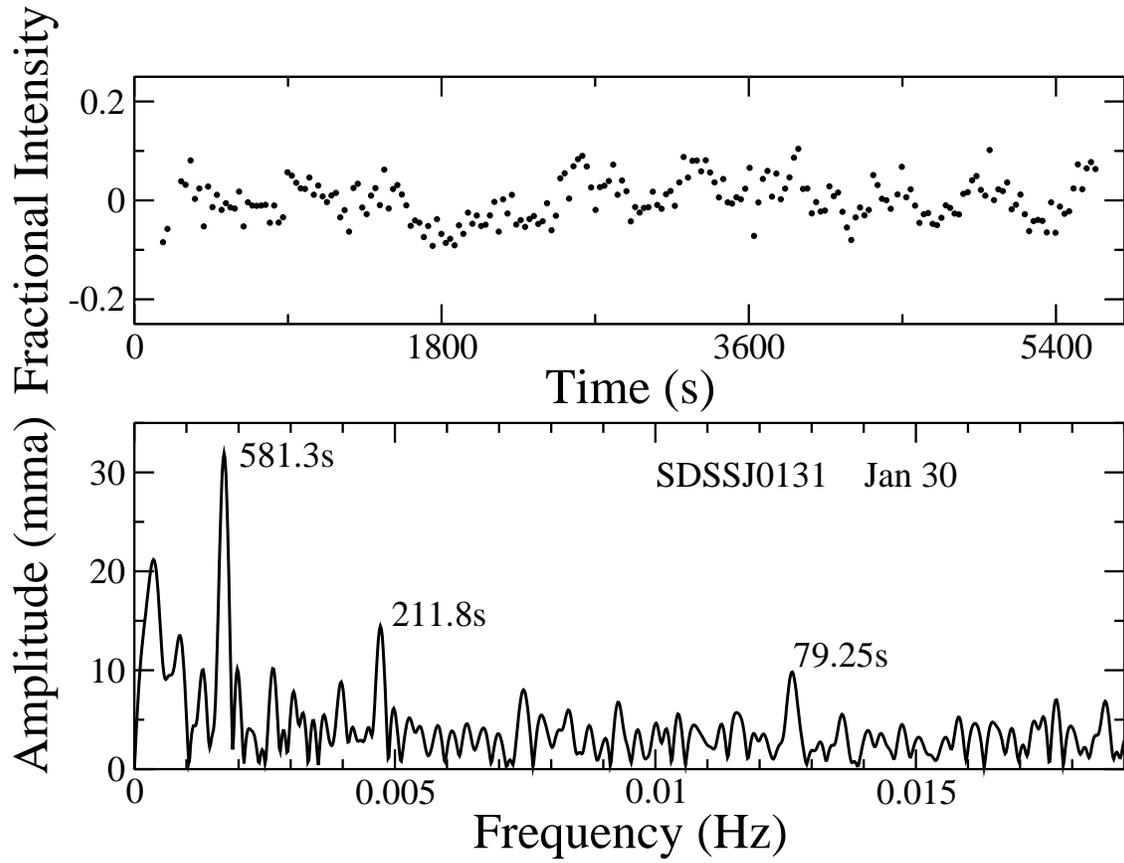


Fig. 6.— Light curve and DFT for 2006 Jan 30 APO data on SDSSJ0131.

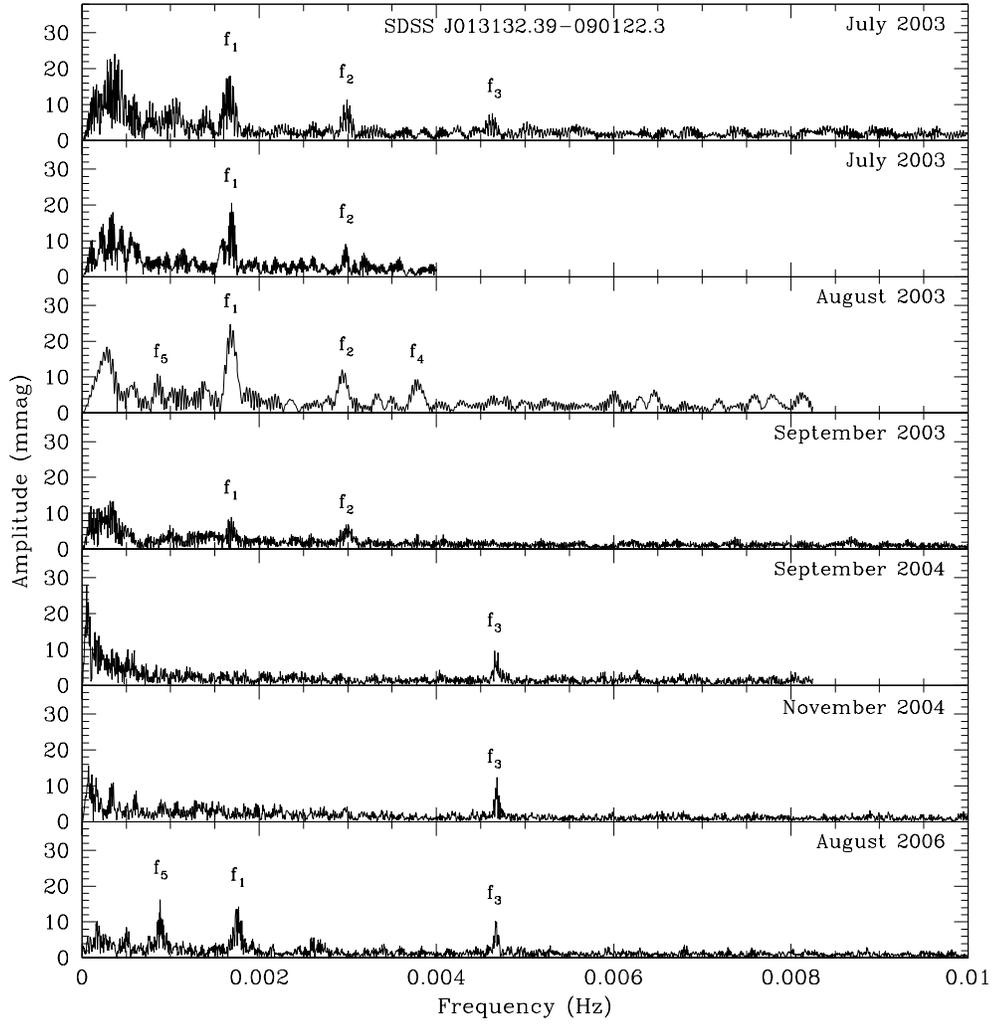


Fig. 7.— FTs for all SAAO data on SDSSJ0131.

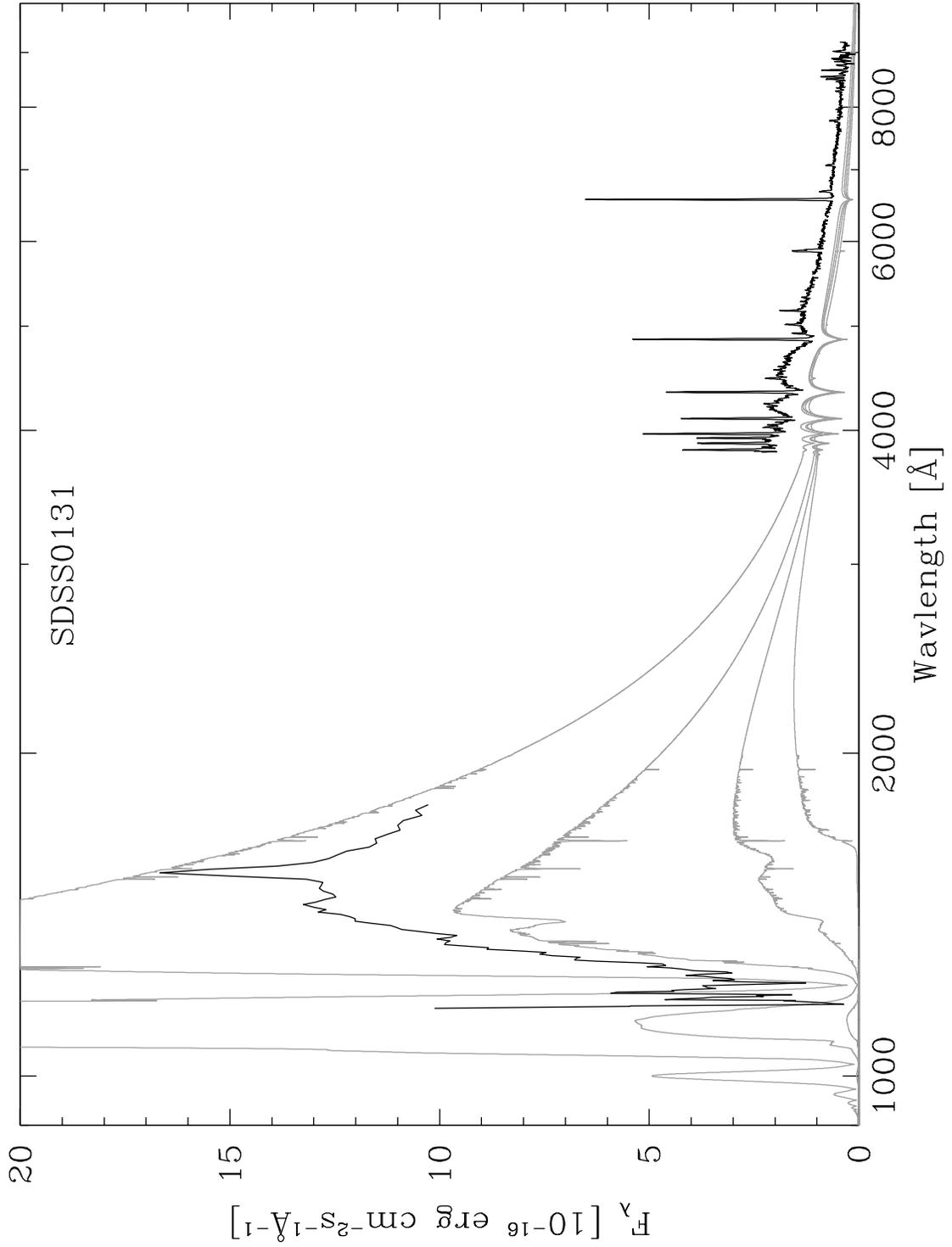


Fig. 8.— SBC and SDSS spectra of SDSSJ0131 (dark lines) compared with model white dwarfs from 10,000K, 12,000K, 14,000K and 20,000K (light grey lines from bottom to top). The best fit was a 14,000K white dwarf.

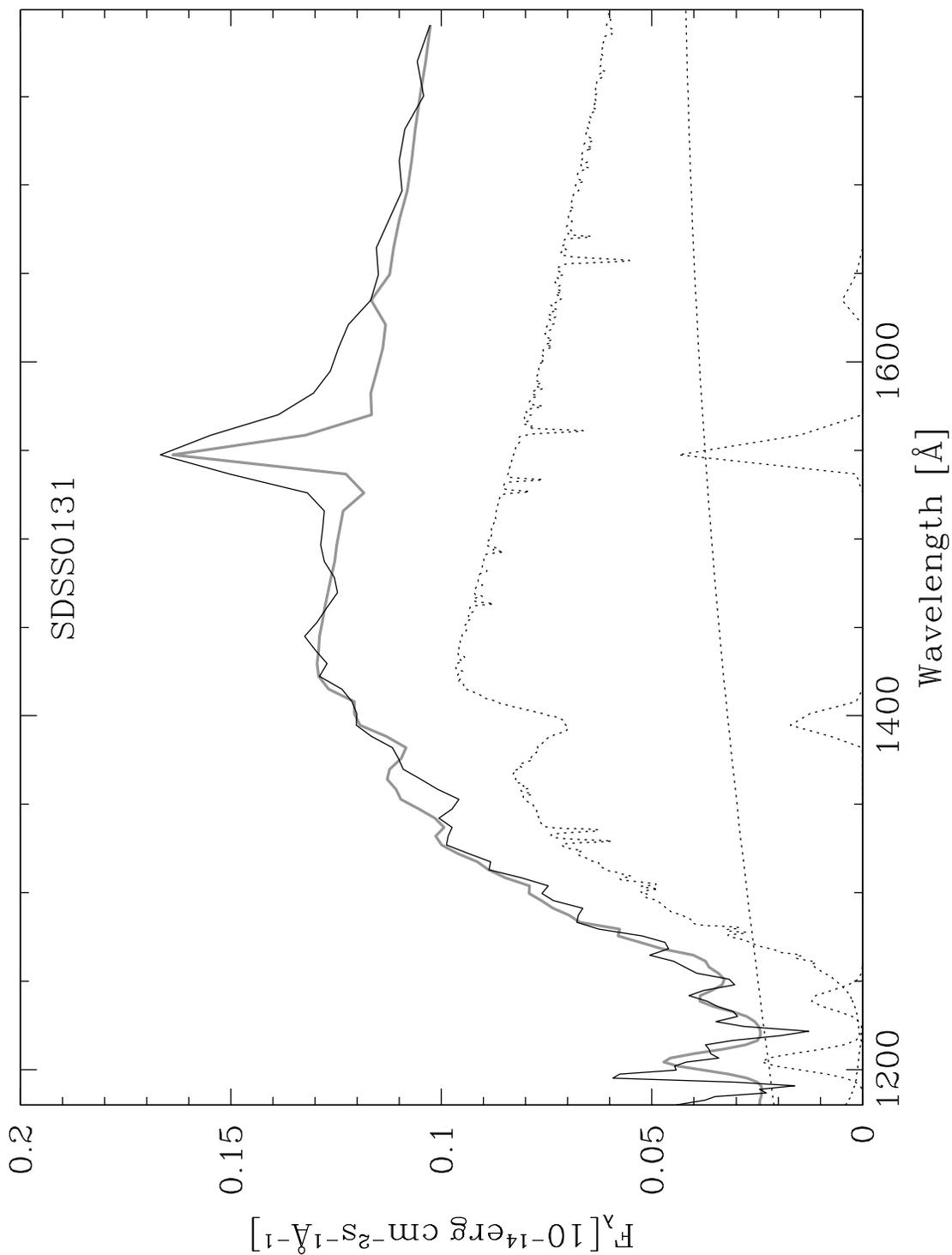


Fig. 9.— Best fit of SBC spectrum of SDSSJ0131 with a 14,400K white dwarf at 350 pc and 15,000K black body component.

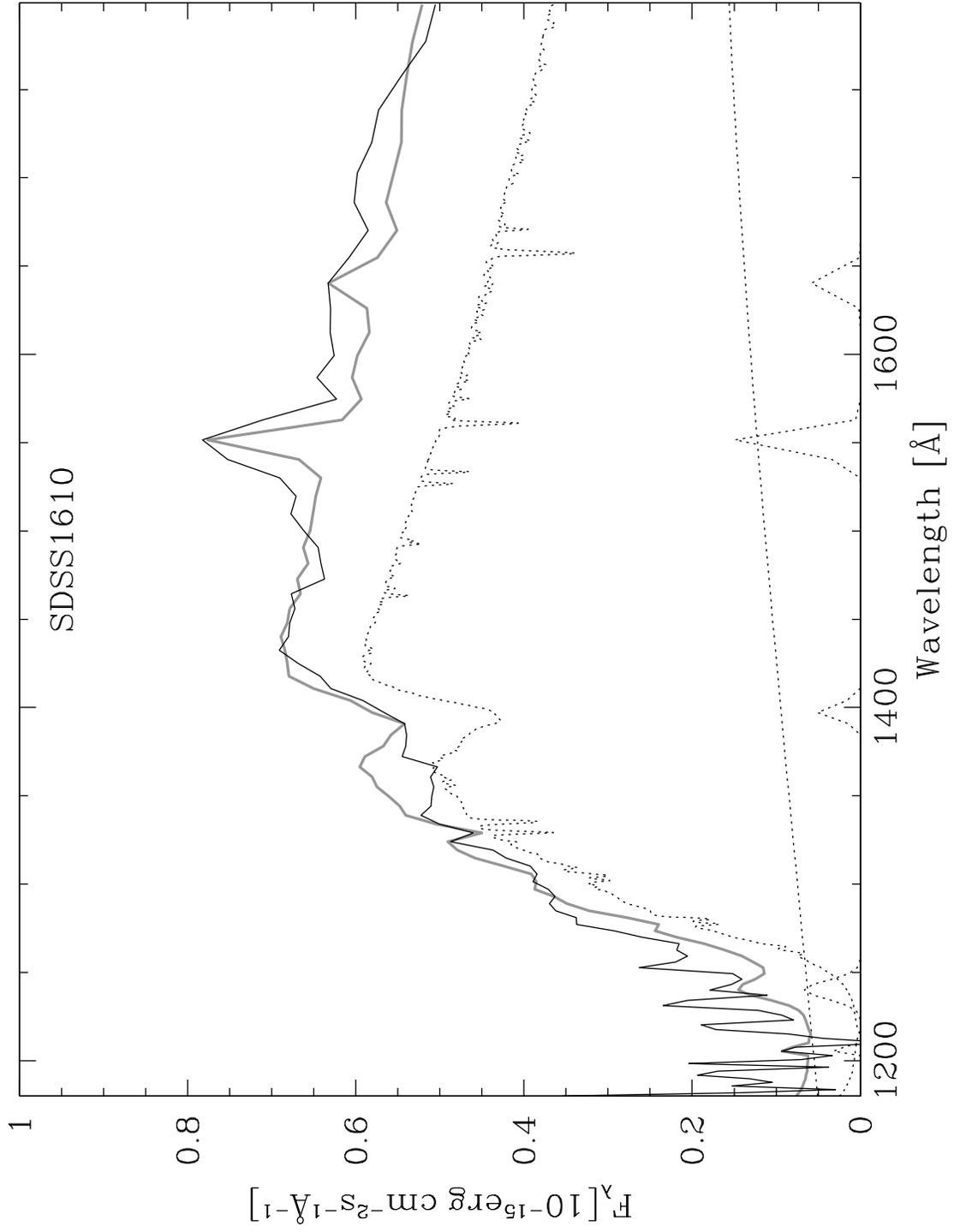


Fig. 10.— Best fit of SDSSJ1610 with a 14,500K white dwarf at a distance of 480 pc and a 12000K black body.

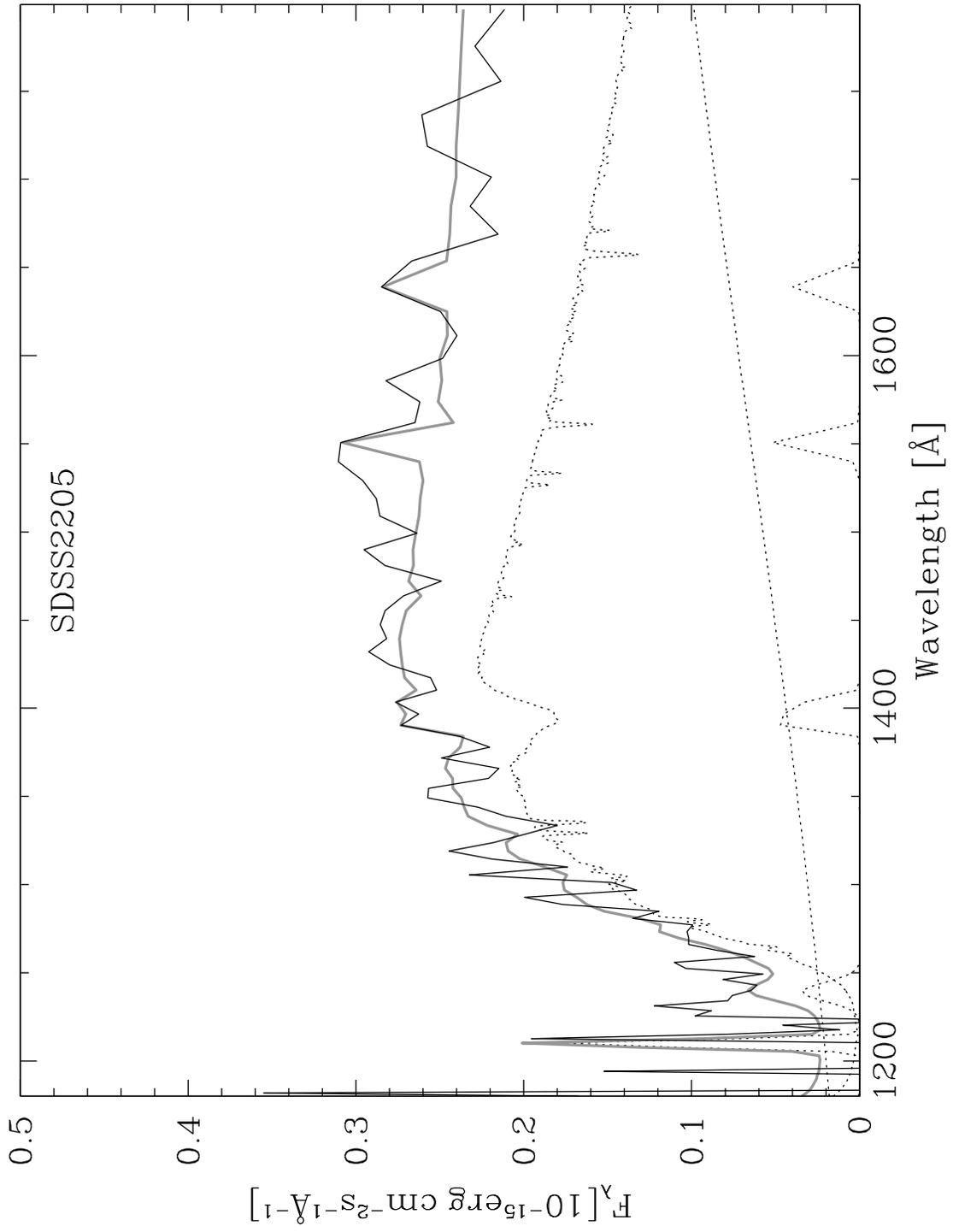


Fig. 11.— Best fit of SDSSJ2205 with a 15,000K white dwarf at a distance of 810 pc and a 11,000K black body.