

## ANALYSIS OF A FRAGMENTING SUNSPOT USING *Hinode* OBSERVATIONS

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### ABSTRACT

We employ high-resolution filtergrams and polarimetric measurements from *Hinode* to follow the evolution of a sunspot for eight days starting on 2007 June 28. The imaging data were corrected for intensity gradients, projection effects, and instrumental stray light prior to the analysis. The observations show the formation of a light bridge at one corner of the sunspot by a slow intrusion of neighboring penumbral filaments. This divided the umbra into two individual umbral cores. During the light bridge formation, there was a steep increase in its intensity from 0.28 to 0.7  $I_{QS}$  in nearly 4 hr, followed by a gradual increase to quiet-Sun (QS) values in 13 hr. This increase in intensity was accompanied by a large reduction in the field strength from 1800 G to 300 G. The smaller umbral core gradually broke away from the parent sunspot nearly two days after the formation of the light bridge, rendering the parent spot without a penumbra at the location of fragmentation. The penumbra in the fragment disappeared first within 34 hr, followed by the fragment whose area decayed exponentially with a time constant of 22 hr. In comparison, the parent sunspot area followed a linear decay rate of 0.94  $\text{Mm}^2 \text{hr}^{-1}$ . The depleted penumbra in the parent sunspot regenerated when the inclination of the magnetic field at the penumbra–QS boundary became within 40° from being completely horizontal and this occurred near the end of the fragment’s lifetime. After the disappearance of the fragment, another light bridge formed in the parent which had similar properties as the fragmenting one, but did not divide the sunspot. The significant weakening in field strength in the light bridge along with the presence of granulation is suggestive of strong convection in the sunspot, which might have triggered the expulsion and fragmentation of the smaller spot. Although the presence of QS photospheric conditions in sunspot umbrae could be a necessary condition for fragmentation, it is not a sufficient one.

*Key words:* Sun: photosphere – sunspots – Sun: surface magnetism – techniques: photometric – techniques: polarimetric

*Online-only material:* color figures

### 1. INTRODUCTION

The solar magnetic field is distributed on a wide range of spatial scales. Sunspots can be regarded as the largest magnetic structures with diameters of 20–40 Mm and field strengths in excess of 2.5 kG at the photosphere (see Solanki 2003 and Borrero & Ichimoto 2011 for reviews on the properties of sunspots). The average lifetime of a sunspot is of the order of several days (Ringnes 1964) and its equilibrium configuration is determined by the balance of forces due to gravity, the magnetic field, and the gas pressure (Priest 1982).

The formation of sunspots is initiated by the coalescence of smaller magnetic elements which rise from the convection zone to the surface due to buoyancy. The systematic merging of these fragments often results in a pore, which is characterized by strong and relatively vertical fields ranging from 1 to 1.5 kG. Partial penumbral formation occurs when the magnetic flux exceeds  $1\text{--}1.5 \times 10^{20}$  Mx (Leka & Skumanich 1998). The penumbra develops very rapidly, with pieces of it being completed within an hour or less (Bumba 1965; Keppens & Martínez Pillet 1996). A newly formed penumbral segment is practically indistinguishable from a more mature one in terms of field strengths, inclination angles, and continuum intensities (Leka & Skumanich 1998). Recent high-resolution observations

by Schlichenmaier et al. (2010b) illustrate the rapid sector-wise formation of the penumbra, which fills half the umbral circumference in about 4 hr. The development of a rudimentary penumbra initiates the Evershed flow (Evershed 1909), a nearly horizontal outflow of plasma that starts in the inner penumbra and returns to the solar surface in the mid-penumbra and beyond. During the formation of sunspots, the umbra is often divided by one or several light bridges which tend to demarcate individual magnetic regions (Bumba 1965). Light bridges can also be seen during sunspot fragmentation (Bumba 1965; García de La Rosa 1987). They are considered “field-free” intrusions in umbrae (Parker 1979; Choudhuri 1986) or manifestations of magnetoconvection (Rimmele 1997; Hinzberger et al. 2002).

According to McIntosh (1981), the decay of sunspots starts almost as soon as they are formed. Observations by Wallenhorst & Topka (1982) indicated the absence of spreading or diffusion during the decay, suggesting the in situ disappearance of the magnetic field. However, it is believed that the decay of sunspots could occur through ohmic diffusion across a current sheet around the sunspot (Gokhale & Zwaan 1972) or through moving magnetic features (MMFs)—magnetic elements that stream from the outer penumbra into the surrounding moat. MMFs are observed as extensions of penumbral filaments (Sainz Dalda & Martínez Pillet 2005; Cabrera Solana et al. 2006; Ravindra 2006; Sainz Dalda & Bellot Rubio 2008), sometimes driven by “Evershed clouds” (Shine et al. 1994; Cabrera Solana et al. 2007, 2008)—velocity patches that appear inside the penumbra

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and propagate radially outward along penumbral filaments. The net flux transported out of a sunspot by MMFs is estimated to be  $(0.4\text{--}6.2) \times 10^{19} \text{ Mx hr}^{-1}$  (Hagenaar & Shine 2005). Thus, a sunspot of  $10^{22} \text{ Mx}$  would disintegrate within one or several weeks if the sunspot flux is removed by MMFs alone. Recent observations of a sunspot region by Kubo et al. (2008) indicate a flux change rate and a transport rate of  $1.2 \times 10^{19} \text{ Mx hr}^{-1}$  and  $2.8 \times 10^{19} \text{ Mx hr}^{-1}$ , respectively. Although MMFs are indicative of sunspot decay (Harvey & Harvey 1973), their origin could be related to the interaction between penumbral field lines, the Evershed mass flow, and the moat flow, not to the decay process itself (Martínez Pillet 2002; Kubo et al. 2007).

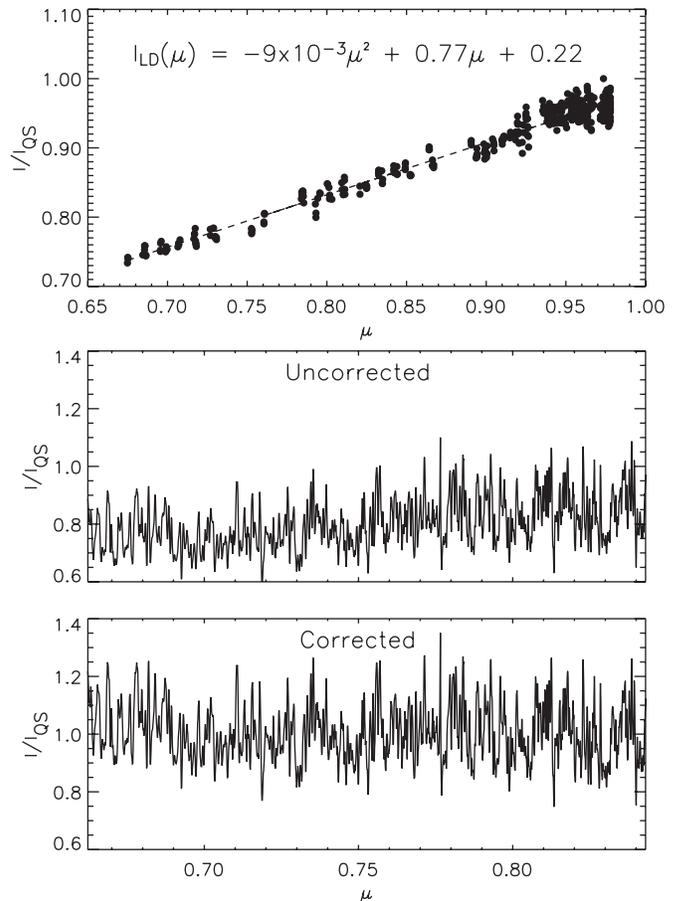
It is important to understand the nature and role of small-scale instabilities in the decay of sunspots. While detailed three-dimensional MHD simulations of active region formation are available (e.g., Cheung et al. 2010), the process of sunspot fragmentation has not been modeled yet and is therefore poorly understood. In this paper, the evolution of a fragmenting sunspot is investigated using high-resolution filtergrams and spectropolarimetric (SP) measurements from the Japanese satellite *Hinode* (Kosugi et al. 2007). The long time sequences provided by *Hinode* are ideally suited for this kind of study. We observed and analyzed the following processes: (1) the formation of a light bridge in the umbra, (2) the fragmentation of the sunspot at the position of the light bridge, (3) the decay of the fragment, (4) the restoration of the penumbra in the parent sunspot, (5) the area and flux decay of the spot, (6) the rotation of the fragment about the parent, and (7) the change in horizontal proper motions caused by the fragmentation. The aim of this work is to identify in situ conditions leading to fragmentation and the different phases of sunspot evolution so as to provide useful constraints for three-dimensional MHD models. Section 2 describes the data processing and Section 3 presents the main results of the analysis. We summarize the sunspot evolution and discuss our findings in Section 4.

## 2. OBSERVATIONS AND DATA PROCESSING

### 2.1. Imaging Data

We employ high-resolution *G*-band filtergrams of NOAA AR 10961 recorded from 2007 June 28 to July 5 by the Broadband Filter Imager of the Solar Optical Telescope (SOT; Tsuneta et al. 2008) on board *Hinode*. The filtergrams have a spatial sampling of  $0''.109$  and were taken in the  $2 \times 2$  binning mode with sizes varying from  $1\text{k} \times 1\text{k}$  to  $2\text{k} \times 1\text{k}$ . The Level-0 *G*-band images were corrected for dark current, flat field, and bad pixels using the *fg\_prep* routine in SolarSoft. Additional corrections are described below.

*Intensity gradients.* The *G*-band data set covers the sunspot's transit across the solar disk, from heliocentric angles ( $\theta$ ) of  $\approx 50^\circ$  close to the East limb to  $\approx 55^\circ$  near the West limb. Consequently, intensity gradients in the image arising from "limb darkening" (Milne 1921; Böhm-Vitense 1997) as well as residual flat-field errors ought to be corrected. A procedure for correcting such intensity variations on *Hinode* *G*-band images has been described by Tan et al. (2009). They selected a horizontal line sampling the quiet-Sun (QS) granulation and used a fifth-degree polynomial to fit the intensity as a function of  $\mu = \cos \theta$  and derive the limb darkening coefficients. A different approach is described below. The filtergrams were first normalized to the average QS intensity of the images taken close to disk center ( $\theta \sim 12^\circ$ ). A  $100 \times 100$  pixel area was then used to compute the mean QS intensity in all the images. The variation of intensity

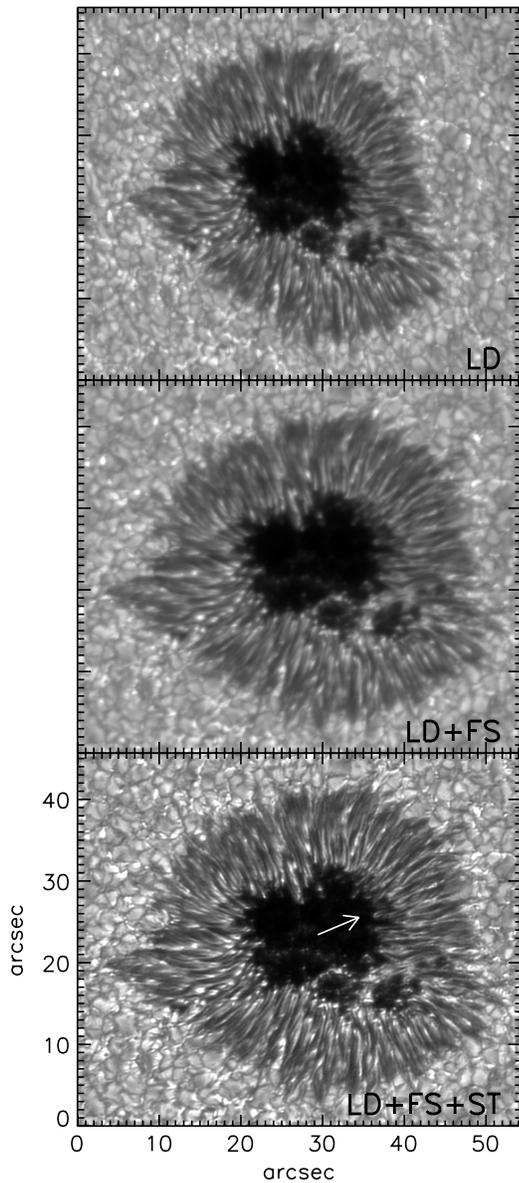


**Figure 1.** Limb darkening and its correction. Top panel: normalized QS intensity averaged over a  $100 \times 100$  pixel area in the Level-1 *G*-band images (black circles) as a function of  $\mu$ . The dashed line is a second-order polynomial fit with the coefficients inscribed in the plot. Middle panel: *G*-band intensity along a horizontal cut in the QS, for a filtergram acquired on 2007 June 28 at  $\theta = 38^\circ$ . Bottom panel: same as above, but with limb darkening removed.

as a function of the  $\mu$  corresponding to the central pixel of the selected area is shown in the top panel of Figure 1. The coefficients given in the figure were obtained from a second-degree polynomial fit. Using these coefficients, the intensity from each pixel was corrected as  $I_{\text{obs}}/I_{\text{LD}}$ . The effect of the correction is illustrated in Figure 1 for a horizontal cut across the QS in a filtergram taken on 2007 June 28, when the sunspot was located at a heliocentric angle of  $38^\circ$ .

*Geometric foreshortening.* Geometric foreshortening produces a projected image of the sunspot as it traverses the solar disk from east to west. A coordinate transformation described by Gary & Hagyard (1990) was performed to render the filtergrams into the heliographic plane ( $x_{\text{H}}, y_{\text{H}}$ ) from the observed image plane ( $x_{\text{I}}, y_{\text{I}}$ ). The middle panel of Figure 2 depicts the sunspot after correction for geometric foreshortening.

*Instrumental stray light.* The presence of stray light in the *Hinode* broadband filtergrams arising from instrumental scattering and its correction have been described by Mathew et al. (2009) using transit observations of Mercury on 2006 November 8. In that paper, the Level 1 images were deconvolved using a point-spread function (PSF) consisting of a weighted linear combination of 4 Gaussians. As a result, the rms contrast of bright points in the QS improved by a factor of 1.7 in the *G* band (430 nm). The PSF derived by Mathew et al. (2009) is used here to remove stray light from the *G*-band filtergrams assuming



**Figure 2.** *G*-band images of sunspot in NOAA AR 10961 located at S14 E38 on 2007 June 28. The acronyms refer to the various corrections applied to the Level-1 images—LD: limb darkening; FS: foreshortening; ST: stray light. All images have been scaled identically. Solar north points up and west is to the right. The white arrow in the bottom panel points to disk center.

that the images corrected for foreshortening are equivalent to those obtained with the telescope pointing to disk center (the conditions of the transit observations). The improvement in the image contrast after removal of stray light is shown in the bottom panel of Figure 2, when the sunspot was located at S14E38 on 2007 June 28. The implication of scattered light on the fine structure of sunspots using *Hinode* filtergrams is detailed in a separate work (Louis et al. 2012).

### 2.2. Spectropolarimetric Data

In addition to the *G*-band filtergrams, we also use observations taken by SOT SP (Lites et al. 2001; Ichimoto et al. 2008). This instrument measured the four Stokes profiles of the iron lines at 630 nm with a spectral sampling of 21.55 mÅ, a pixel size of 0'32, and an exposure time of 1.6 s per slit position (fast map mode). The SP data were corrected for dark current,

flat field, thermal flexures, and instrumental polarization using *sp\_prep.pro*. The continuum maps at 630 nm were also corrected for intensity gradients in a similar manner as the *G*-band images. Maps of field strength, inclination, and azimuth were obtained from a Milne–Eddington inversion of the observed Stokes profiles using the MERLIN code.<sup>5</sup> The procedure described by Crouch et al. (2009) was applied to the vector magnetograms to resolve the 180° azimuth ambiguity. The resulting vector magnetic fields were subsequently transformed to the local reference frame.

## 3. RESULTS

### 3.1. General Description of Sunspot Evolution

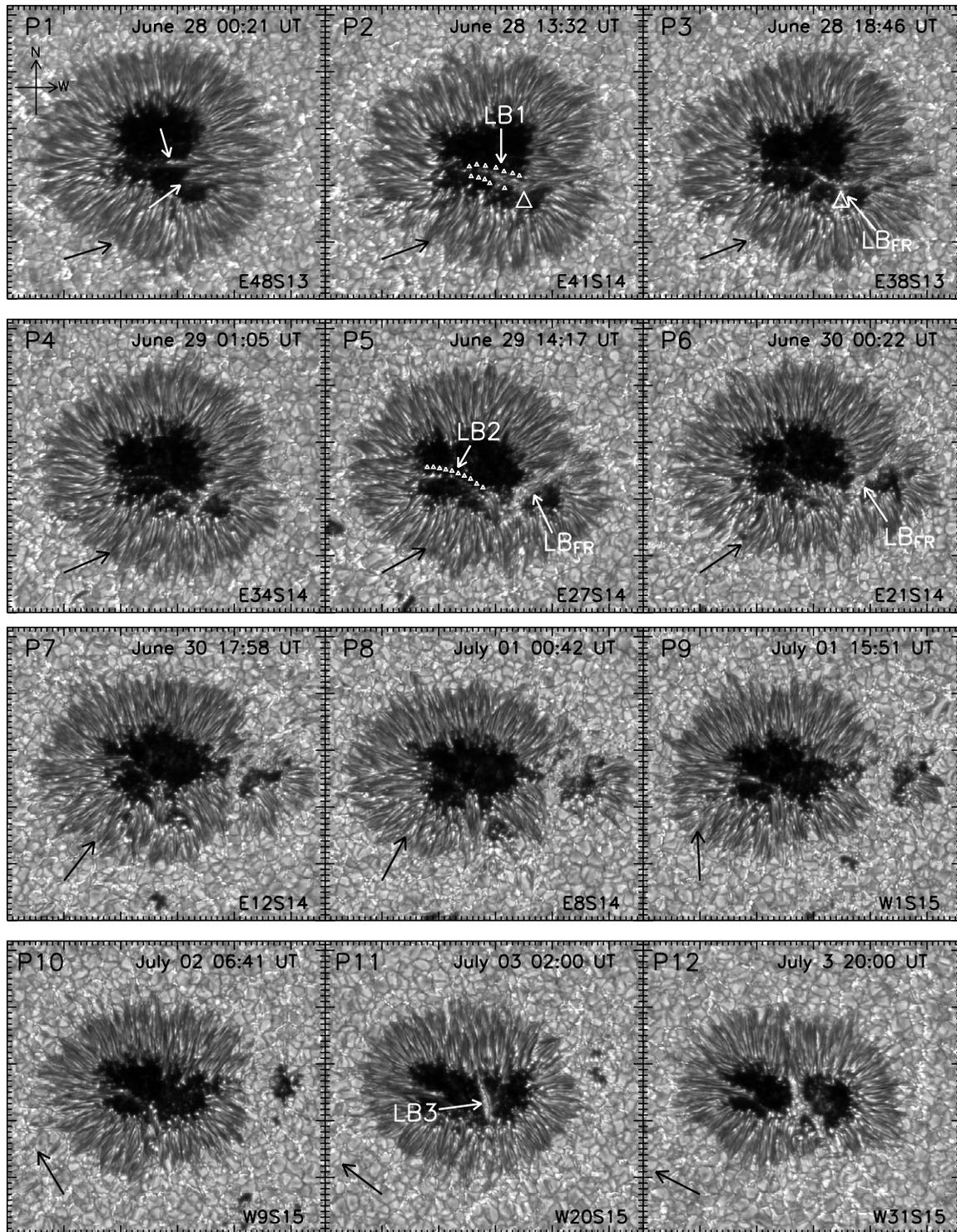
NOAA AR 10961 appeared close to the East limb during the early part of 2007 June 25. The spot moved across the disk at about 14°S and disappeared beyond the West limb on July 7.

Our data set begins on June 28 and shows a fairly regular sunspot with a well-developed penumbra (panel P1 of Figure 3). The western side of the umbra–penumbra boundary (UPB) shows two penumbral filaments extending into the umbra (white arrows in P1). These filaments progress deeper into the umbra eastward giving rise to a light bridge labeled “LB1” in P2 and outlined by tiny white triangles. The apex of the large white triangle indicates a narrow umbral region between the intruding filament and the southern UPB (base of the triangle). This location witnesses the formation of another light bridge, labeled “LB<sub>FR</sub>” in the following panel P3. The time taken for the formation of LB<sub>FR</sub> is approximately 5 hr. During this period, LB1 becomes narrower and fainter in comparison to its newly formed counterpart. Panel P5 shows LB<sub>FR</sub> increasing significantly in width from 485 km to 1295 km by the end of June 28, nearly 20 hr after it was formed. The panel also shows another light bridge “LB2” (outlined with small white triangles), which can be regarded as a successor of LB1. While LB1 was a penumbral structure, LB2 predominantly consisted of umbral dots along its length. We choose to treat LB1 and LB2 as distinct structures based on their different morphology and time of formation.

During the early part of June 30 (P6 in Figure 3), the symmetrical arrangement of the penumbral filaments near the southern section of the UPB gets further disturbed. This location coincides with one of the anchorage points of the light bridge LB<sub>FR</sub>. In addition, there are traces of disruption in the penumbra with the absence of penumbral filaments in the northwestern sector of the sunspot, tracing outward to the photosphere from the entrance of LB<sub>FR</sub>. The part of the sunspot encompassed by the LB<sub>FR</sub> starts to separate from the parent sunspot around 18:00 UT on June 30, as can be seen in P7. At this stage, the two umbral boundaries marked by the light bridge are separated by nearly 1.7 Mm and the morphology of the light bridge resembles the QS photosphere. By the end of June 30, the fragmentation is complete with the sunspot and its fragment being well separated. They remain within a distance of 10" from each other.

The penumbra in the fragment gradually gets depleted as shown in P8–P9. By the early part of July 2, there are only minor traces of it in the remnant pore (P10). The penumbra in the parent sunspot gradually becomes symmetrical and ordered with the exception of the region closest to the fragment where

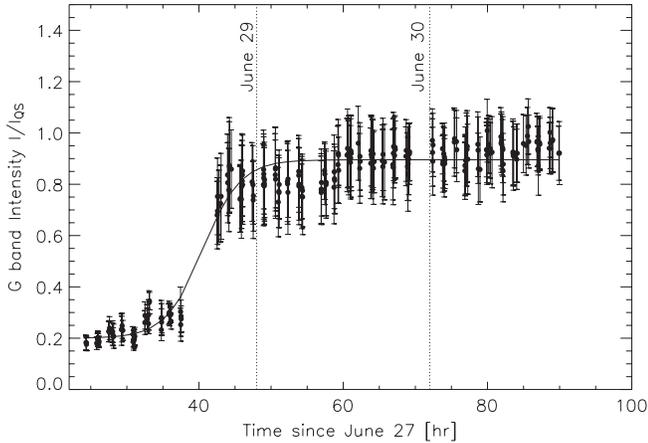
<sup>5</sup> The inversion results can be found at the Community Spectropolarimetric Analysis Center, [http://sot.lmsal.com/data/sot/level2hao\\_new/](http://sot.lmsal.com/data/sot/level2hao_new/).



**Figure 3.** Evolution of sunspot in NOAA AR 10961 from 2007 June 28 to July 3. The position of the sunspot on the solar disk is shown at the bottom of each panel. The image orientation is indicated in the legend at the top left corner of panel P1. Each tickmark corresponds to 1". The black arrow in the bottom left corner points to disk center. See the text for explanation of labels, symbols, and other arrows.

the penumbral filaments are much shorter in length and can be considered rudimentary. The decay of the fragment and restoration of the penumbra in the parent are further elaborated in Section 3.2.3. The decaying pore is only seen as traces of smaller fragments during the early part of July 3 (P11). While the parent spot has a well-developed penumbra uniformly structured all around the umbra by July 3, a conspicuous light bridge

“LB3” is seen running from the northern to southern end of the UPB, nearly splitting the umbra into two equal halves. LB3 progressively becomes broader and exhibits large granules along its length. By June 5, the penumbra close to the southern part of LB3 gets gradually depleted (P12). The sunspot retains this configuration until it advances over the Western limb on 2007 July 7.



**Figure 4.** Evolution of  $G$ -band intensity in the light bridge. The black circles represent the mean  $G$ -band intensity of the region where the light bridge was formed. The vertical bars indicate the rms fluctuations of the  $G$ -band intensity. The solid line is a Boltzmann sigmoid fit.

### 3.2. Evolution of Light Bridge and Penumbrae

This section describes the small-scale changes in the light bridge and penumbrae associated with the evolution and fragmentation of the sunspot.

#### 3.2.1. Formation and Evolution of $LB_{FR}$

The fragmentation of the sunspot is heralded by the formation of a light bridge in the southwest sector of the sunspot umbra ( $LB_{FR}$ , also called LB in this section). The time sequence of  $G$ -band images reveals that the LB formed as a result of a slow incursion of penumbral filaments from the north toward the southern UPB. This intrusion is accompanied by an increase in the intensity by a factor of two over a time duration of 12 hr starting at 1:49 UT on June 28. The speed of the filament motion into the umbra was estimated to be  $\sim 0.08 \text{ km s}^{-1}$ , which is consistent with Katsukawa et al. (2007). During the latter half of June 28, a coherent narrow LB is formed, isolating the smaller umbral core from the parent umbra (P3 of Figure 3). At this time, the LB consists of several bright, grain-like structures with a faint dark boundary separating individual cells (not shown). Intensities of these grains are comparable to those of bright penumbral grains at the southern end of the LB. We put an upper limit of  $\sim 10$  hr from the intrusion of the penumbral filament to the formation of the LB. The following 24 hr witnesses an increase in the width of the LB as well as in the size of bright cells/grains on it, with the motion of the latter closely resembling those of photospheric granules.

The morphological transformation of the LB is accompanied by an increase in its intensity. The temporal variation of the mean intensity of the region where the LB formed is shown in Figure 4. The isolation of the smaller umbral core by the LB is seen as a jump in the  $G$ -band intensity, increasing from 0.28 to  $0.7 I_{QS}$  in 4 hr. The steep increase is followed by a more gradual rise touching near QS—photospheric intensities close to midday of June 29. A Boltzmann sigmoid of the form

$$I(t) = a_0 + \frac{a_1 - a_0}{1 + \exp((a_2 - t)/a_3)}$$

was used to fit the observed time variation of the mean  $G$ -band intensity. The coefficients  $a_0$  and  $a_1$  correspond to the top and bottom of the curve, respectively,  $a_2$  is the time taken to reach halfway between the top and bottom values, and  $a_3$  is the slope

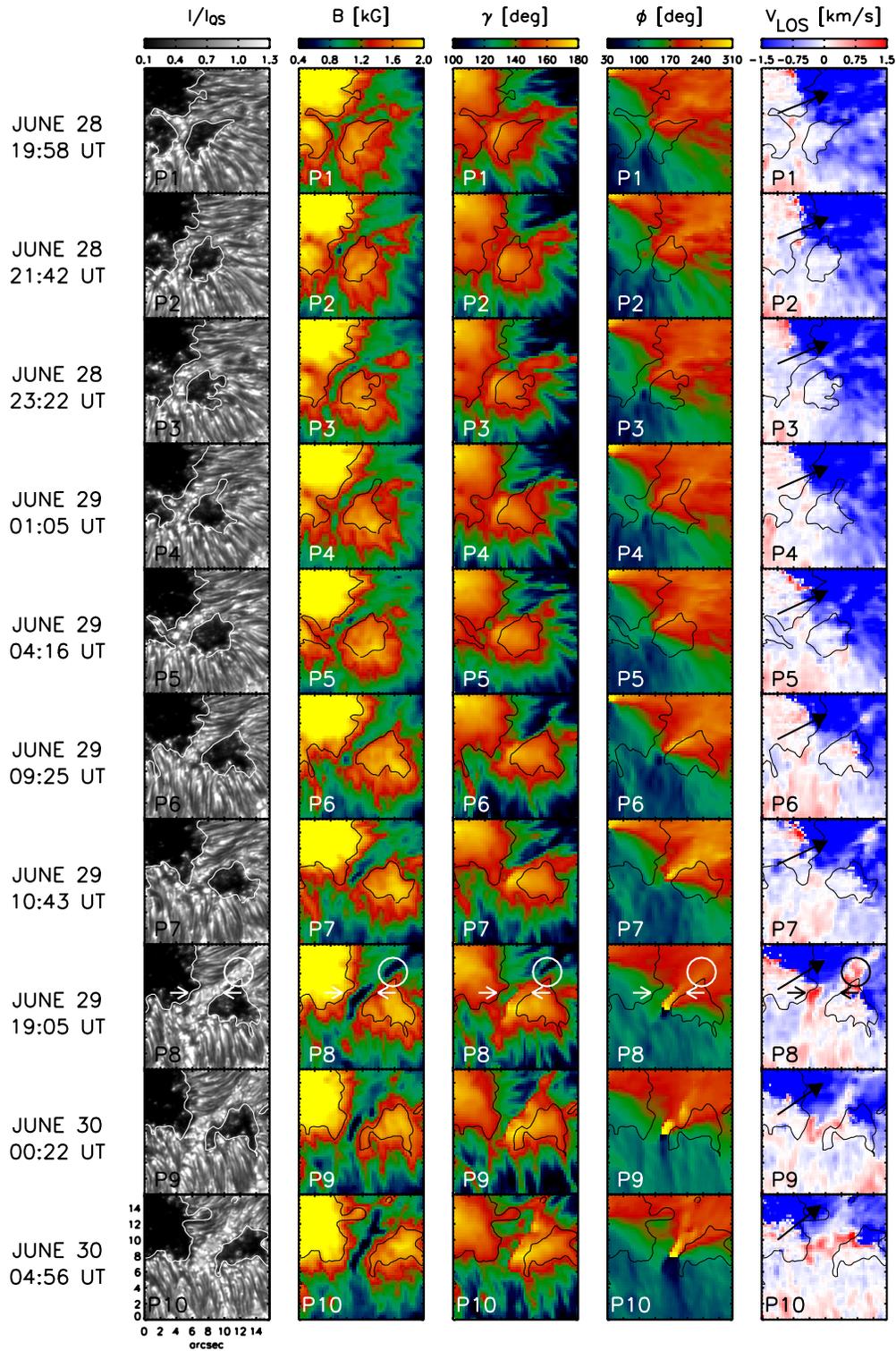
of the function (the larger its value, the shallower the curve). The best fit yields  $a_0 = (0.2 \pm 0.01)I_{QS}$ ,  $a_1 = (0.895 \pm 0.004)I_{QS}$ ,  $a_2 = (16.58 \pm 0.14) \text{ hr}$ , and  $a_3 = (2.6 \pm 0.05) \text{ hr}$ .

Figure 5 shows the temporal evolution of the  $G$ -band intensity, magnetic field strength, field inclination, field azimuth, and line-of-sight (LOS) velocity in the LB. The contours represent the umbral border and outline the LB. The time sequence of maps begins when the LB is already formed. While the main umbra consists of fields in excess of 2 kG, the smaller umbral region bounded by the LB is marginally weaker with an average field strength of 1.5 kG. Compared to the adjacent umbrae, the field strength in the LB is weaker by  $\sim 500 \text{ G}$  as seen in panels P1–P6. The magnetic field in the LB is relatively inclined with respect to the vertical, reaching values of  $120^\circ$ – $150^\circ$ . As a consequence of the smaller core being bounded by the LB, one azimuth center is located close to the central part of the LB on its western edge. The field is predominantly oriented across the LB. This is different to the configuration found in other LBs (e.g., Louis et al. 2008). The LOS velocity in the LB (panels P1–P6) shows weak upflows in the range  $0.1$ – $0.35 \text{ km s}^{-1}$ , as well as downflows of  $0.25$ – $0.57 \text{ km s}^{-1}$ . These flows are quite weak in comparison to the penumbra whose disk and limb sides show velocities greater than  $\pm 1 \text{ km s}^{-1}$ .

Nearly 20 hr after its formation, the field strength along the axis of the LB drops to 650 G (P7). Panels P8–P10 indicate the following: (1) the penumbra close to the northern entrance of the LB appears to be sidelined to the eastern edge while bright photospheric material is seen intruding from the periphery of the spot all the way to the LB, (2) the minimum field strength in the LB reduces to 300 G, (3) a narrow upflowing lane is seen along the central axis of the LB with downflows of  $1.5 \text{ km s}^{-1}$  adjacent to it, and (4) strong downflows are also seen in photospheric regions of the LB near penumbral filaments (circle in P8).

The LOS velocity and the field strength along a cut (arrow head across the LB in P8) indicate a reduction in field strength that coincides with the upflowing lane and is flanked by downflows on either side. The amplitude of Stokes  $V$  in the pixels exhibiting the upflows is weaker than in the downflows. Downflows of up to  $5 \text{ km s}^{-1}$  are seen in the region marked with a circle in P8. The Stokes  $V$  profiles emerging from those pixels show an extended hump characteristic of very large velocities. The magnitude of the downflows was derived from a two-component SIR inversion (Ruiz Cobo & Del Toro Iniesta 1992) with height-independent parameters (except for temperature). The fill fraction of the fast component was estimated to be 16%.

The evolution of the mean and minimum field strengths in the region where the LB formed is depicted in Figure 6. The weakening of these quantities is similar to the trend seen in the  $G$ -band intensity. A Boltzmann sigmoid fit to the field strength reveals that the lower knee of the curve coincides with the weak fields detected in the LB on June 28, 19:05 UT (P8 of Figure 5). The morphology of the LB from this point onward also appears distinct from the nearby penumbral features. The time taken for the field strength (and the continuum intensity) to reach halfway between the bottom and top values is  $\sim 16$  hr, starting from June 28. By comparison, the field strength of the neighboring umbral region remained unaffected during the evolution and morphological transformation of the LB. With a significant field weakening in the LB and the onset of convection that is revealed by the upflows and downflows, the sunspot split close to midday of June 30, nearly 45 hr after the formation of the LB. The following section describes the evolution of the fragment and its subsequent decay.

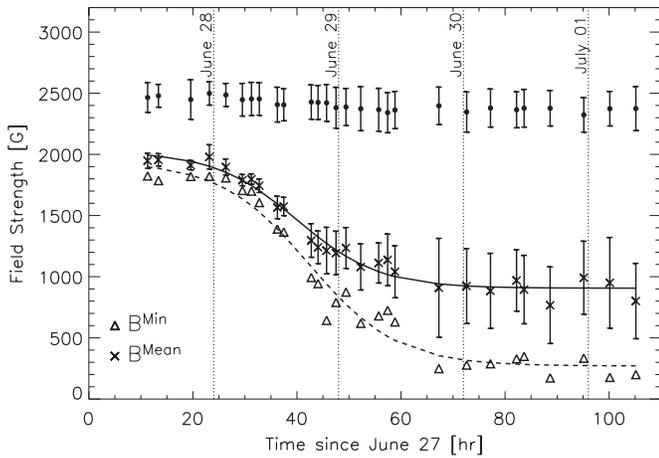


**Figure 5.** Evolution of vector magnetic field and LOS velocity in the light bridge. From left to right: *G*-band intensity, field strength, inclination, azimuth, and LOS velocity. All the maps have been scaled as shown in their respective color bars. Solar east is to the left and north is up. Inclinations and azimuths are expressed in the local reference frame. Azimuths are measured counterclockwise from solar west (right). Thus, an azimuth of  $0^\circ$  implies a magnetic field pointing along the positive *x*-axis. Blue/red correspond to up/downflows in the last column. See the text for explanation of arrow head and region marked by circle in P8. (A color version of this figure is available in the online journal.)

### 3.2.2. Evolution of Fragment

Figure 7 depicts the temporal evolution of the sunspot fragment and the disappearance of its penumbra. The first panel shows the fragment during the early part of June 30 when it was still part of the parent sunspot. The angular span of the penumbra

around the smaller umbral core that broke away is  $\approx 120^\circ$ . Fragmentation commences with the depletion of penumbra close to the anchorage points of the light bridge. The penumbra is seen to reduce symmetrically on either side of the fragment, disappearing clockwise and counterclockwise in the northern and southern penumbra, respectively. The angular span reduces



**Figure 6.** Evolution of field strength in the light bridge. The crosses and triangles represent the mean and minimum field strength of the region where the light bridge was formed. The vertical bars represent the rms fluctuations. The solid and dashed lines are Boltzmann sigmoid fits to the observed mean and minimum field strengths. The black circles correspond to the mean umbral field strength close to the LB.

to half its initial value within nearly 2 hr between June 30, 23:01 UT and July 1, 00:34 UT. The sector-wise disappearance of the penumbra is almost a reversal of the formation process observed by Schlichenmaier et al. (2010b). There is no major trace of the penumbra by July 2, with an isolated pore being the remnant of the fragment, although individual strands of the penumbra can be spotted around the pore (panel 12). There is also a reduction in the mean width of the fragment penumbra, which decreased from  $6''.0 \pm 0''.5$  on June 30 to  $4''.6 \pm 0''.5$  on July 1. The above values correspond to a 24 hr average.

At the time of fragmentation on June 30, 13:39 UT, the fragment has an area of  $92 \text{ Mm}^2$  and decays exponentially with a time constant of  $22.1 \pm 0.2 \text{ hr}$  (bottom panel of Figure 9). Panels 11 and 12 of Figure 7 depict the transition of the fragment to a pore devoid of a penumbra. This occurs between July 1, 20:51 UT and July 2, 01:25 UT. The white contour in panel 11 corresponds to the penumbra–QS boundary (PQB). It has also been overlaid in panel 12, which shows a large pore accompanied by a tiny magnetic patch located near the bottom edge of the white contour. The area of the fragment on July 1, 20:51 UT is  $33 \text{ Mm}^2$ , with umbral and penumbral contributions of 14 and  $19 \text{ Mm}^2$ , respectively. The pore and the patch lie within the white contour representing the fragment 4.5 hr earlier. The total area of the pore and its satellite fragment is  $\sim 21 \text{ Mm}^2$  at that moment. If the decrease in the area is attributed to the disappearance of the penumbra and if the natural decay rate of the fragment is neglected, then the pore and the patch ought to have decayed to at least  $14 \text{ Mm}^2$ , which is inconsistent with the measured value. This suggests that the pore gets replenished with a fraction of the earlier decaying penumbra. During penumbral formation, the area of the umbra remains constant and the growth of the penumbra alone accounts for the increase in spot area (Schlichenmaier et al. 2010b; Rezaei et al. 2012). The process we have witnessed suggests a different behavior for penumbral decay.

One also finds evidence of transient penumbrae that appear and disappear close to the southern edge of the fragment. This is illustrated in panels 5–8 of Figure 7 with ellipses, which shows the appearance of rudimentary penumbrae in locations where photospheric granulation was observed earlier (panels 5 and 6). Such structures do not form stable penumbral filaments.

Transient penumbrae can develop quite rapidly over periods of  $\sim 30$  minutes, but their disappearance is more gradual and occurs over periods of 1.5 hr or more. One can identify two distinct penumbral filaments at the position marked by the ellipse in panel 5 of Figure 7. Panel 7 shows several bright granules clustered together at the edge of the umbra–QS boundary with intensities comparable to those of regular penumbral grains. In a span of 15 minutes, a single penumbral filament is observed in this region. The rudimentary penumbra is not visible after July 1, 5:43 UT. Similar structures have been reported by Schlichenmaier et al. (2010b) during penumbral development, so they are not exclusive to penumbral decay.

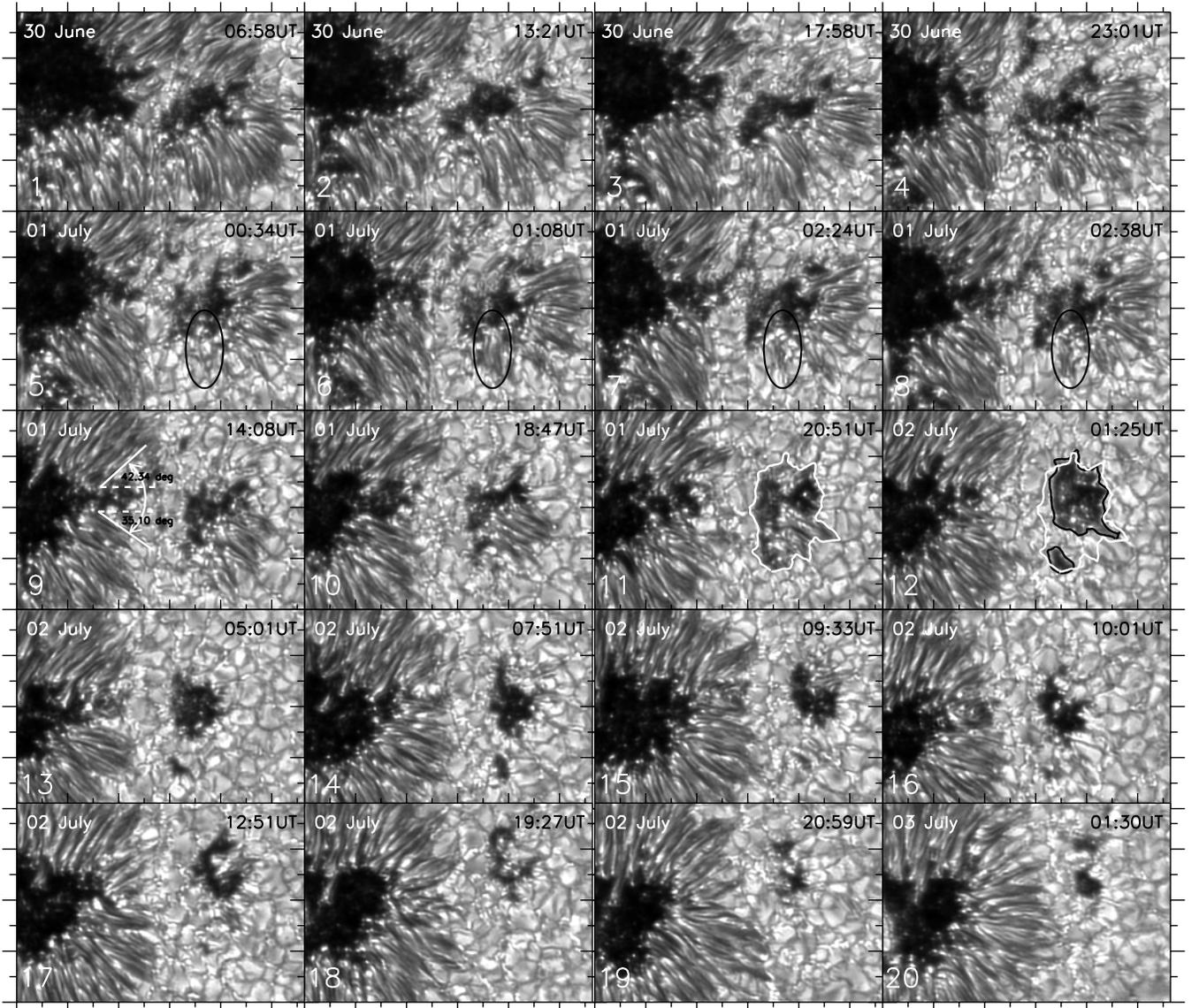
### 3.2.3. Restoration of Penumbra in Parent Sunspot

The separation of the fragment renders a discontinuity in the azimuthal arrangement of the penumbra in the parent sunspot which coincides with  $\text{LB}_{\text{FR}}$ . Panels 9–20 of Figure 7 show the intervening photosphere between the fragment and the parent sunspot. The angular span devoid of the penumbra in the parent is  $\sim 80^\circ$  (panel 9 of Figure 7). The appearance of a complete penumbra in this region is observed only toward the end of July 2, nearly 40 hr after the fragmentation. Furthermore, the length of the penumbral filaments in this region is about 1.5 times shorter than the average filament elsewhere in the sunspot. There is also a strong anticorrelation between the area of the fragment and the width of the parent’s penumbra closest to the fragment. This would suggest that the presence of the fragment hinders the formation of the penumbra in the parent spot.

To determine the threshold of the magnetic field inclination and how it relates to the regeneration of the penumbra in the parent, we compare the field inclination ( $\gamma$ ) at the location of fragmentation with the rest of the sunspot having a regular penumbra, both for the UPB and for the PQB. The inclination in the region of fragmentation is calculated along the portion of the UPB and PQB intensity contours that is closer to the fragment. The left panel of Figure 8 shows the time evolution of  $\gamma$  at the UPB. The mean inclination is  $\sim 150^\circ$  at the location of fragmentation as well as the rest of the sunspot. However, the right panel shows the field becoming more inclined at the PQB, from  $146^\circ$  on July 1, 8:30 UT to  $130^\circ$  on July 2, 7:21 UT. The average value of  $\gamma$  in the rest of the sunspot shows typical values of  $100^\circ$ – $115^\circ$ , i.e.,  $10^\circ$ – $25^\circ$  from the horizontal (the spot is of negative polarity), which is consistent with Jurčák 2011. The increase in inclination coincided with the growth of the penumbra that occurred between 5:00–9:00 UT on July 2 (refer panels 13–16 of Figure 7). Thus, the radial width of the penumbra increased when the average field inclination at the border of the sunspot dropped below  $130^\circ$  (i.e., when the field became within  $40^\circ$  of being purely horizontal).

### 3.3. Area Decay of Sunspot and Fragment

Since the sunspot loses a significant area during the fragmentation process, it is necessary to compare its area and flux decay rates with that of the fragment. Estimation of the sunspot area relies on determining the intensity level separating the penumbra/umbra from the QS. Once the intensity at the interface is known, the area can be calculated from the total number of pixels within the corresponding contour. The intensity of the PQB was determined with the help of the cumulative histogram method of Pettauer & Brandt (1997), which yielded a value of 0.925 for the *G*-band time sequence. In order to obtain a smooth contour, the filtergrams were filtered using a  $7 \times 7$  pixel boxcar.



**Figure 7.** Decay of fragment and formation of penumbra in the parent sunspot in the post-fragmentation phase. The ellipse in panels 5–8 depicts transient penumbra in the decaying fragment. The white contour in panels 11 and 12 outline the fragment on July 1, 20:51 UT, while the black contours in panel 12 mark the pore and the magnetic patch on July 2, 01:25 UT. Each tickmark corresponds to  $5''$ .

The top panel of Figure 9 shows the decay of the sunspot's area with time. The sunspot area decreased from  $697 \text{ Mm}^2$  at 00:21 UT on June 28 to  $636 \text{ Mm}^2$  at 13:39 UT on June 30. A linear fit to the data points yields a decay rate of  $-22 \pm 1.2 \text{ Mm}^2 \text{ day}^{-1}$  or  $-7 \pm 0.4 \text{ MSH day}^{-1}$ <sup>6</sup> with a jump of  $93 \text{ Mm}^2$  at the time of fragmentation.

The bottom panel of Figure 9 shows the area decay of the fragment. In contrast to the parent, the decay appears to be nonlinear. The exponential fit provides a lifetime of 98 hr if the minimum area of the fragment is assumed to be  $2 \text{ Mm}^2$  (the uncertainty of the area measurements). Observations show that the fragment survived for at least 76 hr after fragmentation, as traces of it were seen until the second half of July 3.

Figure 10 shows the temporal evolution of the longitudinal magnetic flux of the parent sunspot and its fragment. Using the inversion results, this quantity is computed as  $\phi = \sum_i f_i B_i \cos \gamma_i S_i$ , where  $f_i$  represents the magnetic filling

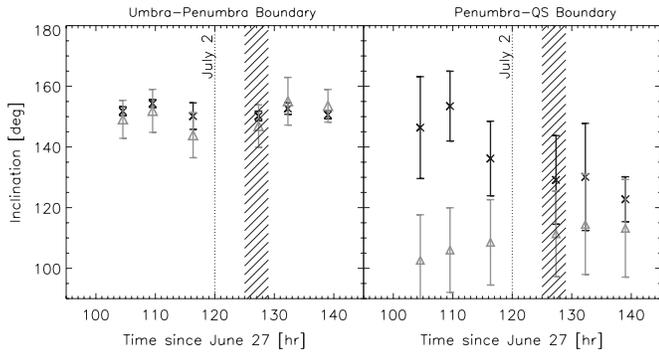
factor,  $B_i$  and  $\gamma_i$  the magnetic field strength and inclination, and  $S_i$  the area of pixel  $i$ . The summation extends over all pixels. The magnetic flux of the parent remains nearly a constant at  $-5.0 \times 10^{21} \text{ Mx}$  over a duration of 61 hr up to the point of fragmentation. The flux is then observed to decay strongly at a rate of  $-4.9 \times 10^{15} \text{ Mx s}^{-1}$  until the disappearance of the fragment. The above value is consistent with that obtained by Deng et al. (2007) and Kubo et al. (2008). After the disappearance of the fragment, the decay rate increases to  $-7.3 \times 10^{15} \text{ Mx s}^{-1}$ .

The magnetic flux of the fragment (bottom panel of Figure 10) amounts to  $-6.3 \times 10^{20} \text{ Mx}$  at the time of separation. By July 2 it decreased down to  $-2.2 \times 10^{20} \text{ Mx}$ . The decay of the penumbra in the fragment could be attributed to the order-of-magnitude deficit in the magnetic flux.

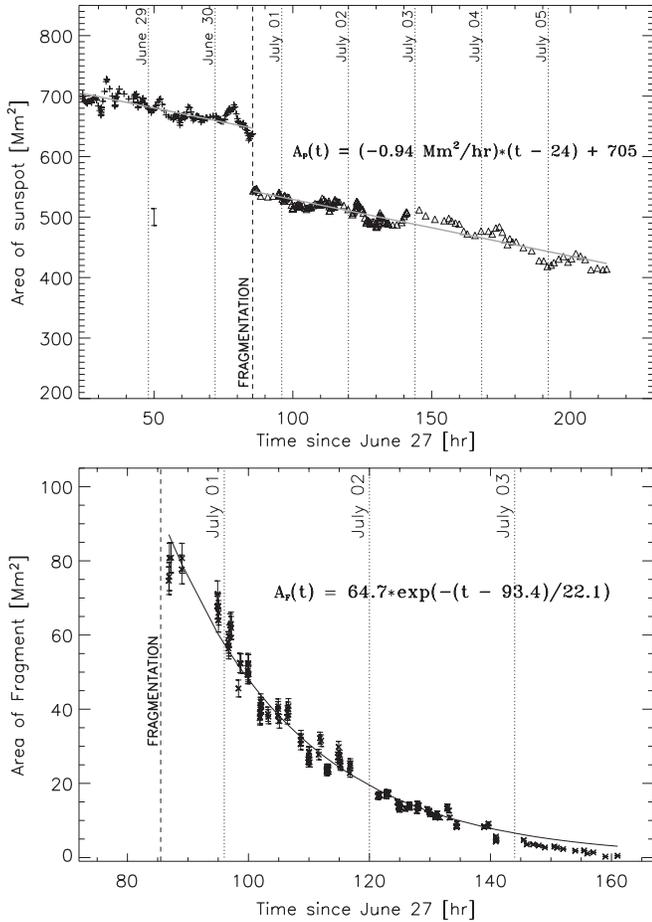
#### 3.4. Rotation of Fragment about the Parent Sunspot

During the evolution of the sunspot, the region enclosed by the light bridge appears to rotate in an anticlockwise direction

<sup>6</sup>  $1 \text{ MSH} = 6.3 \text{ arcsec}^2$ .

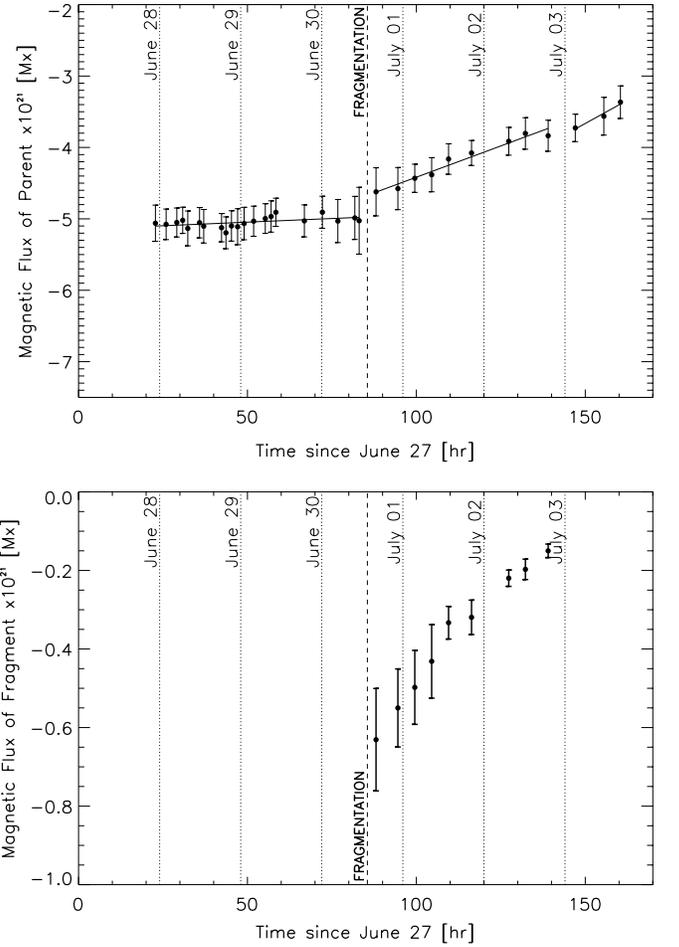


**Figure 8.** Variation of field inclination at the inner and outer penumbral boundaries. Left: time variation of the mean inclination at the umbra–penumbra boundary. The crosses and triangles correspond to the fragmentation site and the rest of the spot, respectively. The vertical error bars represent rms values. The hatched region denotes the interval during which formation of penumbra was seen in the sunspot at the location of fragmentation. Right: same, but for the outer penumbral boundary.



**Figure 9.** Decay of NOAA AR 10961. Top: sunspot area before fragmentation (plus symbols) and after fragmentation (triangles). The area was computed from the  $G$ -band images. Rapid changes in area are neglected by smoothing the data by a three-point average. The typical error in the area estimation is indicated at data point (50, 500). The solid gray line is a linear fit to the data points. Bottom: decay of sunspot fragment. The cross symbols show the fragment area as a function of time. The solid line is an exponential fit to the curve. The vertical bars indicate the error arising from a  $\pm 5\%$  uncertainty in the intensity contour level.

with respect to the main umbral core, as is evident from Figure 3. The rotation described in this section refers to the orbiting motion of the fragment around the center of the parent sunspot (Yan et al. 2008). In order to determine the rotation of the



**Figure 10.** Evolution of the magnetic flux of NOAA AR 10961 derived from SP magnetograms. Top: magnetic flux of the parent. The error bars correspond to an uncertainty of 5% in the iso-intensity contour separating the parent/fragment from the QS. Bottom: same as above, but for magnetic flux of the fragment.

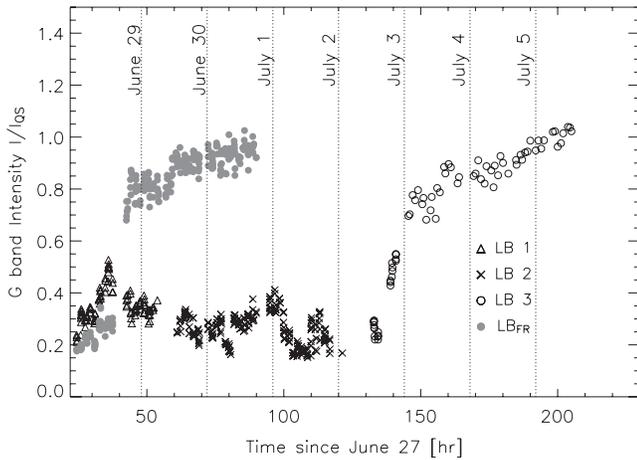
fragment, two points connecting the central axis and the rotating periphery need to be located. The centroid of each umbral core is calculated as

$$x_c = \frac{\sum_{i,j} x_{i,j} I(i,j)}{\sum_{i,j} I(i,j)}, \quad (1)$$

$$y_c = \frac{\sum_{i,j} y_{i,j} I(i,j)}{\sum_{i,j} I(i,j)}. \quad (2)$$

The rotation angle is measured counterclockwise from solar south (vertically downward) to the line joining the two centroids. Over the course of five days the fragment rotated about the parent by nearly  $60^\circ$  with an average rotation speed of  $13^\circ \text{ day}^{-1}$ . This is consistent with that obtained by Brown et al. (2003) for AR 9280.<sup>7</sup> The separation between the parent and fragment, estimated as the distance between their centroid positions, increased almost linearly from approximately  $10''$ – $23''$  in five days. As the rotation of the fragment is more pronounced than its parent, it would suggest that the sunspot as a whole is comprised of several individual magnetic strands rooted to a common flux system rooted deeper in the photosphere (Bello González et al. 2012).

<sup>7</sup> Incidentally, AR 9280 elongated and separated into two pores during its transit.



**Figure 11.** Temporal evolution of the mean  $G$ -band intensity of all sunspot LBs observed in NOAA AR 10961.

### 3.5. Role of other LBs in Sunspot Evolution

In addition to  $LB_{FR}$ , there were several other LBs that were observed during the sunspot’s lifetime but did not cause fragmentation. This section compares various physical properties of the LBs seen in the active region to ascertain why fragmentation occurred. Figure 11 shows the time variation of  $G$ -band intensity in LB1, LB2, and LB3, the three LBs described in Section 3.1 and marked in Figure 3. To allow comparisons, the curve corresponding to  $LB_{FR}$  is also included. The  $G$ -band intensities of LB1 and LB2 remain within  $0.53 I_{QS}$  and do not exhibit any specific trend as a function of time. By comparison, the light curve of LB3 appears very similar to that of  $LB_{FR}$ , with the intensity rising from umbral to near photospheric values in about 35 hr. The photometric properties of both LB3 and  $LB_{FR}$  are consistent with each other as indicated in Table 1.

The average field strengths in LB1 and LB2 exceed 1.5 kG, with a minimum field strength of more than 1.1 kG. In LB3, the corresponding values are 1.6 and 0.7 kG, respectively. This is comparable to  $LB_{FR}$ , whose average and minimum field strengths are 1.3 and 0.6 kG, 34 hr prior to fragmentation. In addition, the field in LB1 and LB2 is mostly vertical with a mean inclination of about  $160^\circ$ . This can be attributed to the fact that LB1 and LB2 morphologically resemble umbral dots and penumbral filament intrusions.

The LB lifetimes range from 1 to 2.5 days. LB3 in particular is quite long lived, since it was observed to be intact during the sunspot’s transit across the Western limb on 2007 July 7, implying a duration of more than 80 hr. The lifetime of  $LB_{FR}$  is estimated to be  $\sim 45$  hr, from the moment it isolated the umbral core until the sunspot fragmentation. The area of the LBs in NOAA AR 10961 varied between  $6.3$  and  $8.6 \text{ Mm}^2$ . While LB3 had an area of  $7 \text{ Mm}^2$  on July 3 at 3:03 UT, this increased significantly to  $11 \text{ Mm}^2$  nearly 8 hr later as the LB developed an additional arm connecting to the penumbra at its southern end. The fraction of the umbra isolated by the LBs is shown in the last row of Table 1. The area bounded by LB1 and LB2 represents 24% and 16% of the total umbral area, while LB3 divided the umbra nearly in half. By comparison,  $LB_{FR}$  segregated only 6% of the umbra. Although LB3 had physical characteristics similar to  $LB_{FR}$ , it did not fragment the spot.

### 3.6. Evolution of Horizontal Motions

In this section we describe the horizontal proper motions of intensity features, in and around the sunspot, and how they

**Table 1**  
Physical Properties of Sunspot LBs in NOAA AR 10961

Parameter		LB1	LB2	LB3	$LB_{FR}$
$G$ -band intensity ( $I_{QS}$ )	Min	0.21	0.15	0.22	0.18
	Max	0.53	0.41	1.09	1.02
	Mean	0.35	0.26	0.75	0.73
$B$ (G)	Min	1160	1390	680	630
	Max	2150	2130	2330	1880
	Mean	1720	1800	1620	1280
$\gamma$ (deg)	Min	139	148	108	110
	Max	172	168	166	176
	Mean	158	159	135	150
Lifetime (hr)		30	60	>80	45
Area ( $\text{Mm}^2$ )		7.3	6.3	11.1	$8.6^a$
Umbral fraction (%)		23.7	15.6	38.4	5.9

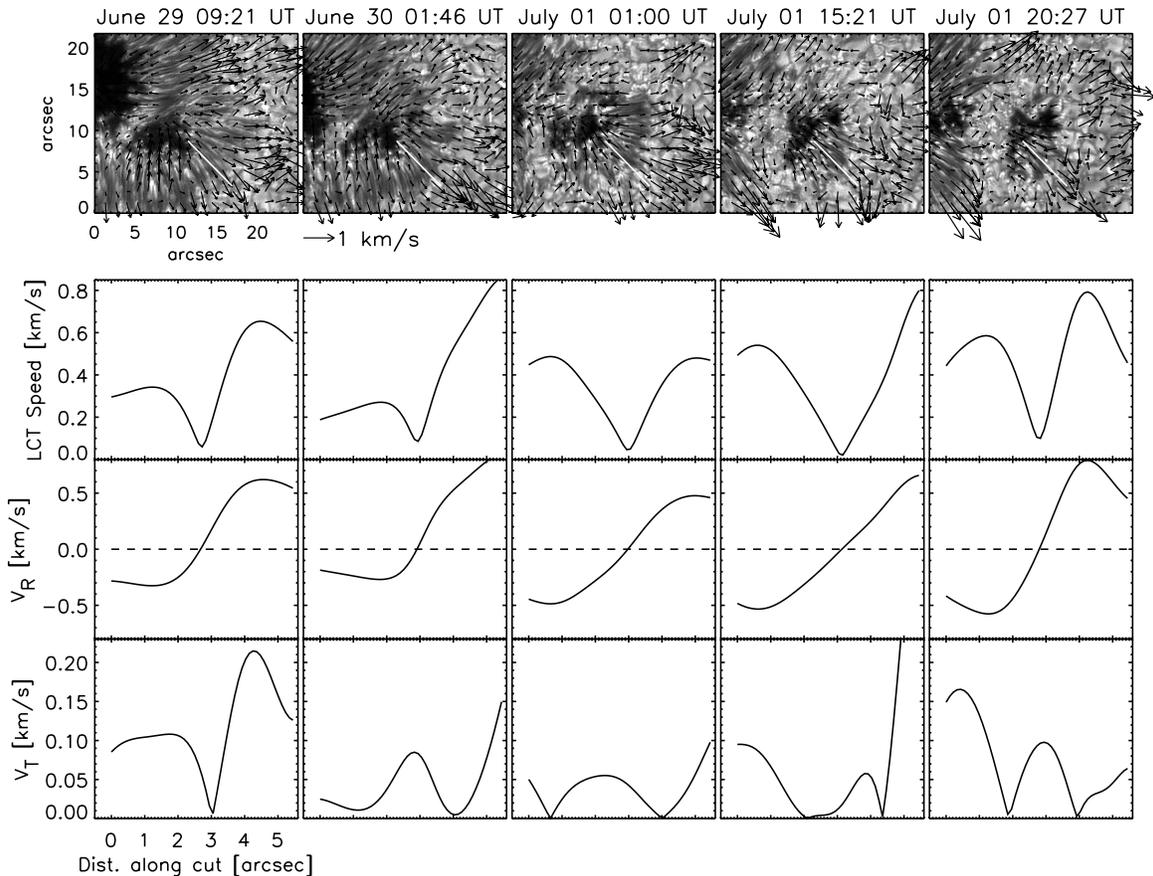
**Note.** <sup>a</sup> Area estimated on 18:31 UT 2007 June 28 after LB was formed.

were affected by the fragmentation process. The horizontal motions were determined using local correlation tracking (LCT; November 1986; November & Simon 1988; Fisher & Welsch 2008; Welsch et al. 2004), a technique which computes the relative displacement of small sub-regions centered on a particular pixel with sub-pixel accuracy using cross-correlation techniques. The sub-regions are apodized by a Gaussian window whose FWHM is roughly the size of the structures that need to be tracked. Knowing the displacement and the time interval, the horizontal speed for each pixel can be determined. First, the aligned  $G$ -band data sets are filtered for acoustic waves using a phase velocity cutoff value of  $6 \text{ km s}^{-1}$  (Title et al. 1989). The filtered images are then subject to the tracking routine. After experimentation, an apodizing window of  $1''$  width and a time difference of 2 minutes between two images were chosen. To reduce the noise in the measurements, five velocity images were averaged. The LCT input parameters are similar to the ones used by Vargas Domínguez et al. (2010) and Verma & Denker (2011) on *Hinode*  $G$ -band filtergrams.

Horizontal proper motions in a sunspot are characterized by the presence of inward and outward directed flows originating from a region of divergence, or dividing line (DL; Sobotka et al. 1999; Sobotka & Sütterlin 2001), located at  $0.6$ – $0.7$  penumbral radii. Since these motions change significantly in the region of fragmentation, remaining fairly uniform in the rest of the AR, we restrict our analysis of the horizontal motions in and around the fragment in the following sections.

### 3.7. Continuance of Horizontal Motions in Fragment

As the width of  $LB_{FR}$  increases on June 30 and the separation between the sunspot and its fragment widens, the average speeds of proper motions in the intermediate granulation region are seen to be less than  $100 \text{ m s}^{-1}$ . Radial motions are seen in the penumbra of the fragment even after separating from the parent. Figure 12 shows horizontal velocity maps at different stages of evolution of the fragment. The horizontal velocity vectors have been overlaid on the  $G$ -band images in the top panel. Below each  $G$ -band snapshot we plot the variation of the horizontal speed, as well as the radial and tangential components of the velocity along a radial cut in the penumbra (white line). The radial component is measured with respect to the cut and is negative and positive for inward and outward motions, respectively. Figure 12 illustrates the following: (1) the inward and outward motions from the DL is present even



**Figure 12.** Continuation of the horizontal motions in the decaying fragment after separation from the parent spot. First row: horizontal motion maps on different days overlaid on the *G*-band images. Arrows have been drawn only for pixels where the speed is greater than  $70 \text{ m s}^{-1}$ . The white line represents the cut along which the horizontal speed and the radial ( $V_R$ ) and tangential ( $V_T$ ) components of the velocity vector are shown in the second, third, and fourth rows, respectively.

after fragmentation, (2) such a radial pattern is only observed in penumbral filaments wherever present in the decaying fragment, and (3) inward motions of  $100\text{--}150 \text{ m s}^{-1}$  are seen at the location of fragmentation which also lacks a penumbra. These inward motions are weaker than those typically seen in the inner penumbra of the sunspot.

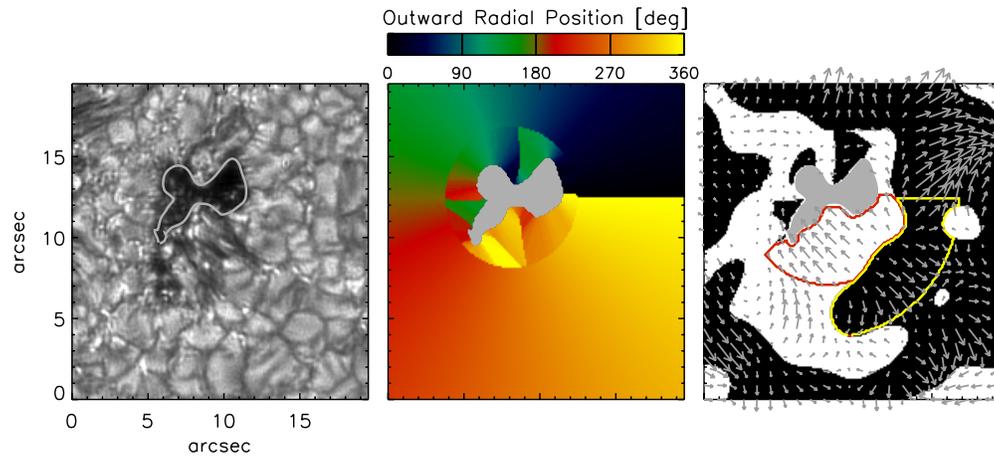
### 3.7.1. Horizontal Motions during the Pore Phase

Here we compare the nature of the horizontal motions in the pore on July 2 and in the fragment comprising a penumbra a day earlier. In order to isolate regions of inward and outward motions surrounding the pore, the angle  $\beta$  between the position vector and the velocity vector at each pixel was determined. The former is directed from the pore centroid or the pore-QS boundary outward. The azimuth of the position vector is the angle it makes with solar west, increasing in the counterclockwise direction. For radial distances greater than  $4''$ , the azimuth is calculated using the pixel location and the pore's centroid. For pixels closer to the pore, the azimuth is assigned the value of the point on the pore-QS boundary lying nearest to the pixel of interest. The above procedure is similar to the one adopted by Vargas Domínguez et al. (2010). Once the position angle is known,  $\beta$  is calculated such that it is  $180^\circ$  and  $0^\circ$  for pure radial and outward motions, respectively. Binary maps are then obtained assigning 1 for  $90^\circ < \beta < 180^\circ$  and 0 for  $0^\circ < \beta < 90^\circ$ . The middle panel of Figure 13 shows the azimuths of the position vector on June 1 at 20:39 UT, while the right panel displays the binary image depicting inward and outward motions. The

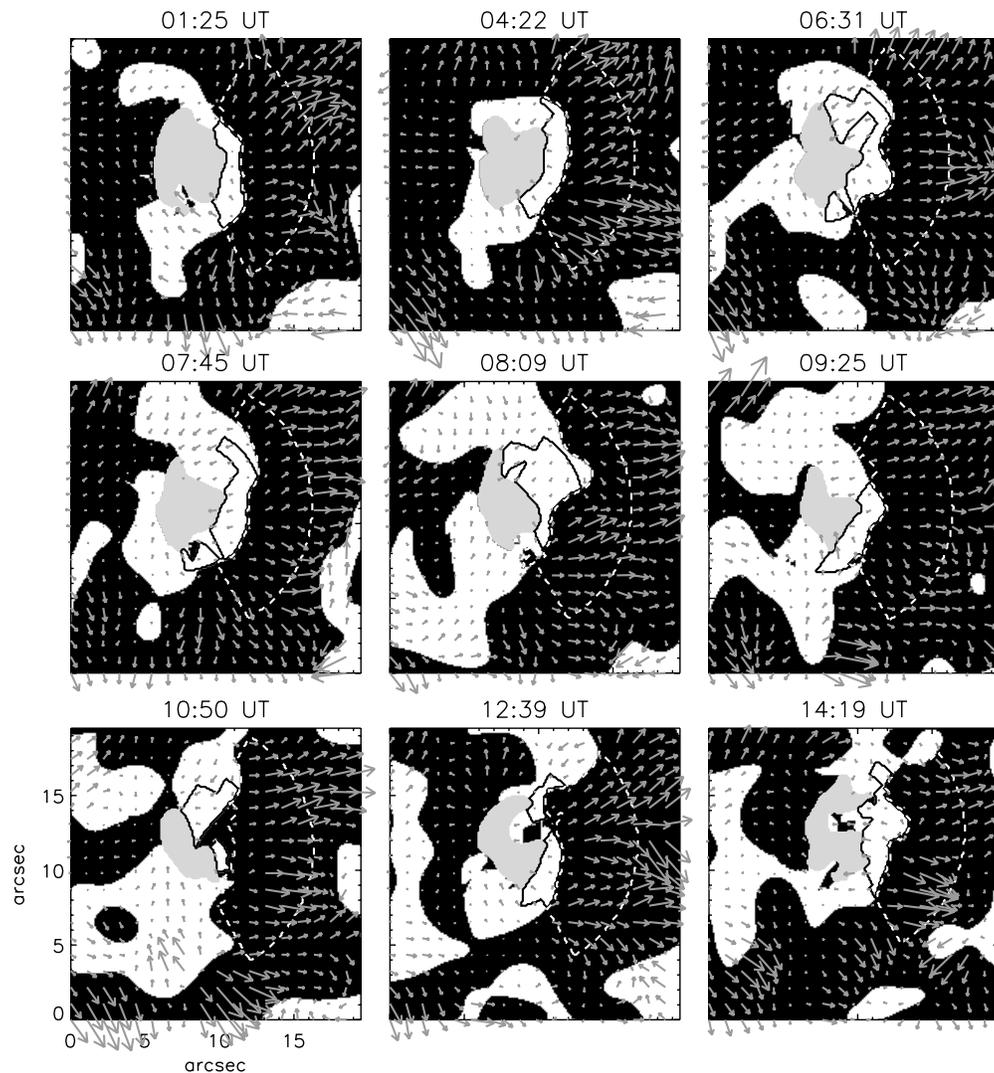
red and yellow contours span the azimuth range from  $220^\circ$  to  $360^\circ$ .

Figure 14 shows a sequence of binary maps on July 2, when the fragment was reduced to a pore. The pore has been masked in light gray and the gray arrows represent the velocity vectors. The figure clearly shows a region of inward motions surrounding the pore on the western side. The inward motions are interrupted by slower outward motions on the eastern side, possibly due to the presence of the larger parent sunspot. The black contour represents a region of inward motions having a  $\pm 60^\circ$ -span in azimuth around the western edge of the pore. This is the region selected for further analysis. Since the binary maps on July 2 do not show clear differences with the ones on July 1 when the penumbra was still present, we compare the magnitude of the inward and outward motions during the pore phase with that of the previous day when the penumbra was intact.

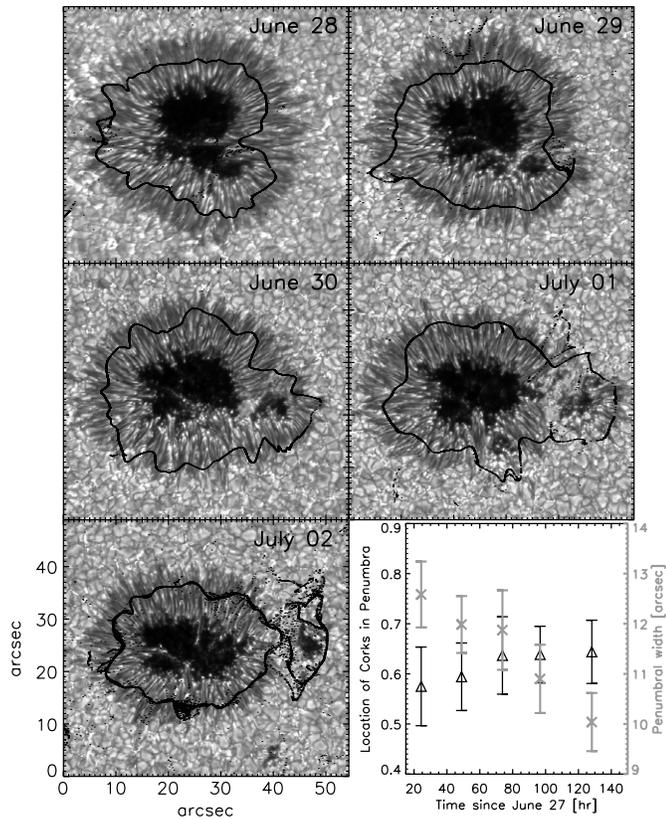
In the pore, the mean speed of the inward motions is  $0.1\text{--}0.3 \text{ km s}^{-1}$ , while the fragment has  $\sim 0.4 \text{ km s}^{-1}$  on July 1. The radial distance of the DL from the pore boundary is  $\sim 1''.6$ . In the case of the fragment, the DL lies further out at  $3''$ . The DL is located at even larger radial distances in the parent sunspot that has a well-formed penumbra, up to  $6''$  from the umbra. These results are in good agreement with the findings of Deng et al. (2007) and Vargas Domínguez et al. (2010). The differences between the horizontal motions surrounding the fragment when it is a pore and when it shows a penumbra are (1) the speed of the inward motions, which is systematically smaller by  $\sim 0.25 \text{ km s}^{-1}$  in the pore as revealed by the LCT



**Figure 13.** Horizontal motions around fragment with final traces of penumbra. Left panel: *G*-band image of fragment on July 1 at 20:39 UT. The gray contours represent the umbra–penumbra boundary of the fragment. Middle panel: azimuth corresponding to the position vector measured in the outward direction from fragment centroid or umbra–penumbra boundary. The azimuth increases in the anticlockwise direction with 0° and 180° pointing toward solar west and east as shown in the color bar. Right panel: binary mask representing inward (white) and outward (black) radial motions surrounding the fragment. The horizontal velocity vectors are indicated with gray arrows. The thick red/yellow contours span the azimuth range from 220° to 360° and represent the inflow/outflow regions. (A color version of this figure is available in the online journal.)



**Figure 14.** Binary image representing inward and outward radial motions surrounding the pore on July 2. The horizontal velocity vectors are indicated with gray arrows except for the pore which has been shaded in light gray. The black contour outlined on the western interface of the pore corresponds to the inflow region chosen for analysis. The contour has a  $\pm 60^\circ$  azimuthal span centered at 0° azimuth.



**Figure 15.** Cork maps depicting the evolution of the sunspot and its horizontal flow field. Each cork map was constructed using 25,600 artificial tracer particles, shown in black dots, that were tracked backward in time for 24 hr. Bottom right panel: evolution of the position of the DL (black triangles) in terms of the normalized penumbral distance and penumbral width (gray crosses). The vertical bars correspond to the rms value.

maps on July 2; and (2) the location of the DL separating the inward and outward motions, which lies much closer to the pore boundary in the former.

### 3.7.2. Persistent Divergence Locations

In the previous section, the inward and outward motions surrounding the pore were compared with those of the fragment on the previous day. A common feature that was observed in both cases is the region of divergence separating the two types of motions, although their nature and properties are distinct from one another. In order to identify the persistent locations of divergence as the sunspot evolved we use the method of corks described below.

We start out with 25,600 artificial trace particles scattered uniformly and allowed to advect in the flow field. These corks are traced backward in time for 24 hr. As a consequence of the long time duration, the corks appear to collect at the strongest centers of divergence which in this case lie in the outer penumbra of the sunspot. The resulting locations of divergent motion are shown in Figure 15. Furthermore, the shape of the DL gets stretched in the direction of separation as the fragmentation commences. Note that the DL remains fairly constant in the rest of the sunspot as described earlier. The regions where the DL appears discontinuous coincide with places which lack a penumbra or where photospheric granulation intervenes. This is seen particularly on July 1 near the southern region of the parent sunspot and the anchorage points of the light bridge where the fragment separated from the spot. The existence of a DL in the

penumbra of the fragment even after the break up process is consistent with the observations shown in Figure 12. On July 2, the DL closed in around the parent spot, hinting at the re-establishment of the penumbra. The DL was also present in the pore on July 2, but at a smaller distance from the pore-QS boundary than in the sunspot.

The bottom right panel of the figure indicates that as the sunspot evolves, the DL has a weak tendency to move outward. The position of the DL in terms of the normalized penumbral distance, varies from  $0.58 \pm 0.16$  to  $0.64 \pm 0.13$  from June 28 to July 2 respectively. There is also a systematic reduction in the width of the penumbra as the sunspot decays, decreasing from  $12.6 \pm 1.3$  arcsec to  $10.0 \pm 1.2$  arcsec over a course of five days. Those pixels that lie outside the sunspot or close to the fragment umbral core were not considered in the above calculation.

## 4. DISCUSSION

We followed the evolution of a sunspot in NOAA AR 10961 for eight days in 2007 June/July using high-resolution imaging and SP observations from *Hinode*. The observations show the formation of an LB that fragmented the sunspot into two parts. The LB formed as a consequence of penumbral filament intrusion into the umbra. The fragmentation process began with the depletion of the penumbra at the anchorage points of the light bridge. Consequently, the characteristic pattern of proper motions in and around sunspots was also disrupted. The presence of the fragment inhibited its restoration. This could be due to the inability of magnetic fields to rise to the surface by buoyancy because of the intervening convective wall that extended beneath the photosphere between the parent sunspot and the fragment.

The precursors to penumbral formation are elongated granules that develop structures resembling penumbral filaments close to the edge of the sunspot boundary, as observed by Schlichenmaier et al. (2010a, 2010b) and Lim et al. (2011). These features could be tiny  $\Omega$ -shaped magnetic loops. In addition, one finds the formation of the penumbra in the sunspot to be closely related to the magnetic field inclination at the PQB as seen in the post fragmentation phase.

So, does the presence of a light bridge herald the fragmentation of a sunspot and what are the conditions for this to happen? The significant field strength reduction in the LB along with the presence of granulation is suggestive of strong convection which might have triggered the expulsion and fragmentation of the smaller umbral core. Katsukawa et al. (2007) point out that the process of LB formation, fragments the sunspot by injecting hot, weakly magnetized gas into the umbra. However, the LB analyzed by Katsukawa et al. (2007) remained an extension of the penumbra and did not separate individual umbral cores. Furthermore, the intensity and field strength did not evolve to QS values and the sunspot remained a single coherent structure during its transit. This does not imply however, that LBs which have matured into photospheric structures will in turn fragment the sunspot. LB3, which was seen in the sunspot after the fragmentation, had similar characteristics to LB<sub>FR</sub> and in this regard it ought to have split the sunspot within a timescale of three to four days after it had reached QS conditions, but it did not. The failure of fragmentation might be attributed to the fraction of the umbra that LBs isolate in general. This could decide whether individual fragments can remain stable to convective motions. In our case, we cannot rule out the possibility that the sunspot could have fragmented soon after its passage across the western limb. Although the presence of photospheric conditions in a

light bridge is a necessary condition for fragmentation, it is not a sufficient one.

The fragmentation of sunspots, seen as a transition of their sub-structures, demonstrates a strong interplay between convection and magnetic fields extending over a large range of spatial and temporal scales. The next step would be to carry out a similar investigation on the formation of sunspots, which would be crucial to understanding the organization and stability of magnetic fields at the solar surface.

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